# **Cosmological Scattering Potentials and** their Quantum Simulation



# theory

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## experiment







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# Cosmological Particle Production

**Quantum Simulation** 

#### of QFT

#### in Curved Space





# **Analogue Quantum Simulation of Kinematics in Curved Spacetime**

Horizons Hawking Radiation	<b>Black Hole Spacetimes</b>
Class. Fluids	
Weinfurtner et al. (2011)	
<b>Bose-Einstein-Condensate (BEC)</b>	
Lahav et al. (2010)   Steinhauer (2016)	de Nova et al. (2019)
<b>Optical Fibre</b>	
Philbin et al. (2008)   Choudhary et al.	(2012)   Jaquet (2018)
Microcavity Polaritons	
Nguyen et al. (2015)   Jacquet et.al (202	(2) New avenues for fermionic field
Superfluid 3-He Človečko et al. (2019)	Simeón et.al (2023)
	Haller, Meng et.al (2023)

## BEC

Hung et al. (2012) | Eckel et al. (2018) | Tajik et al. (2023) 2+1 dim: Viermann et al. (2022), Simeón et al. (2022)

#### **False Vacuum Decay**

BEC

Berti et al.(2023), Cominoti et al. (2023) | Jenkins et.al(2024)

Rotating Spacetimes Black Hole Superradiance acetimes **Class.** Fluids Torres et al. (2017) | Cromb et.al (2020) **Photon Superfluid** 9) Vocke et al. (2017) **Microcavity Polaritons** (8)Falque et al. (2023) | Delhom et al. (2023) **Superfluid 4-He** s for fermionic fields: Svancara et al. (2023) ón et.al (2023)

## **Cosmological Particle Production**

**Trapped Ions** Wittemer et al. (2018)

Laser Pulse Steinhauer et al. (2022)

### **Dynamical Casimir effect**

BEC Jaskula et.al (2011) **Superconducting Circuit** Wilson et.al (2011)











# Quantum Simulation of Scalar Field in Curved Spacetime

Effective action for weakly interacting BECs

$$\Gamma[\Phi] = \int \mathrm{d}t \,\mathrm{d}^D x \left\{ \hbar \Phi^*(t, \mathbf{x}) \left[ \mathrm{i} \frac{\partial}{\partial t} - V(t, \mathbf{x}) \right] \Phi(t, \mathbf{x}) - \frac{\hbar^2}{2m} \nabla \Phi^*(t, \mathbf{x}) \nabla \Phi(t, \mathbf{x}) - \frac{\lambda(t)}{2} \left[ \Phi^*(t, \mathbf{x}) \Phi(t, \mathbf{x}) \right]^2 \right\}$$

1. Expand to 2nd order in **fluctuations** 

$$\Phi(t, \mathbf{x}) = e^{iS_0(t, \mathbf{x})} \left( \sqrt{n_0(t, \mathbf{x})} + \frac{1}{\sqrt{2}} [\phi_1(t, \mathbf{x}) + i\phi_2(t, \mathbf{x})] \right) \quad \text{[Simeon et al.(2022)]}$$

2. Evaluate background on Gross-Pitaevskii equation (in hydrodynamic form)

$$\partial_t n_0 + 
abla (n_0 oldsymbol{v}) = 0$$

3. Employ acoustic approximation





$$\hbar\partial_t S_0 + V + \lambda n_0 + rac{\hbar^2}{2m} igg[ (
abla S_0)^2 - rac{
abla^2 \sqrt{n_0}}{n_0} igg] = 0$$

Neglect quantum pressure

$$q=rac{\hbar^2}{2m}rac{
abla^2\sqrt{n_0}}{n_0}$$

$$r \sqrt{g} g^{\mu\nu} \partial_{\mu} \phi \partial_{\nu} \phi$$

# Acoustic FLRW-Spacetime

#### **BEC as stationary background**

$$oldsymbol{v}(t,oldsymbol{x})=(oldsymbol{\hbar}/m)
abla S_0(t,x)=0$$

$$n_0(\mathbf{x}) = \bar{n}_0 \left[1 + \frac{\kappa}{4}r^2\right]^2$$



#### **FLRW line-element**

# Analogue Cosmological Particle Production

#### **Density-Contrast**

$$\delta_c(t,u,arphi) = \sqrt{rac{n_0(u)}{ar{n}_0^3}} \left[ n(t,u,arphi) - n_0(u) 
ight] ~~ \sim \partial_t \phi + \mathcal{O}(\phi^2)$$

$$\langle \delta_c(t,u,arphi) \delta_c(t,u',arphi') 
angle = rac{oldsymbol{\hbar}^2 m}{\lambda_f^2 ar{n}_0^3} \langle \dot{\phi}(t,u,arphi) \dot{\phi}(t,u',arphi') 
angle$$

Quasi-particle occupancies measurable through full density-contrast correlations

#### **Power Spectrum**

$$\frac{1}{2}\left\langle 0\right|\left\{ \dot{\phi}(t,x),\dot{\phi}(t,x')\right\}\left|0\right\rangle_{c} = \int_{k}\mathcal{F}(k,L)\frac{\sqrt{-h(k)}}{a_{\mathrm{f}}^{3}}S_{k}(t)$$



# (Analogue) Cosmological Particle Production as a Scattering Problem

#### **Mode equation**

$$igg[-rac{\mathrm{d}^2}{\mathrm{d}\eta^2}+V(\eta)igg]\psi_k(\eta)=-h(k)\psi_k(\eta)\qquad \eta=\intrac{\mathrm{d}}{a(\eta)} d\eta$$

$$V(\eta) = -a^2(\eta) \left[m^2 + \xi R(\eta)
ight] + rac{D-1}{2} \Bigg[rac{a''(\eta)}{a(\eta)} - rac{3-3}{2} \Bigg]$$





#### Eigenenergies

$$h(k) = \begin{cases} -k \left[ k + (D-1)\sqrt{|\kappa|} \right] & \text{for} \\ -k^2 & \text{for} \end{cases}$$

$$\left[-\left[k^2 + \left(\frac{D-1}{2}\right)^2 |\kappa|\right]$$
 for

#### **Region III**

$$r_k \cdot e^{i \omega_k \eta}$$

$$e^{-i\omega_k\eta}$$

outgoing

 $\eta$ 

vacuum

#### [CFS, Lopez, Simeón, Flörchinger, **Oberthaler group (in preparation)]**





# (Analogue) Cosmological Particle Production as a Scattering Problem

**Power Spectrum** 

$$S_k(t) = rac{1}{2} + N_k + \Delta N_k^0 \cos(2\omega_k\eta(t) + artheta_k)$$



#### Offset

$$N_k = \left|rac{r_k}{t_k}
ight|^2$$

Amplitude

$$\Delta N_k^0 = igg| rac{r_k}{t_k^2}$$



Phase

$$artheta_k = rg(-r_k \mathrm{e}^{-2\mathrm{i}\omega_k\eta_\mathrm{f}})$$

[CFS, Lopez, Simeón, Flörchinger, **Oberthaler group (in preparation)]** 



# **Quantum Simulation of Cosmological Scattering Potentials**

## Focus on minimally coupled, massless fields in D = 2 spatial dimensions



#### **Discontinuous transitions imply singular** contributions

$$V_{
m s}(\eta) = rac{\dot{a}(t(\eta))}{2} [\delta(\eta-\eta_{
m i})-\delta(\eta-\eta_{
m f})]$$









# **Results:**





# **Results:**





# Periodical Cosmological Scattering Potentials



#### **Transfer matrix method**

# Periodical Cosmological Scattering Potentials



 $\omega_k \eta_b$ 





 $\omega_k \eta_b$ 



#### Invariant Cosmological Vacua through Transparant Potentials **Poeschl-Teller potential** $V(\eta) = -\alpha^2 \lambda (\lambda - 1) \operatorname{sech}^2(\alpha \eta)$ Spectrum



$$a(\eta) = [c_1 P_{\lambda-1}( anhlpha\eta) + c_2 Q_{\lambda-1}( anhlpha\eta)]^{2/(D-1)}$$



**Continuum shift:**  $V(\eta) = \alpha^2 \lambda^2 - \alpha^2 \lambda (\lambda - 1) \operatorname{sech}^2(\alpha \eta)$ 

**Constant eff. mass** (asympt.):

$$\omega_k^2 o \omega_k^2 + (a)$$





# Invariant Cosmological Vacua through Transparent Potentials

Generalized transparent potentials [Kay, Moses (1956)]

$$egin{aligned} V(\eta) &= -2rac{\mathrm{d}^2}{\mathrm{d}\eta^2}\mathrm{log}\,\mathrm{det}[\mathbbm{1}+A(\eta)] \ &= -4\sum_{m=1}^N\kappa_m\psi_m^2(\eta) \end{aligned}$$



Scattering amplitudes +Bound state energies

 $\hat{A}_{nm} = \frac{\sqrt{A_n A_m}}{\kappa_n + \kappa_m} \exp\{(\kappa_n + \kappa_m)\eta\}$ 

Bound-state energies

#### Correspond to Solitons of Korteweg-deVries-hierarchy [e.g. Gardner et al. (1974)]

#### Transparent property related to integrability of inverse scattering transform:

Scattering potential

## Gelfand, Levitan (1955); Marchenko (1955)

# Future avenues

## **Inverse Scattering Theory and Isospectrality:**

1. Infer cosmological evolution from power spectrum

### **Full Mode Dispersion:**

[Corley, Jacobson (1996); Martin, Brandenberger (2002)]

## Analogue rainbow metric

Cosmological Particle Production [e.g. Weinfurtner et al. (2008)] Hawking radiation [e.g. Coutant, Weinfurtner (2017)]

# **Non-linear mode interaction in BEC:**

Dissipation effects [e.g. Micheli, Robertson (2023)]

Quantum entanglement of two-mode-squeezed states

- 2. Isospectral scattering potentials [Cooper,Khare,Sukhatme (1995), Dunne, Feinberg (1998)] **Iso-spectral QFTCS (different coupling to background)**

- Beyond healing-length: Bogoliubov dispersion [Bogoliubov (1946), Volovik (2009)]
  - Corresponds to Corley-Jacobson dispersion as UV-completion Transplanckian Problem

# Primordial Cosmological Perturbations



Source: ESA/Planck Collaboration 2018

Generation of both scalar and tensor perturbations to metric is essentially captured by massless scalar on FLRW-spacetime [Martin,Brandenberger(2001); Mukhanov,Winitzki (2013); Brandenberger, Peter (2016)]

Seeds for cosmic structure formation were generated through dynamic background Universe

e.g Inflation: Miniscule vacuum fluctuations were amplified and stretched beyond cosmological horizon and ceased oscillating [Mukhanov, Feldman, Brandenberger (1992]

Alternatives are possible, e.g. **Bouncing Cosmologies** [Brandenberger,Peter (2016); ljjas,Steinhardt (2018)]







# **BackUp: Primordial Cosmological Perturbations**

#### Scalar and tensor perturbations to the metric obey (Martin (2008))

## $rac{\mathrm{d}^2 \mu_{S,T}(k,\eta)}{\mathrm{d} n^2} + \omega_{S,T}^2(k,\eta) \mu_{S,T}(k,\eta) = 0$ wit

 $\bullet$ 

th 
$$\omega_S^2(k,\eta)=k^2-rac{(a\sqrt{\gamma})''}{a\sqrt{\gamma}}$$
  $\omega_T^2(k,\eta)=k^2-rac{a}{a\sqrt{\gamma}}$ 



approximately constant for a wide class of inflation models (Power-law-inflation)



## **Backup: Scattering Analogy of Cosmological Particle Production**

#### **Generalize theory**

$$\Gamma[\phi] = -\frac{1}{2} \int dt \, d^D x \sqrt{g} \left[ \partial_\mu \phi \partial^\mu \phi + \left( m^2 + \xi R \right) \phi^2 \right]$$
  
Ricci scalar

#### Line-element

$$\mathrm{d}s^2 = -\mathrm{d}t^2 + a^2(t) \left[\frac{\mathrm{d}u^2}{1 - \kappa u^2} + u^2 \mathrm{d}\Omega_D^2\right]$$

1. Conformal time

2. Rescale

 $\mathrm{d}\eta = \frac{\mathrm{d}t}{a(t)}$ 

$$\chi(x) = a^{\frac{D-1}{2}}(\eta)\phi(x)$$

To wit

$$\Gamma[\chi] = -\frac{1}{2} \int \mathrm{d}\eta \, \mathrm{d}^D x \sqrt{\gamma} \, \chi \left[ \frac{\mathrm{d}^2}{\mathrm{d}\eta^2} - \Delta + m_{\mathrm{eff}}^2(\eta) \right] \chi$$

#### Mode expansion

$$\chi(\eta, \boldsymbol{x}) = \int_{\boldsymbol{k}} \left[ a_{\boldsymbol{k}} \mathcal{H}_{\boldsymbol{k}}(\boldsymbol{x}) \psi_{k}(\eta) + a_{\boldsymbol{k}}^{*} \mathcal{H}_{\boldsymbol{k}}^{*}(\boldsymbol{x}) \psi_{k}^{*}(\eta) \right]$$

**Eigenfunctions of Laplace-Beltrami** 

$$\Delta \cdot \mathcal{H}_{k}(\boldsymbol{x}) = -h(k)\mathcal{H}_{k}(\boldsymbol{x})$$

$$h(k) = \begin{cases} -k\left[k + (D-1)\sqrt{|\kappa|}\right] & \text{for } \kappa > 0\\ -k^{2} & \text{for } \kappa = 0\\ -\left[k^{2} + \left(\frac{D-1}{2}\right)^{2}|\kappa|\right] & \text{for } \kappa < 0 \end{cases}$$

#### Mode equation has Schrödinger form

#### **Energy eigenvalue**

$$\left[-\frac{\mathrm{d}^2}{\mathrm{d}\eta^2} + V(\eta)\right]\psi_k(\eta) = E_k\psi_k(\eta)$$

$$E_k=\sqrt{-h(k)}$$

#### Scattering potential

$$egin{aligned} V(\eta) &= -m_{ ext{eff}}^2(\eta) \ &= -a^2(\eta)\left[m^2 + \xi R(\eta)
ight] + rac{D-1}{2} \Bigg[rac{a''(\eta)}{a(\eta)} - rac{3-D}{2}igg(rac{a'(\eta)}{a(\eta)}igg)^2 \Bigg] \end{aligned}$$

#### **Examples:**

$$V(\eta) = -a_0^2 m^2 + \xi D(D-1)\kappa \quad \text{Stationary space} \quad a(t) = a_0$$
$$V(\eta) = -\frac{(D-1)^2}{4}\kappa \quad \text{Conformally coupled, massless} \quad \xi = \frac{D-1}{4D}, \ m = 0$$

Focus on minimally coupled, massless fields from now on  $(\xi = 0, m = 0)$ 

#### **Vanishing potential:** $V(\eta) = 0$

$$\left[\frac{D-1}{2} - q(t)\right]\dot{a}(t)^2 = 0 \qquad \text{Deceleration parameter} \quad q(t) = -\frac{\ddot{a}(t)a(t)}{\dot{a}^2(t)}$$

Implies either stasis ( $\dot{a} = 0$ ) or radiation domination q = (D - 1)/2











$$egin{split} H_0 &= rac{1}{c_{
m s}(a_{
m s}^{
m max})^{q+1}\Delta t} rac{(a_{
m s}^{
m max}/a_{
m s}^{
m min})^{(q+1)/2}-1}{q+1} \ H_0 &= rac{1}{c_{
m s}(a_{
m s}^{
m min})\Delta t} \left[rccos\left(2\sqrt{rac{a_{
m s}^{
m min}}{a_{
m s}^{
m max}}}-1
ight)+2\left(\sqrt{rac{a_{
m s}^{
m min}}{a_{
m s}^{
m max}}}
ight] \end{split}$$





# BackUp: PowerLaw (linear)

#### Spectrum for $\gamma = 1$ , $\alpha_i = 400$ , $\alpha_f = 50$ ; $c_s^{\text{fin}} = 1.1 \,\mu\text{m}/\text{ms}$















 $\Delta t = 3 ms$ 

 $\Delta t = 1.5 ms$ 

















 $\Delta t = 3 ms$ 





BackUp: Power Law (2/3)

#### Spectrum for $\gamma = 2/3$ , $\alpha_i = 400$ , $\alpha_f = 50$ ;

#### $c_s^{\text{fin}} = 1.1 \,\mu\text{m}/\text{ms}$



 $\Delta t = 3 ms$ 

 $\Delta t = 1.5 ms$ 







# **Thermal Initial State**

#### Initial spectrum (linear ramp)







Find further partners with

 $y(\eta) \propto a^{(D-1)/2}(\eta)$ 

Or: Find different potentials with equal scattering coefficients

### minimally coupled, massless

$$a_{\rm lin}(t) = 1 + H_0 t$$
$$D = 3$$

$$a_{\text{quad}}(t) = \left(1 + \frac{H_0}{2}t\right)^2$$
$$D = 2$$

$$y''(\eta) - V(\eta)y(\eta) = 0$$



# Iso-spectral cosmological backgrounds

### Factorize Scattering Hamiltonian (after shift to zero energy solution)

$$H_1 = -rac{\mathrm{d}^2}{\mathrm{d}\eta^2} + V_1(\eta) = A^\dagger A$$

with 
$$A = \frac{\mathrm{d}}{\mathrm{d}\eta} + W(\eta),$$

$$V_1(\eta) = W^2(\eta) - W'(\eta)$$
 is

#### In scattering analogy

$$W(\eta)=-rac{y'(\eta)}{y(\eta)}=-rac{D-1}{2}rac{a'(\eta)}{a(\eta)}$$

$$H_2=-rac{\mathrm{d}^2}{\mathrm{d}\eta^2}+V_2(\eta)=AA^\dagger$$

$$A^{\dagger} = -\frac{\mathrm{d}}{\mathrm{d}\eta} + W(\eta)$$

s iso-spectral to [Dunne, Feinberg (1998)]

$$V_2(\eta) = W^2(\eta) + W'(\eta)$$

with

$$y(\eta) \propto a(\eta)^{(D-1)/2}$$

zero-energy-state

# Iso-spectral cosmological backgrounds

### **Iso-spectrality**

$$V_1(\eta) = rac{D-1}{2} igg[ rac{D-3}{2} igg( rac{a'(\eta)}{a(\eta)} igg)^2 + rac{a''(\eta)}{a(\eta)} igg] ext{ } = rac{D-1}{2} igg[ rac{D-1}{2} igg[ rac{D+1}{2} igg( rac{a'(\eta)}{a(\eta)} igg)^2 - rac{a''(\eta)}{a(\eta)} igg]^2 + rac{a''(\eta)}{a(\eta)} igg]$$

### **Realize** $V_2$ with different QFT on cosmological background

$$egin{aligned} V_2(\eta) &\equiv - ilde{a}^2(\eta)[m^2+ ilde{\xi} ilde{R}(\eta)]+rac{ ilde{D}-1}{2}iggl[rac{ ilde{a}''(\eta)}{a(\eta)}-rac{3- ilde{D}}{2}iggl(rac{ ilde{a}'(\eta)}{ ilde{a}(\eta)}iggr)^2iggr]\ &=rac{1}{2}iggl[rac{D-1}{2}iggl[rac{D+1}{2}iggl(rac{a'(\eta)}{a(\eta)}iggr)^2-rac{a''(\eta)}{a(\eta)}iggr] \end{aligned}$$

