

WIs and Baryonic States in $\mathcal{N}=1$ SUSY Yang-Mills theory on the lattice

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The Model

$\mathcal{N}=1$ SUSY Yang-Mills Theory

The action:

$$S_{\text{SYM}} = \text{Re} \int d^4x d^2\theta \text{Tr}[W^\alpha W_\alpha] = \int d^4x \left\{ \frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a + \frac{1}{2} \bar{\lambda}^a \gamma_\mu \mathcal{D}_\mu \lambda^a \right\}$$

- Field strength tensor: $F_{\mu\nu} = -ig F_{\mu\nu}^a T^a = \partial_\mu A_\nu - \partial_\nu A_\mu + [A_\mu, A_\nu]$
- Covariant derivative in adjoint representation: $\mathcal{D}_\mu \lambda^a = \partial_\mu \lambda^a + g f_{abc} A_\mu^b \lambda^c, \quad a = 1, \dots, N_c^2 - 1$
- Vector supermultiplet:
 - Gauge field $A_\mu^a(x)$, "Gluon"
 - Majorana-spinor field $\lambda^a(x)$, $\bar{\lambda} = \lambda^T C$, "Gluino"
- SUSY transformations (on-shell):

$$\delta A_\mu^a = -2g \bar{\lambda}^a \gamma_\mu \varepsilon, \\ \delta \lambda^a = -\frac{i}{g} \sigma_{\mu\nu} F_{\mu\nu}^a \varepsilon$$

- In contrast to QCD:
 - λ is Majorana spinor field, " $N_f = \frac{1}{2}$ "
 - adjoint representation of $SU(N_c)$
- Gluino mass term $m_{\tilde{g}} \bar{\lambda}^a \lambda^a$ breaks SUSY softly

Motivation

- SYM: simplest model with SUSY and local gauge invariance
- Part of the supersymmetrically extended Standard Model
- Possible connection to ordinary QCD
- Similar to QCD:
 - Asymptotic freedom
 - Confinement
 - Numerical lattice simulation of bound states

Solution of non-perturbative Problems:

- Spontaneous breaking of chiral symmetry $Z_{2N_c} \rightarrow Z_2$
↳ Gluino condensate $\langle \lambda \lambda \rangle \neq 0$
- Spectrum of bound states → Supermultiplets
- Confinement of static quarks
- Spontaneous breaking of SUSY?
- SUSY restoration on the lattice
- Check predictions from effective Lagrangeans (Veneziano, Yankielowicz ...)

Spontaneous breaking of chiral symmetry

$U(1)_\lambda: \lambda' = e^{-i\varphi \gamma_5} \lambda, \bar{\lambda}' = \bar{\lambda} e^{-i\varphi \gamma_5} \leftrightarrow$ R-symmetry, $J_\mu = \bar{\lambda} \gamma_\mu \gamma_5 \lambda$

Anomaly: $\partial^\mu J_\mu = \frac{N_c g^2}{32\pi^2} \epsilon^{\mu\nu\rho\sigma} F_{\mu\nu}^a F_{\rho\sigma}^a$ breaks $U(1)_\lambda \rightarrow Z_{2N_c}$

Spontaneous breaking $Z_{2N_c} \rightarrow Z_2$
by Gluino condensate $\langle \lambda \lambda \rangle \neq 0$
↳ first order phase transition at $m_{\tilde{g}} = 0$

$N_c = 2: \langle \lambda \lambda \rangle = \pm C \Lambda^3$

Spectrum of bound states

Expect colour neutral bound states of gluons and gluinos
Predictions from effective Lagrangeans:

Chiral supermultiplet (Veneziano, Yankielowicz)

- 0^- gluinoball $a - \eta' \sim \bar{\lambda} \gamma_5 \lambda$
- 0^+ gluinoball $a - f_0 \sim \bar{\lambda} \lambda$
- spin $\frac{1}{2}$ gluino-glueball $\sim \sigma_{\mu\nu} \text{Tr}(F_{\mu\nu} \lambda)$

Generalization (Farrar, Gabadadze, Schwetz):

additional chiral supermultiplet

- 0^- glueball
- 0^+ glueball
- gluino-glueball

possible mixing and Baryonic states

Simulations

SUSY on the lattice

Lattice breaks SUSY. Restoration in the continuum limit?
Curci, Veneziano: use Wilson action, search for continuum limit with SUSY

$$S = -\frac{\beta}{N_c} \sum_p \text{Re} \text{Tr } U_p \\ + \frac{1}{2} \sum_x \left\{ \bar{\lambda}_x^a \lambda_x^a - \kappa \sum_{\mu=1}^4 \left[\bar{\lambda}_{x+\hat{\mu}}^a V_{ab,x\mu} (1 + \gamma_\mu) \lambda_x^b + \bar{\lambda}_x^a V_{ab,x\mu}^t (1 - \gamma_\mu) \lambda_{x+\hat{\mu}}^b \right] \right\} \\ \beta = \frac{2N_c}{g^2}, \quad \kappa = \frac{1}{2m_0 + 8} \quad \text{hopping parameter}, \quad m_0: \text{bare gluino mass}$$

$$V_{ab,x\mu} = 2 \text{Tr} (U_{x\mu}^\dagger T_a U_{x\mu} T_b), \quad \text{adjoint link variables}$$

We study gauge group SU(3).

Fermion integration

Fermionic action

$$S_f = \frac{1}{2} \bar{\lambda} Q \lambda = \frac{1}{2} \lambda M \lambda, \quad M \equiv CQ$$

Pfaffian

$$\int [d\lambda] e^{-S_f} = \text{Pf}(M) = \pm \sqrt{\det Q}$$

Effective gauge field action

$$S_{\text{eff}} = -\frac{\beta}{N_c} \sum_p \text{Re} \text{Tr } U_p - \frac{1}{2} \log \det Q[U]$$

Include sign Pf(M) in the observables.

Monte Carlo algorithm

- Monte Carlo Simulation has been used to generate configurations.
- Each configuration contains numerical values of link variables (U).
- These configurations are used to compute correlation functions.

Sign Problem:
monitoring of sign Pf(M)

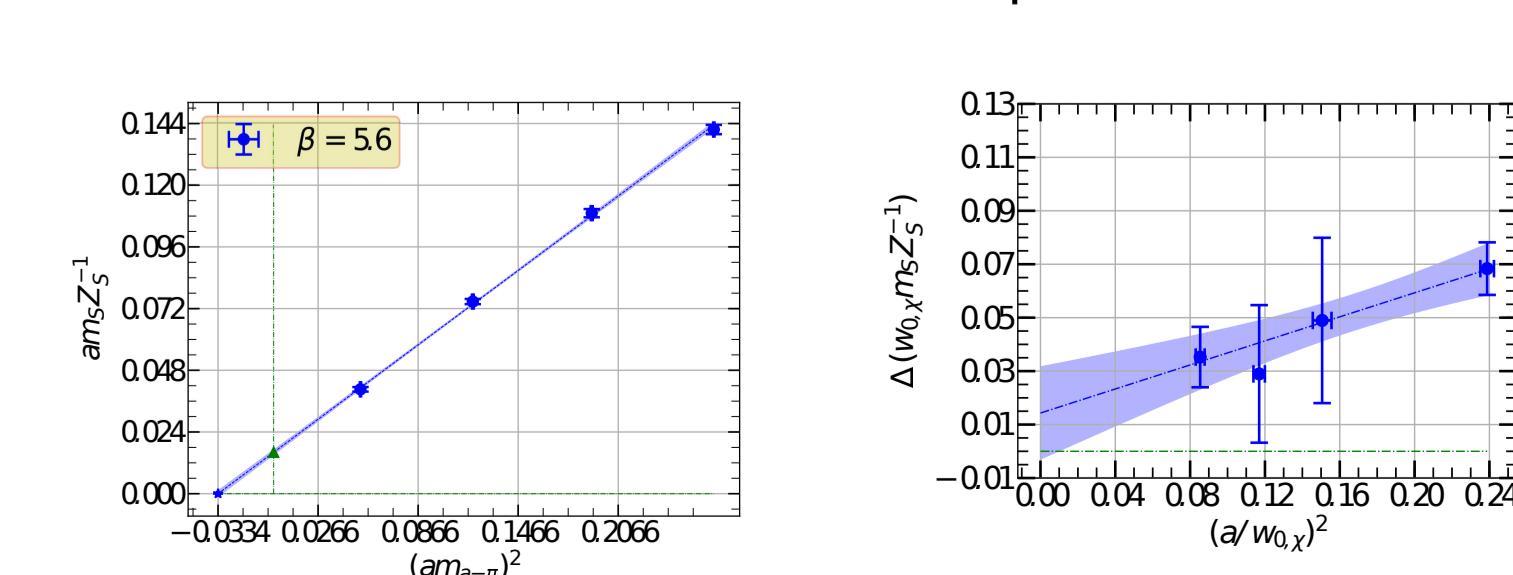
- through spectral flow
- by calculation of real negative eigenvalues of Q with Arnoldi
- Negative Pfaffians occur in our simulations near κ_c , but rarely.

SUSY Ward identities on lattice

Ward identity is the quantum version of Noether Theorem
Renormalised SUSY WI on the lattice with non-zero gluino mass

$$\langle (\nabla_\mu S_\mu(x)) Q(y) \rangle + Z_T Z_S^{-1} \langle (\nabla_\mu T_\mu(x)) Q(y) \rangle \\ = m_S Z_S^{-1} \langle \chi(x) Q(y) \rangle + \langle X_S(x) Q(y) \rangle - \left\langle \frac{\delta Q}{\delta \varepsilon(x)} \right\rangle.$$

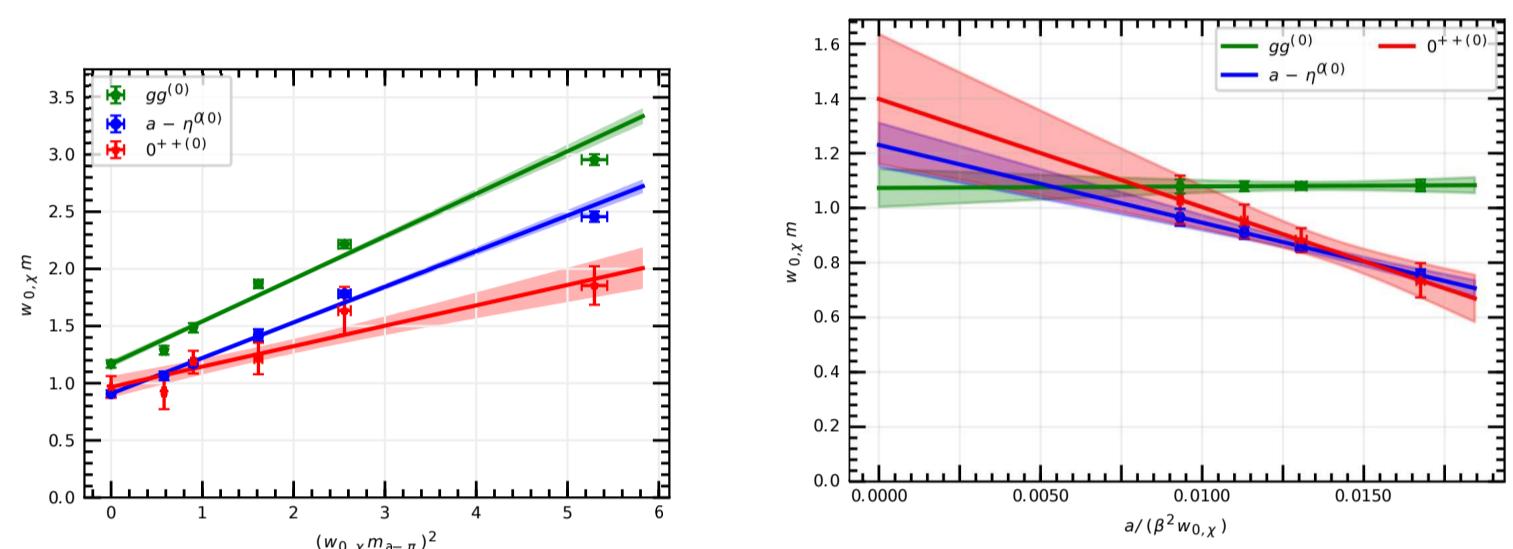
- Z_S and Z_T are renormalization coefficients
- $S_\mu(x)$ and $T_\mu(x)$ are the Super and the Mixing currents.
- Q is an insertion operator.
- $\varepsilon(x)$ is the parameter of infinitesimal symmetry transformations.
- $\left\langle \frac{\delta Q}{\delta \varepsilon(x)} \right\rangle$ is contact term, which is zero if Q is localised at space-time points different from x .
- $\langle X_S(x) Q(y) \rangle$ is introduced by the lattice regulator and vanishes in the continuum limit
- $\langle m_S Z_S^{-1} \langle \chi(x) Q(y) \rangle \rangle$ is the mass term which break SUSY softly
- $\langle a m_S Z_S^{-1} \rangle$ is determine for each gauge ensemble
- Remnant gluino mass $\Delta(w_{0,\chi} m_S Z_S^{-1})$ in physical units is determined in the chiral limit
- The continuum extrapolation is performed, $\Delta(w_{0,\chi} m_S Z_S^{-1})$ is consistant with zero within error as expected from the theory



Results

Light particle spectrum

- Mass of light bound states is determined
- Variational analysis is used
- Mixing is considered
- Chiral extrapolations are performed
- Mass gap between different states due to soft breaking of SUSY is observed
- Mass degenerate chiral supermultiplet is formed as expected from predictions
- Gauge group SU(3) is considered

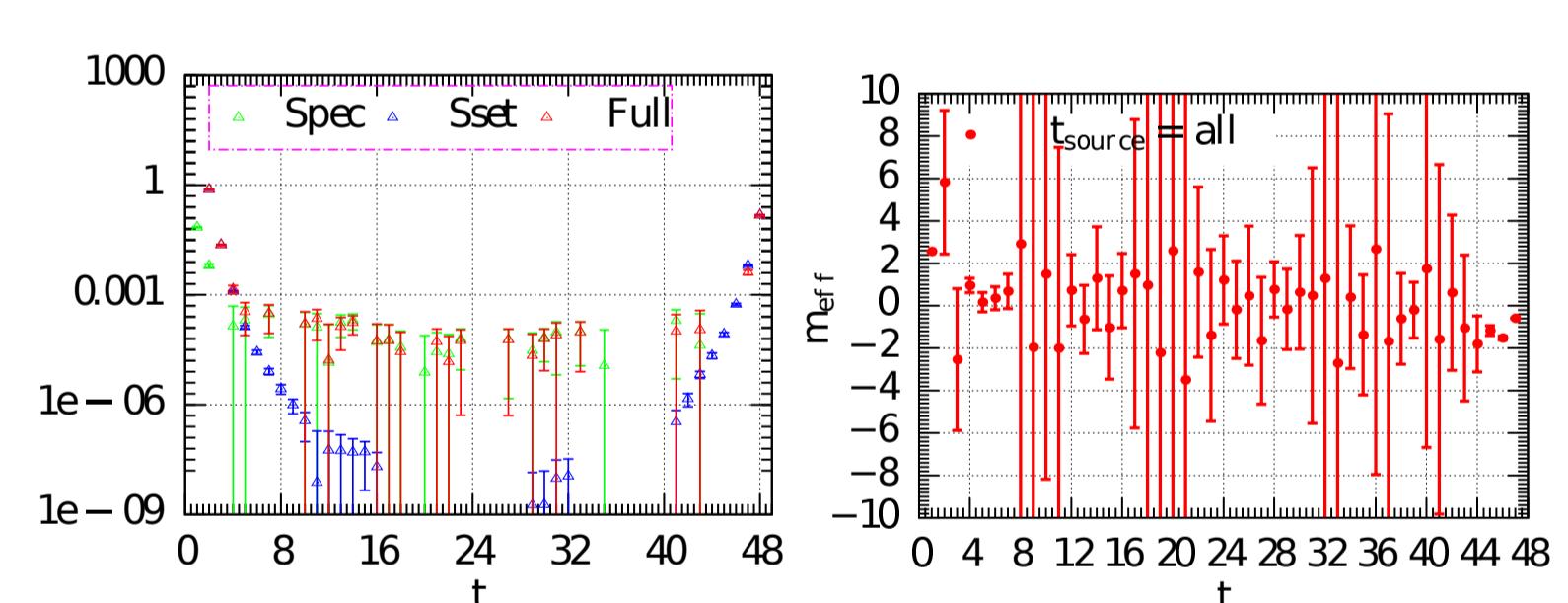


Baryonic States

Baryonic correlation functions can be constructed from Rarita Schwinger field

$$W_\mu(x) = t_{abc} \lambda_a(x) \left(\lambda_b^T(x) C \gamma_\mu \lambda_c(x) \right).$$

- Baryonic states are not predicted by the effective actions
- Constructed from three gluino fields in analogy with QCD
- Consists of two parts; the sunset piece and the spectacle piece
- Stochastic estimator technique is used for spectacle piece
- The lowest eigenvalues for the inverse of the Wilson-Dirac operator is used
- Computation of sunset piece is easy whereas the spectacle piece is rather challenging
- Preliminary results of correlation functions and the effective mass



Conclusion and Outlooks

- First order phase transition at $m_{\tilde{g}} = 0$
- Determination of the mass spectrum of light bound states
- Chiral supermultiplet is formed
- Baryonic states are investigated
- SUSY WIs are being analysed
- SUSY is restored
- More states in SYM theory can be formulated

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