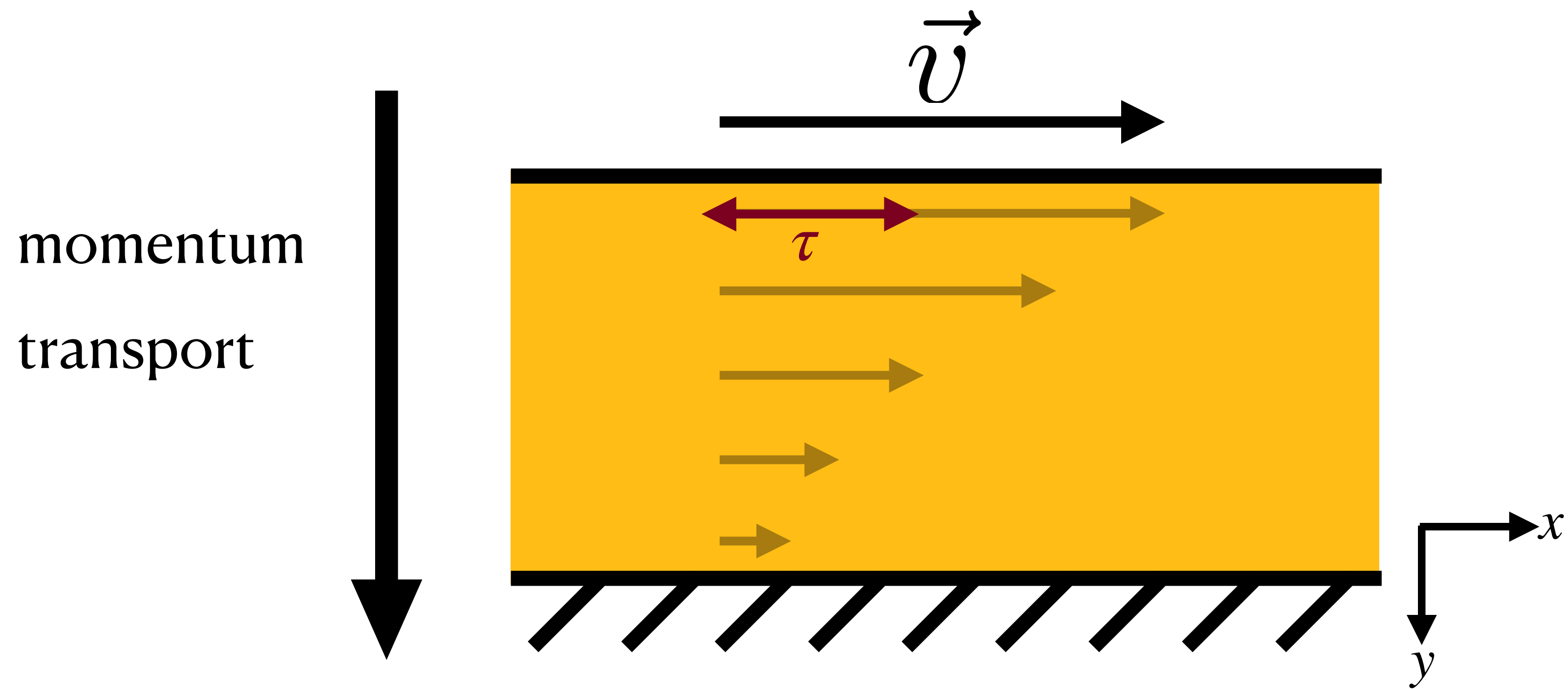


RTG 2522

DFG Research Training Group

From QFT to Hydrodynamic Transport Coefficients

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Newton's law of viscosity

shear stress $\tau = -\eta \frac{dv_x}{dy}$

Motivation

- QGP, neutron stars or dark matter can use hydrodynamical description
- off-equilibrium requires knowledge of macroscopic properties like:
 - equation of state
 - viscosities, conductivities, etc.. $\sigma(T, \mu)$, $\eta(T, \mu)$, $D(T, \mu)$, ...
- interactions determine coefficients

Microscopic Theory

Macroscopic Theory

$$\langle \varphi(x_1) \cdots \varphi(x_N) \rangle \longleftrightarrow \{\sigma, D, \eta, \zeta, \dots\}$$



$$\mathcal{L} = \frac{1}{2} \varphi [\partial_\mu \partial_\mu + m^2] \varphi + \frac{\lambda}{4!} \varphi^4$$

$$\nabla_\mu T^{\mu\nu} = 0, \quad \nabla_\mu J^\mu = 0$$

Linear Response

Classical:

$$\langle X \rangle(t) = \int dt' G_{\text{R}}(t - t') F(t') + \mathcal{O}(F^2)$$

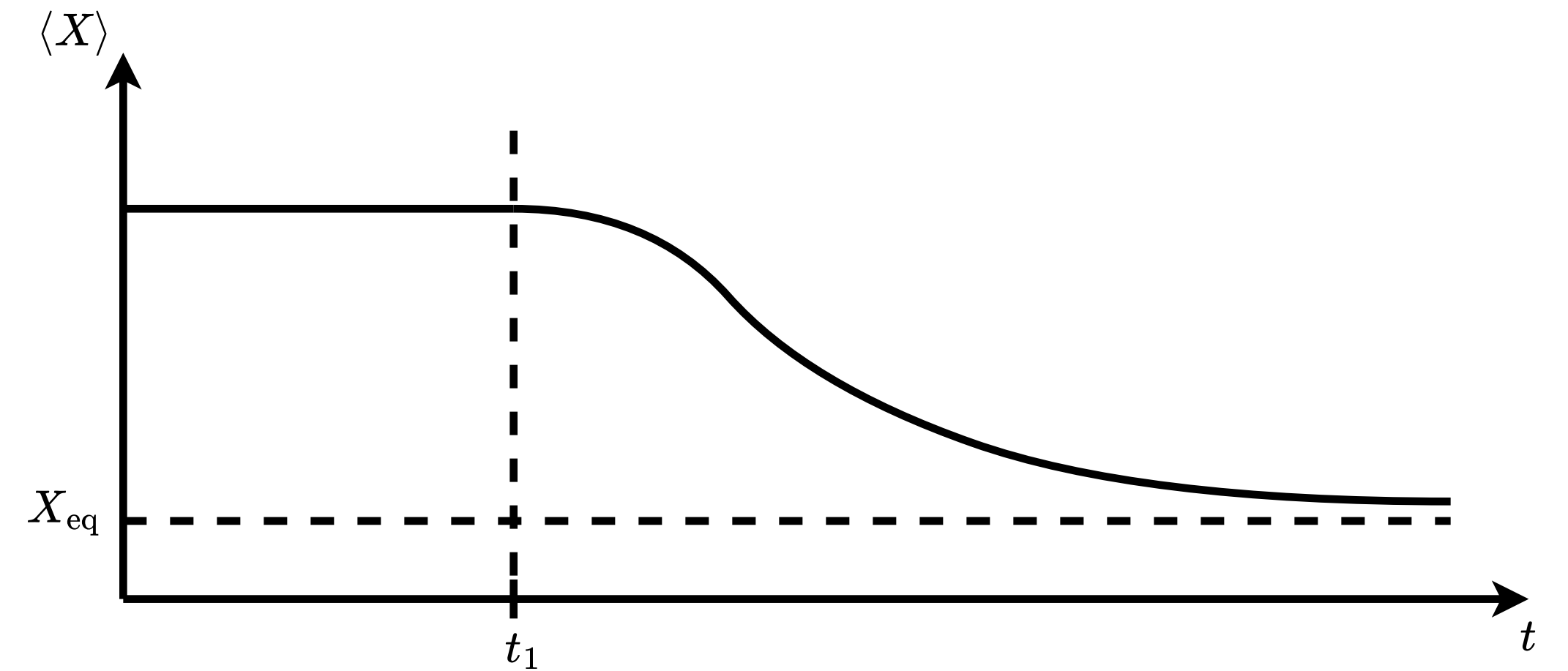
response kernel

QFT:

$$\langle \hat{T}^{\mu\nu} \rangle(x) = \int_y G_{\text{R}}^{\mu\nu\alpha\beta}(x - y) h_{\alpha\beta}(y) + \mathcal{O}(\hbar^2)$$

external force

$$h_{\mu\nu}(x) = g_{\mu\nu}(x) - \eta_{\mu\nu}$$



$$\hat{T}^{\mu\nu}(x) = \frac{1}{2} \frac{\delta S[\varphi; g_{\mu\nu}]}{\sqrt{g} \delta g_{\mu\nu}(x)}$$

Response Kernel:

$$G_{\text{R}}^{\mu\nu\alpha\beta}(x - y) = i\theta(x^0 - y^0) \langle [\hat{T}^{\mu\nu}(x), \hat{T}^{\alpha\beta}(y)] \rangle$$

Hydrodynamics & Transport

- Energy momentum tensor in *Landau frame* (no heat current)

$$T^{\mu\nu} = \epsilon u^\mu u^\nu + (p + \boxed{\Pi}) \Delta^{\mu\nu} + \boxed{\pi^{\mu\nu}}$$

bulk pressure

shear stress

- constituent equations at 1st order

$$\boxed{\Pi = \zeta(T, \mu) \nabla_\mu u^\mu}$$

$$\boxed{\pi^{\mu\nu} = -\eta \Delta^{\mu\alpha} \Delta^{\nu\beta} \left(\nabla_\alpha u_\beta + \nabla_\beta u_\alpha - \frac{1}{2} g_{\alpha\beta} \nabla_\mu u^\mu \right)}$$

- next order: relaxation times $\tau_\Pi, \tau_\pi, \dots$

Hydrodynamics & Transport

- use EoM, constituent equation
- Expand around external source $h_{\mu\nu}$
- e.g. variational approach [Kovtun]:

$$G_{\text{R}}^{\mu\nu\alpha\beta} = -2 \frac{\delta \sqrt{g} T^{\mu\nu}}{g_{\alpha\beta}} \Big|_{g_{\mu\nu} = \eta_{\mu\nu}}$$

Kubo formula

$$G_{\text{R}}^{xyxy}(\omega, \vec{p} = 0) = \text{const} - i\omega\eta(T, \mu) + \mathcal{O}(\omega^2)$$

QFT

- temperature dependence: Imaginary time formalism

$$\int_{-\infty}^{\infty} dt \rightarrow \int_0^{\beta} d\tau, \quad \int_{-\infty}^{\infty} d\omega f(\omega) \rightarrow T \sum_{\omega_n} f(i\omega_n)$$

- retarded Greensfunction:

$$G_R^{\mu\nu\alpha\beta}(\omega, \vec{p}) \sim \frac{1}{2} \frac{\delta^2 \log Z}{\delta^2 g_{\mu\nu} \delta g_{\alpha\beta}} \Big|_{g_{\mu\nu} = \eta_{\mu\nu}}$$

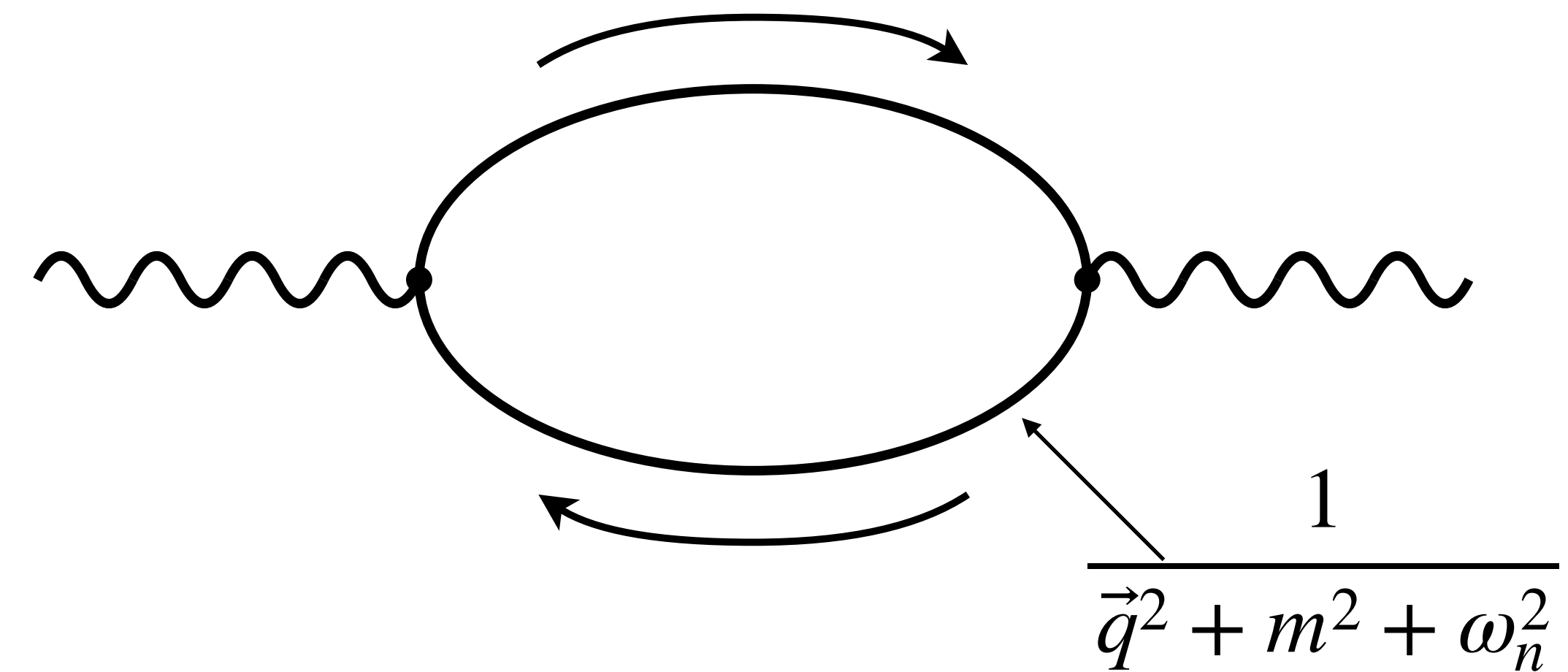
Analytic continuation $i\omega_m \rightarrow \omega + i0^+$

- viscosity in terms of effective action

$$\eta = - \lim_{\omega \rightarrow 0} \partial_{\omega} \text{Im} \frac{\delta^2 \Gamma[\Phi; g_{\mu\nu}]}{\delta g_{\mu\nu} \delta g_{\alpha\beta}} \Big|_{g_{\mu\nu} = \eta_{\mu\nu}}$$

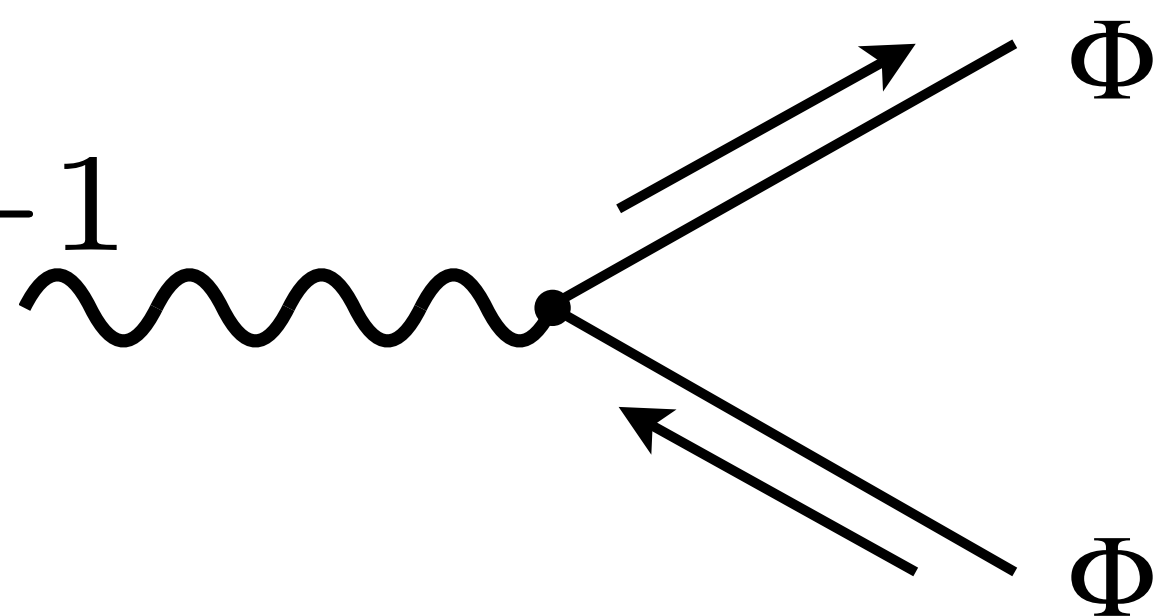
$\langle T^{\mu\nu} T^{\alpha\beta} \rangle$

Viscosity Diagram

$$\eta \sim \lim_{\omega \rightarrow 0} \omega^{-1} \text{Im} \left[\text{Diagram} \right]$$


The diagram shows a loop of two fermions (represented by straight lines with arrows) connected to two external wavy lines. The loop is connected to the wavy lines at two vertices. The propagator is labeled with the expression $\frac{1}{\vec{q}^2 + m^2 + \omega_n^2}$.

Using Cutting rules:

$$\eta \sim \lim_{\omega \rightarrow 0} \omega^{-1} \left[\text{Diagram} \right]$$


The diagram shows a wavy line entering a vertex from the left, which then splits into two outgoing straight lines labeled with the Greek letter Φ .

FRG Flow Equation

- Introduce regulator and RG-scale

$$S[\varphi] \rightarrow S[\varphi] + \int \frac{1}{2} \varphi R_k \varphi$$

- Exact evolution equation [Wetterich '93]

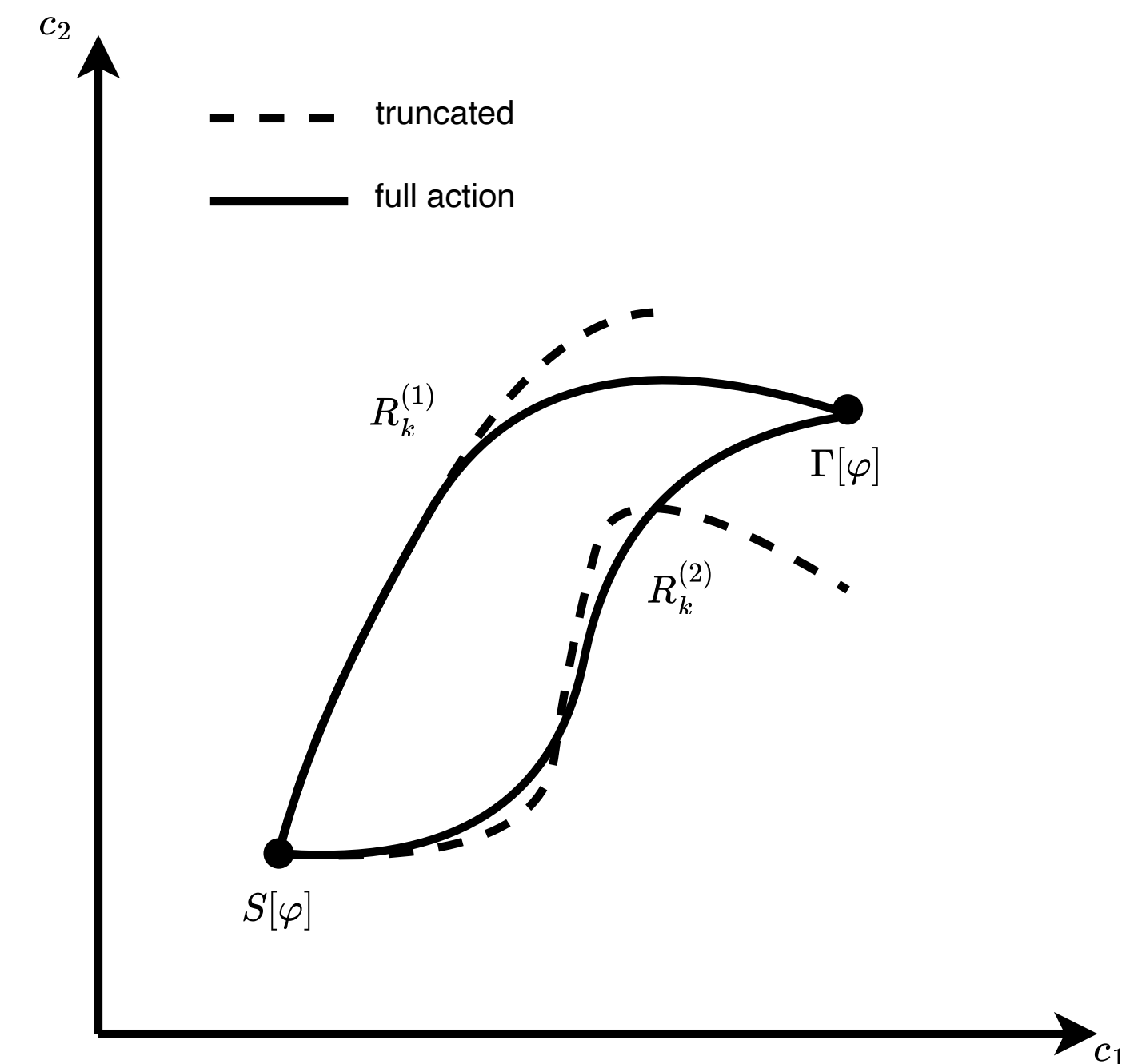
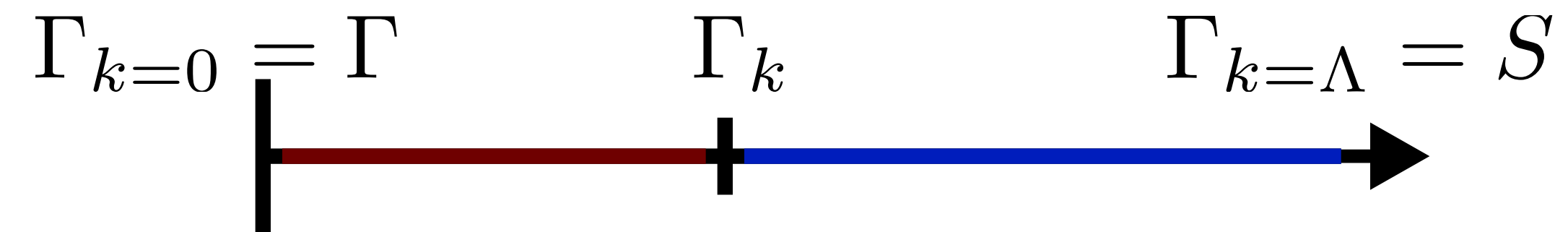
$$k \partial_k \Gamma_k[\Phi; g_{\mu\nu}] = \frac{1}{2} \text{tr} \left[\frac{k \partial_k R_k}{\Gamma_k^{(2)}[\Phi] + R_k} \right]$$

- Truncation needed, e.g.:

$$\Gamma_k[\Phi] \approx \int_x \frac{1}{2} \Phi [\partial_\mu \partial^\mu + m_k^2] \Phi + \frac{\lambda_k}{4!} \Phi^4$$

- with some UV-condition $m_\Lambda^2, \lambda_\Lambda^2, \dots$

- Kubo formula: $\eta(T) \rightarrow \eta_k(T), \quad \eta_{\text{UV}} = 0$



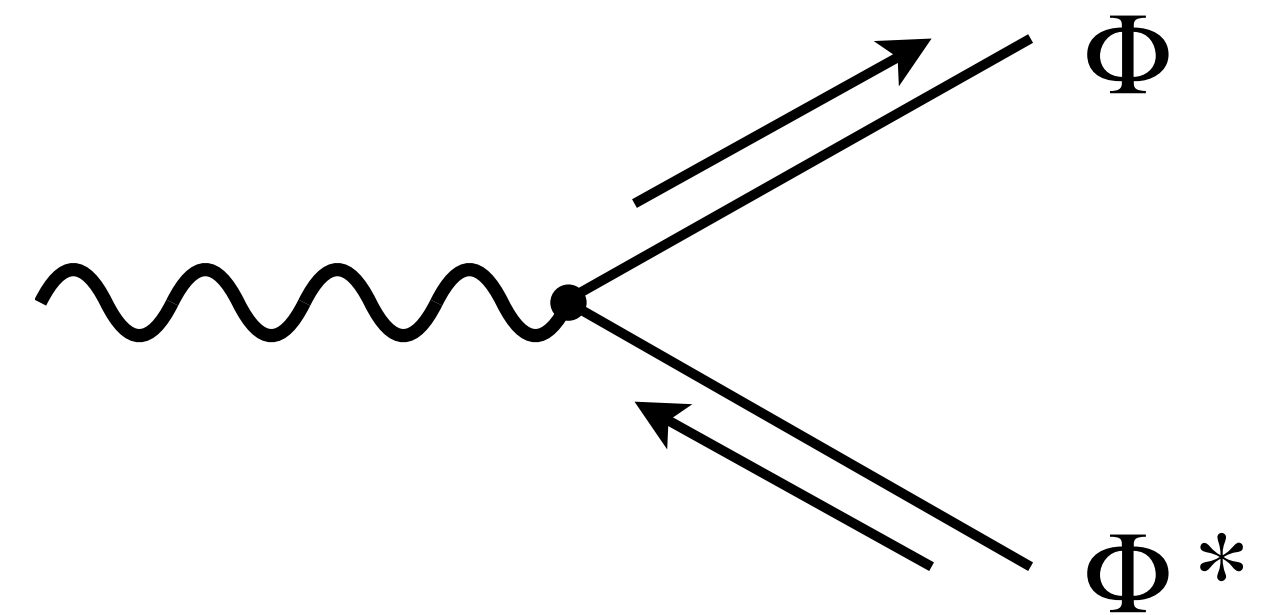
Truncation

- Introduce external gauge field [Rosé, Dupuis '16]

- interactions, thermal mass $\Gamma_k[\Phi; g_{\mu\nu}] = \int d^4x \sqrt{g} \left[\frac{1}{2} \phi (\nabla_\mu \nabla^\mu + m_k^2) \phi + \frac{\lambda}{4!} \right]$

- Damping coefficient [Floerchinger '16]:

$$G(z, \vec{p}) = \frac{1}{m_k^2 + \vec{p}^2 - z^2 - i\gamma_k \text{sgn}(\text{Im}(z))}$$



- In $d=1+0$ dimensions, forward time evolution:

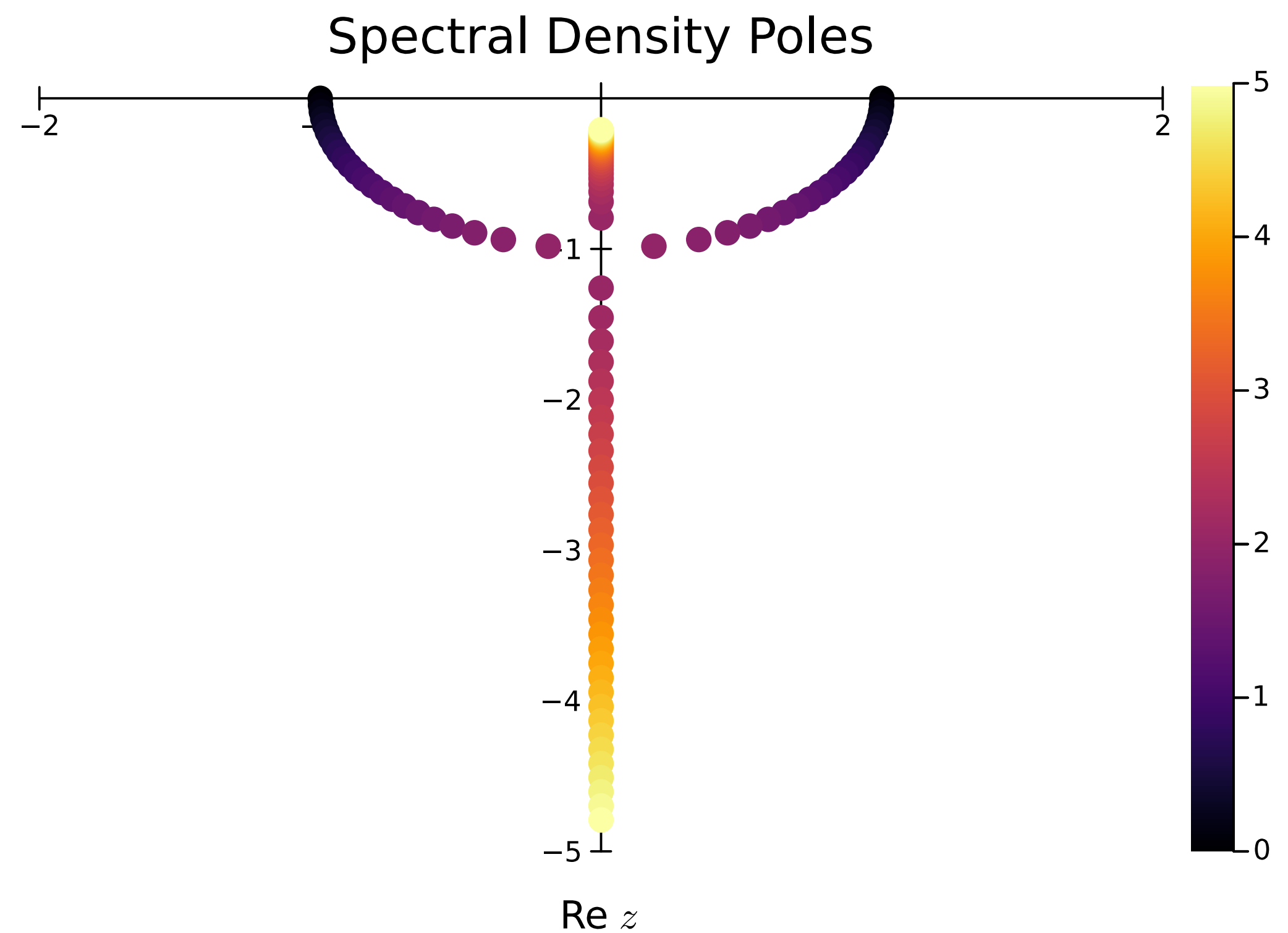
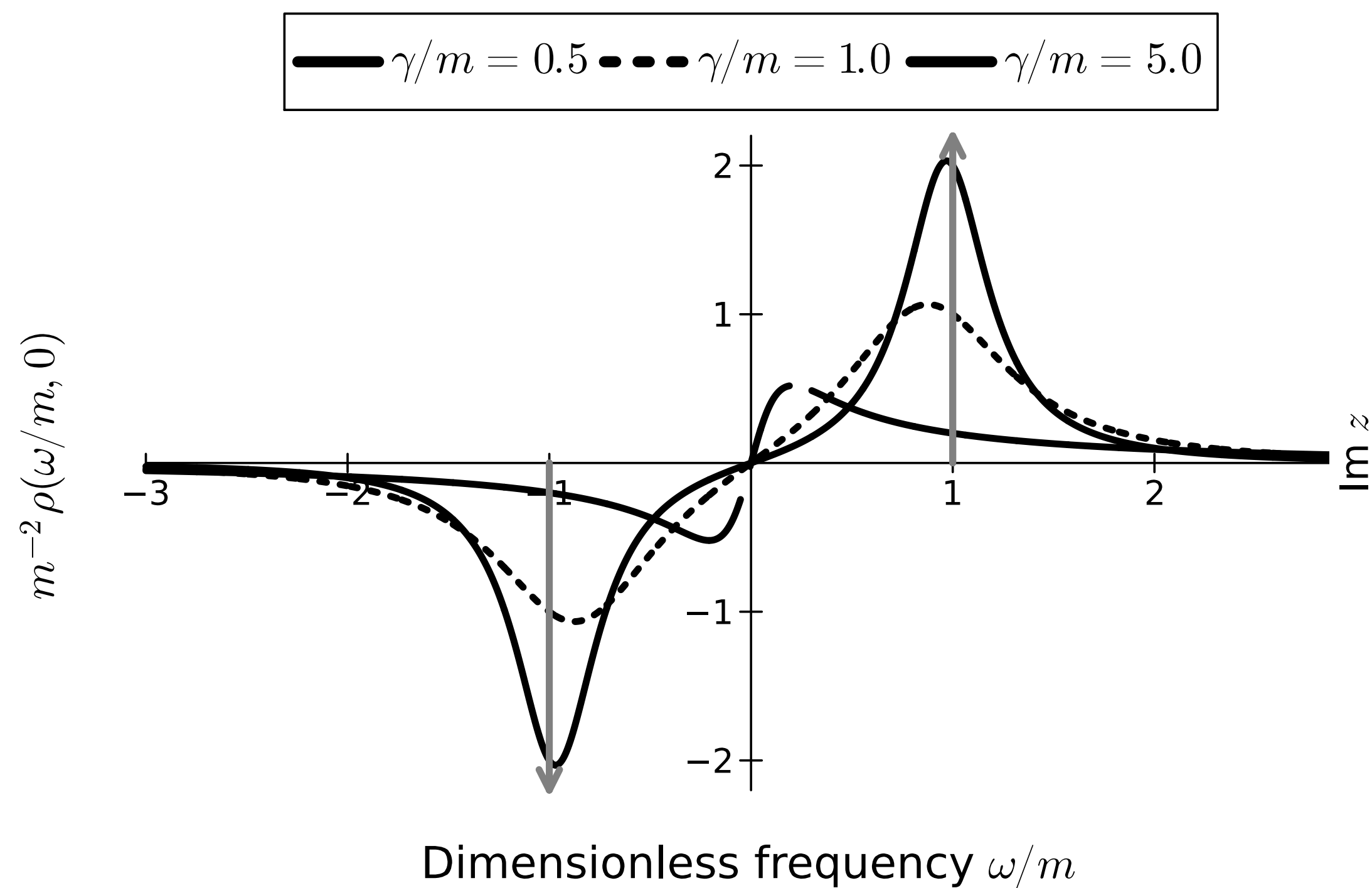
$$[\omega^2 + i\gamma_k \omega - m_k^2] \tilde{\Phi}(\omega) = 0$$

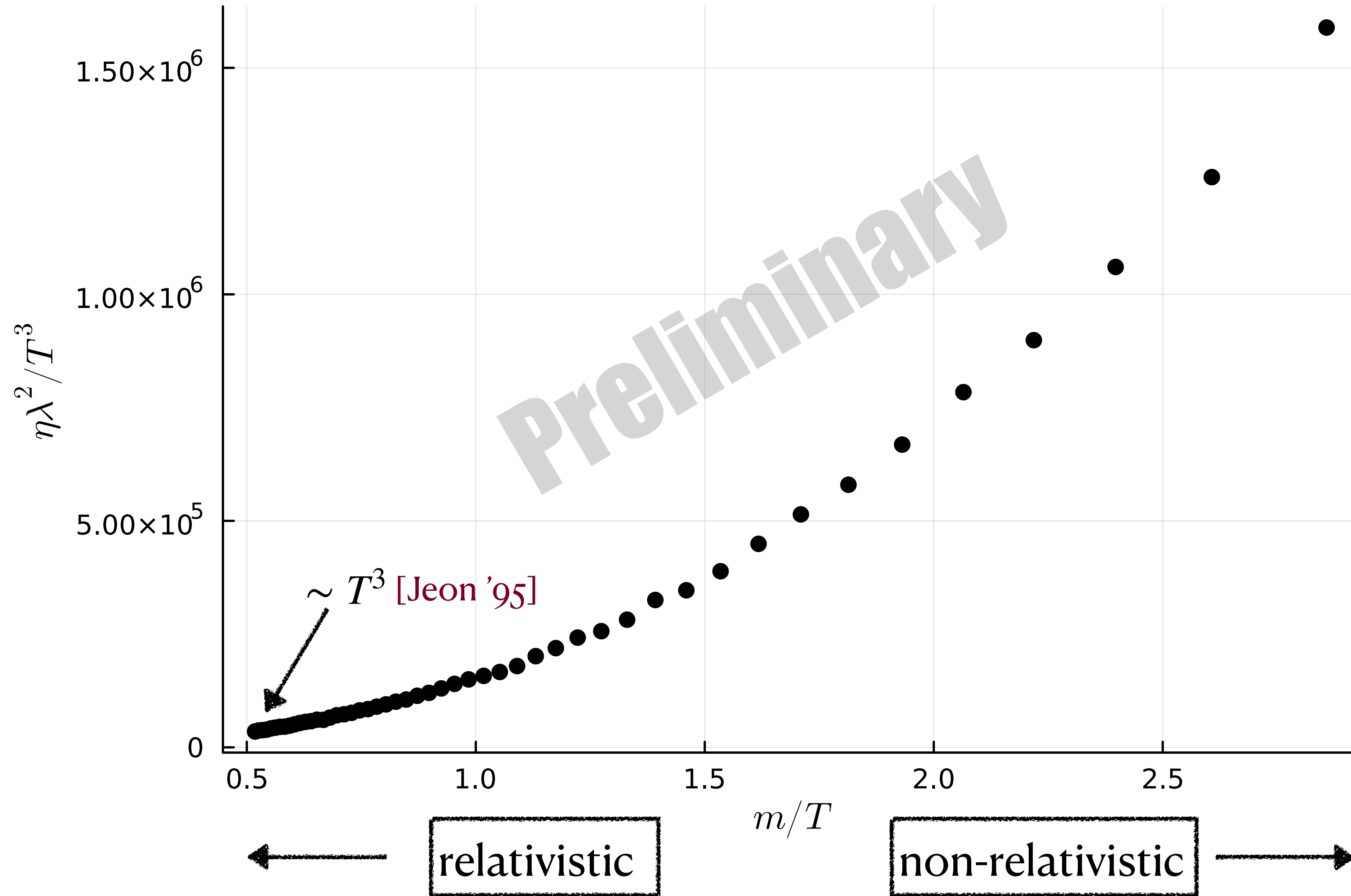
dissipation

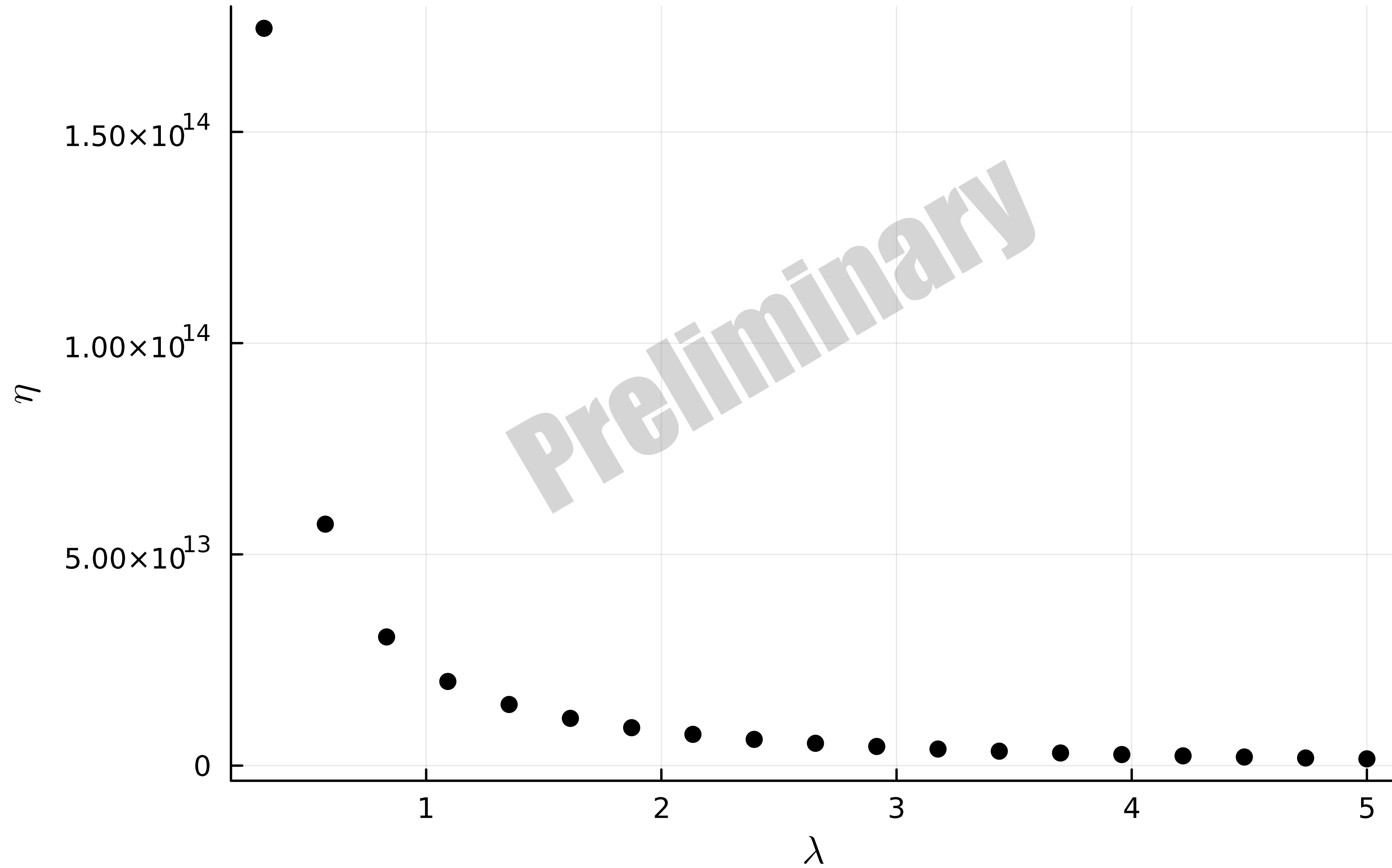
Analytic Structure

- mass pole shifted, decay width generated

$$\rho(z, \vec{p}) = -\frac{1}{2} \text{Im} G(\omega + i0^0, \vec{p})$$







Hydro:
 $\eta \sim l_{\text{mfp}}$

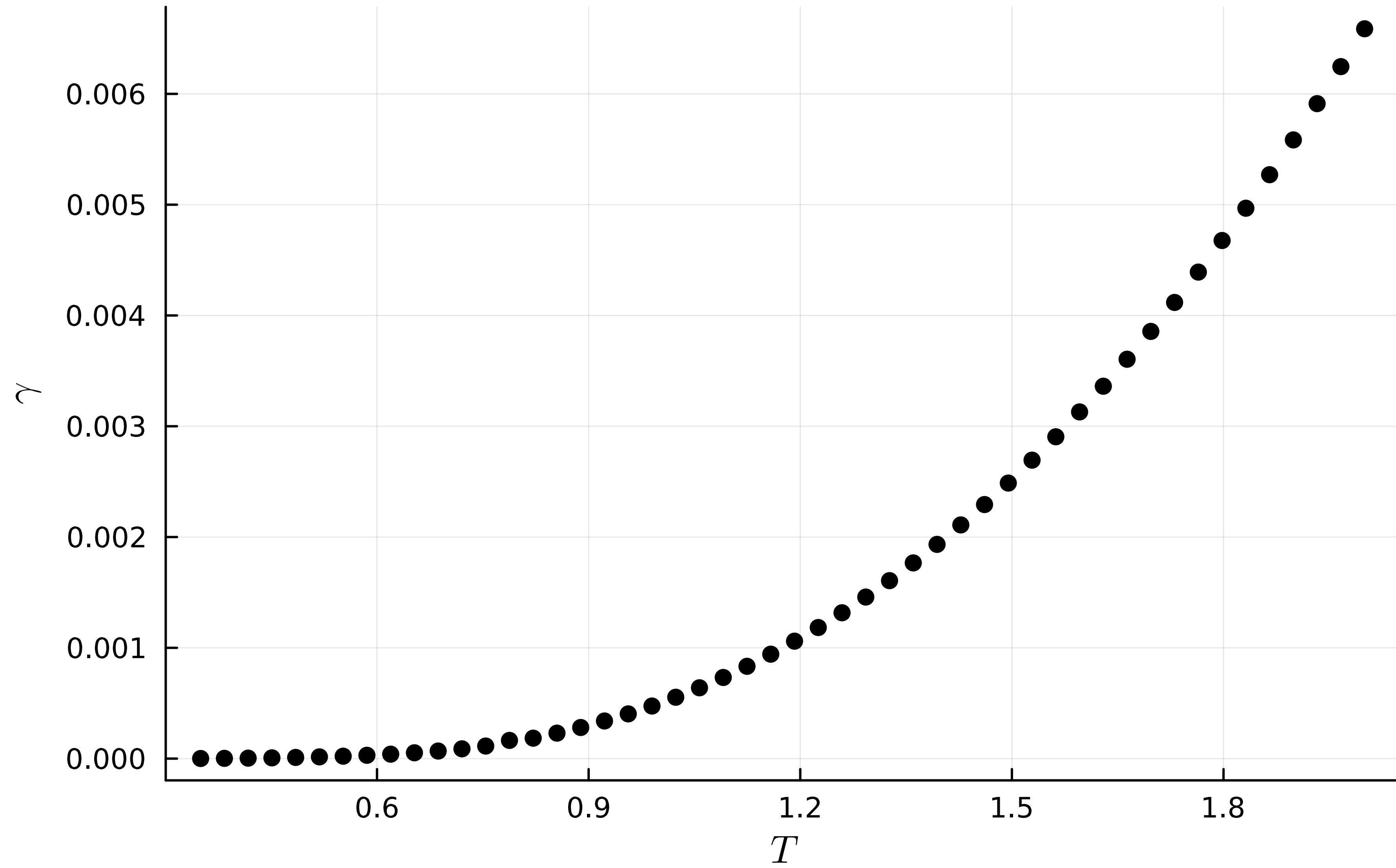
Micro:
 $\frac{1}{\lambda} \sim l_{\text{mfp}}$

Conclusion

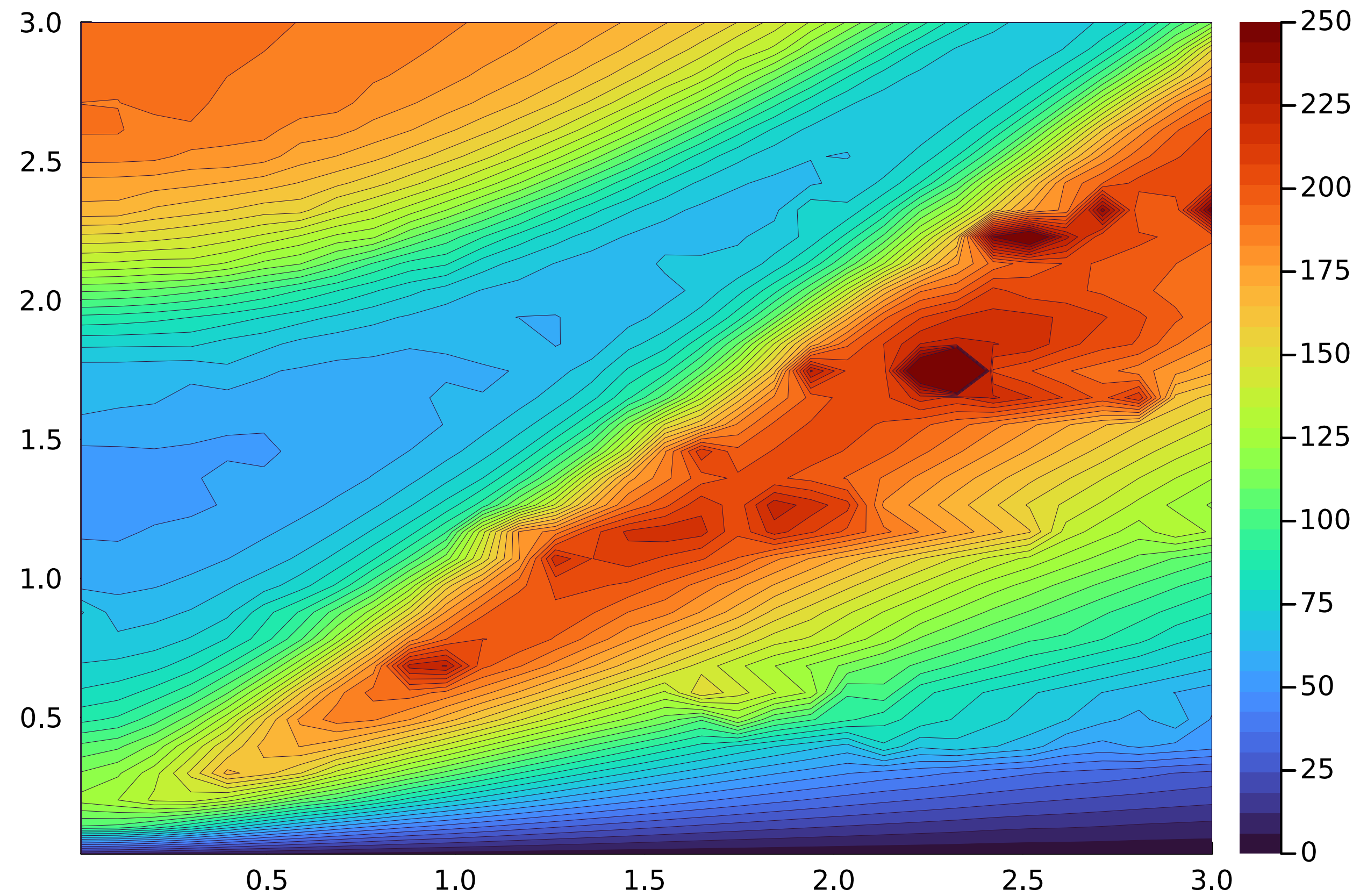
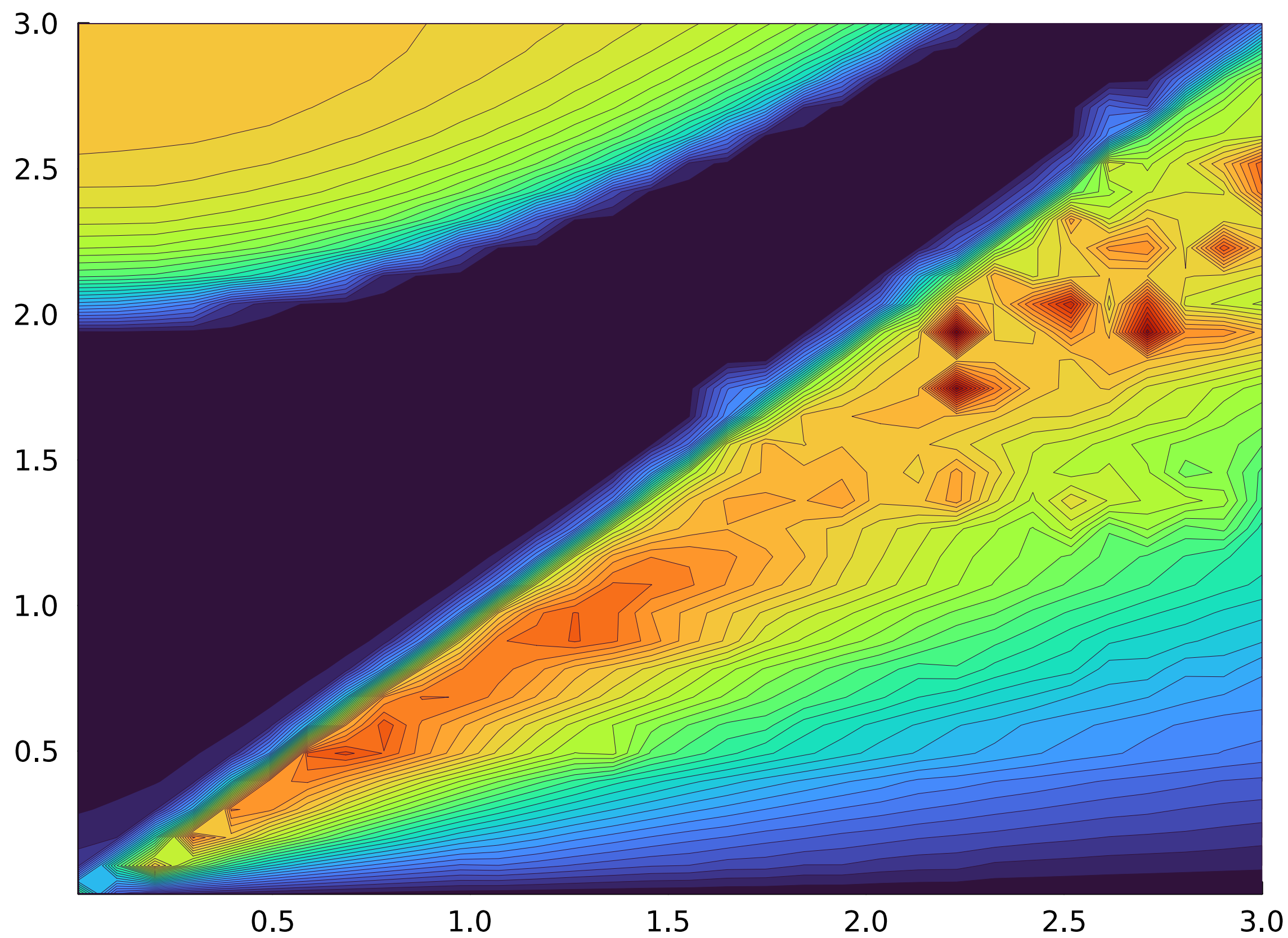
- Prediction of transport coefficients close to equilibrium
- Non-perturbative methods: coefficient flow equation
- Truncation introduces damping coefficients
- Viscosity scaling at small coupling: $\eta(T) \sim T^3$
- diverging at vanishing interaction $\eta \rightarrow \infty$

Thanks for your attention!

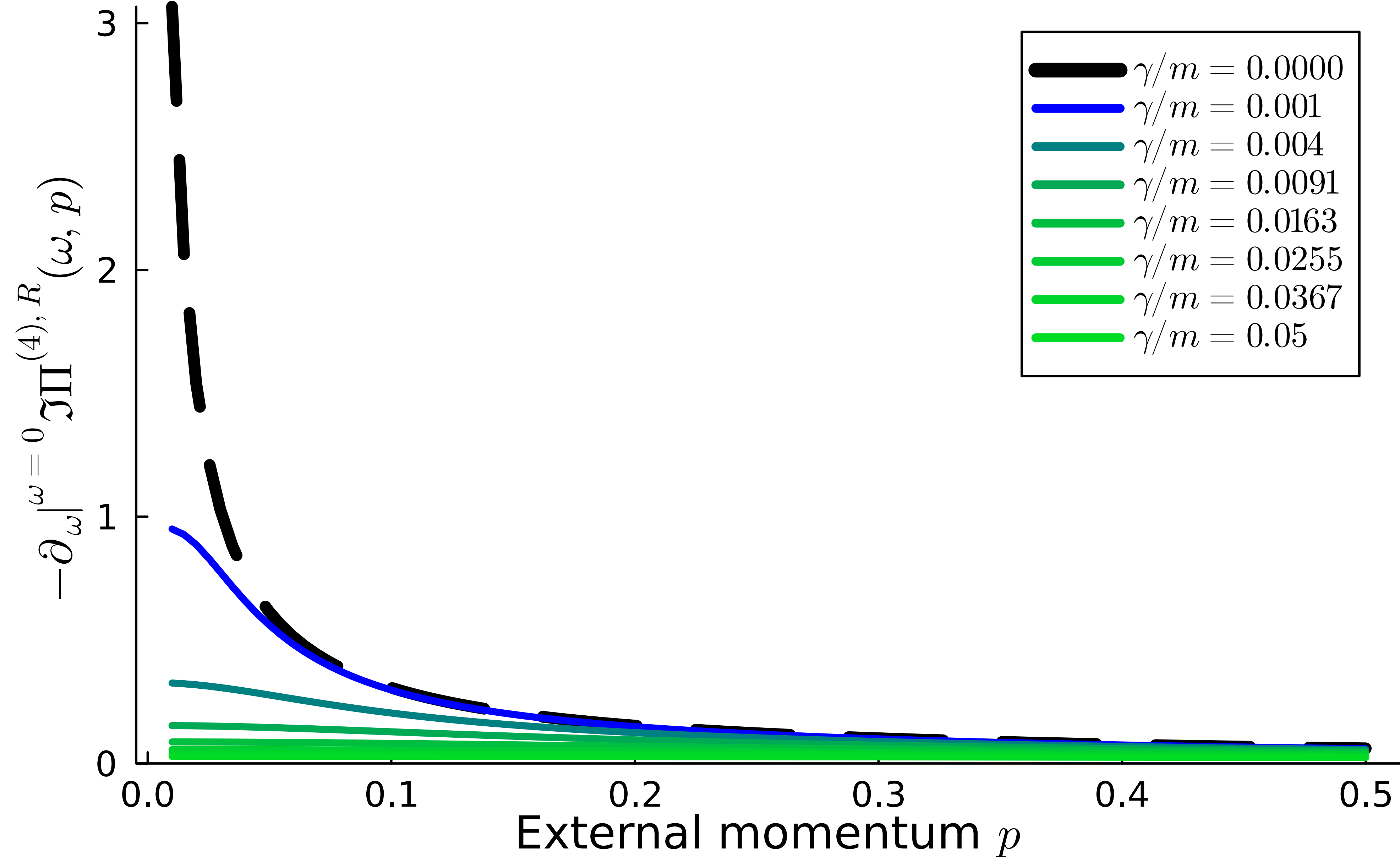
Damping γ



Fourpoint Function - Production Rates



Temperature $m/T = 0.5$



Viscosity Flow Equation

$$\eta = - \lim_{\omega \rightarrow 0} \partial_{\omega} \text{Im} \frac{\delta^2 \Gamma[\Phi; g_{\mu\nu}]}{\delta g_{\mu\nu} \delta g_{\alpha\beta}} \Big|_{g_{\mu\nu} = \eta_{\mu\nu}}$$

Pole Structure of Correlation Functions