

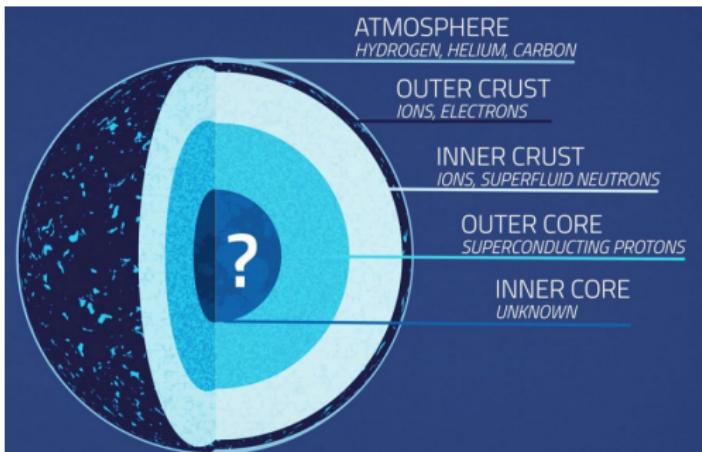
Adibatic Ground States in Non-smooth Spacetimes

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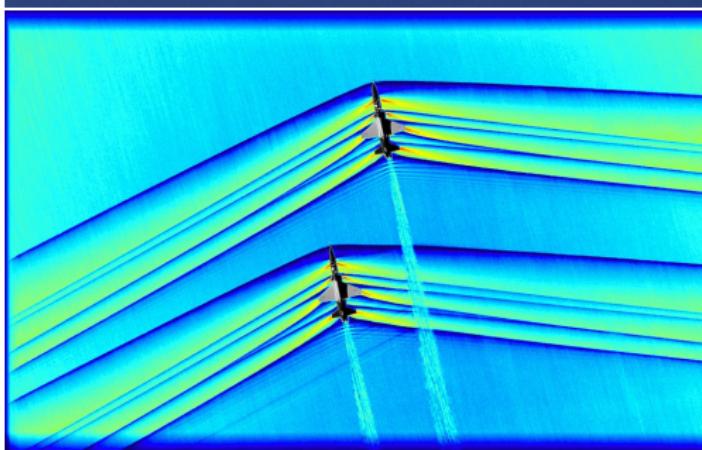
Quantum Effects in Gravitational Fields, Leipzig, 23

Waves Breaking against the Wind, J.M.W. Turner, 1840, Photo © Tate, CC BY-NC-ND

Why **non-smooth** spacetimes?



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THÉORÈME D'EXISTENCE POUR CERTAINS SYSTÈMES
D'ÉQUATIONS AUX DÉRIVÉES PARTIELLES NON
LINÉAIRES.

Par

Y. FOURÈS-BRUHAT.

*Self-Gravitating Relativistic Fluids:
A Two-Phase Model*

DEMETRIOS CHRISTODOULOU

³George M. Bergman, Oberwolfach Photo Collection

⁴George M. Bergman, Oberwolfach Photo Collection

QFT in Non-smooth Spacetimes - general regularity assumptions

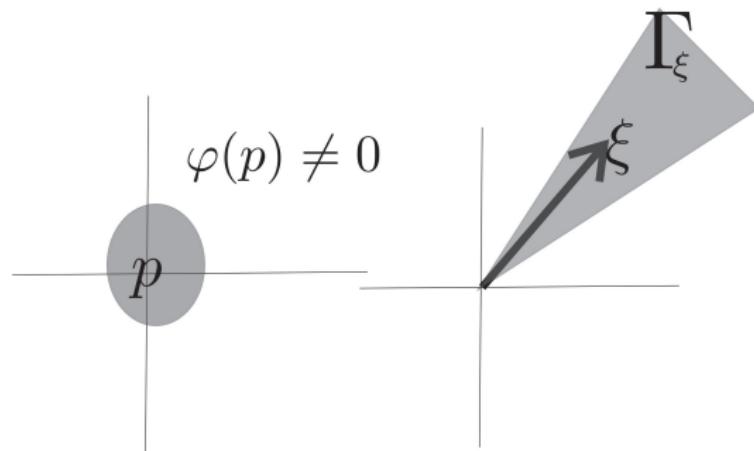
- (Dereziński, Siemssen; 18,19)
 - Constructions of inverses and bisolutions.
 - $L_{loc}^\infty \cap H_{loc}^1$ regularity + local in time assumptions.
- (Vickers,-; 18)
 - Construction of causal propagator based on well-posedness.
 - $C^{1,1}$ regularity.
- (Hörmann, Spreitzer, Vickers,-; 19)
 - Construction of quantization functors.
 - Haag-Kastler Axioms.
 - $C^{1,1}$ regularity.
- (Schrohe,-; 22)
 - Adiabatic ground states.
 - Microlocal structure of causal propagator.
- Future Work:
 - Adiabatic states in non-smooth globally hyperbolic spacetimes.
 - Improvement of adiabatic order.
 - Physical signatures.

Sobolev Wavefront Set and

Non-Smooth Ψ DOS.

\mathbf{H}^s – wavefront set :

$$(p, \xi) \notin WF^s(u) \iff \int_{\Gamma_\xi} (1+|\chi|^2)^s |\widehat{\varphi u}(\chi)|^2 d^n \chi < \infty$$



$$(p, \xi) \notin WF^s(u) \iff u \in H_{mcl}^s(p, \xi)$$

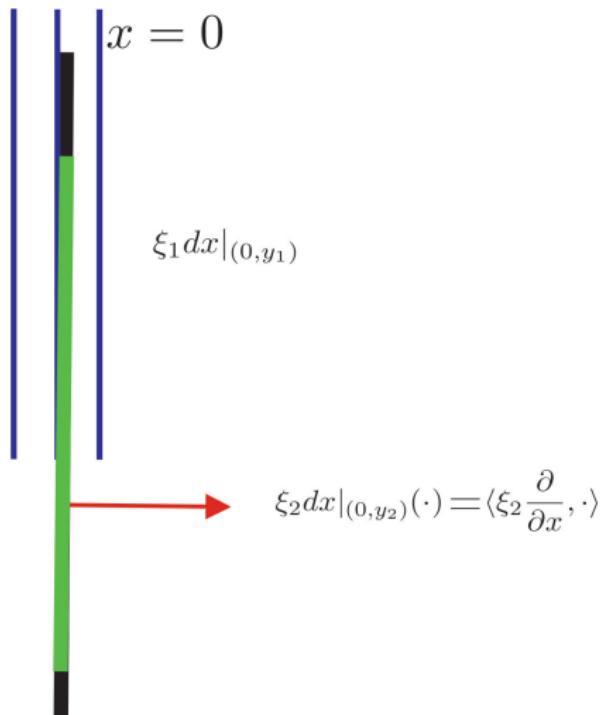
Properties:

- $WF^s(u) \subset T^*(M) \setminus \{0\}$
- $WF^s(u) = \emptyset \iff u \in H_{loc}^s$
- for $s_1 < s_2 \quad WF^{s_1}(u) \subset WF^{s_2}(u) \subset WF(u)$

$$\varphi_{\delta_{x=0}}(\phi) = \int_{\mathbb{R}} \varphi(0, y) \phi(0, y) dy, \varphi \in C_c^\infty$$

$$WF^s(\varphi_{\delta_{x=0}}) =$$

$$\begin{cases} \emptyset & s < -\frac{1}{2} \\ \{(0, y) \in \mathbb{R}^2 \cap \text{supp}(\varphi)\} \times \{(\xi, 0) \in \mathbb{R}^2 \setminus \{0\}\} & s \geq -\frac{1}{2} \end{cases}$$



Non-smooth symbols

A **symbol** $p(x, \xi) \in \textcolor{red}{C}^\tau S_{1,\delta}^{\textcolor{green}{m}}(\mathbb{R}^n \times \mathbb{R}^n)$ if and only if,

$$\|D_\xi^\alpha p(\cdot, \xi)\|_{\textcolor{red}{C}^\tau} \leq C_\alpha \langle \xi \rangle^{\textcolor{green}{m}-|\alpha|+\tau\delta} \text{ for } |\alpha| \geq 0.$$

$$|f|_{C^{0,\beta}} = \sup_{x \neq y} \frac{|f(x) - f(y)|}{\|x - y\|^\beta}, \beta \in (0, 1)$$
$$\|f\|_{C^{\tau=k+\beta}} = \|f\|_{C^k} + \max_{|\rho|=k} |D^\rho f|_{C^{0,\beta}}$$

Non-smooth Ψ DO

The Ψ DO $p(x, D)$ associated to $p(x, \xi) \in C^\tau S_{1,\delta}^m$ is

$$p(x, D)u = \int_{\mathbb{R}^n} e^{ix \cdot \xi} p(x, \xi) \hat{u}(\xi) d\xi$$

Mapping properties:

$$p(x, D) : H^{s+m} \rightarrow H^s$$

for $-\tau(1 - \delta) < s < \tau$.

Symbol decomposition: $p(x, \xi) \in C^\tau S_{1,0}^m$

$$p(x, \xi) = \underbrace{p^\#(x, \xi)}_{\text{smooth}} + \underbrace{p^b(x, \xi)}_{\text{non-smooth, but lower order}}$$

$$p^\#(x, \xi) \in S_{1,\delta}^m$$

$$p^b(x, \xi) \in C^\tau S_{1,\delta}^{m-\tau\delta}; \delta \in (0, 1)$$

Example: The Klein-Gordon Operator

$$p(t, x, \xi_0, \xi) = (-\xi_0^2 + h^{ij}(x)\xi_i\xi_j) + i\frac{1}{\sqrt{h}}\partial_{x^i}(h^{ij}\sqrt{h}(x))\xi_j + m^2$$

$$h^{ij} \in C^\tau, p(t, x, \xi_0, \xi) \in C^{\tau-1}S^2_{1,0}$$

$$p(t, x, D) = \partial_{tt} - \Delta_h + m^2 := P_{KG}$$

Example: The Klein-Gordon Operator

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$$\text{Char}(P_{KG}) := \{(t, x, \xi_0, \xi) \in T^*M \setminus \{0\} : \mathcal{H}(t, x, \xi_0, \xi) = 0\}$$

$$X_{\mathcal{H}} = (\partial_{\xi_0} \mathcal{H}, \partial_\xi \mathcal{H}, -\partial_t \mathcal{H}, -\partial_x \mathcal{H})$$

A single flow line in $\text{Char}(P_{KG})$ is called a **null bicharacteristic strip**.

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A single flow line in $\text{Char}(P_{KG})$ is called a **null bicharacteristic strip**.

The **bicharacteristic relation** C is defined as:

$$C = \{(\tilde{x}, \tilde{\xi}, \tilde{y}, \tilde{\eta}) \in \text{Char}(P_{KG}) \times \text{Char}(P_{KG}), \quad (1)$$

$(\tilde{x}, \tilde{\xi})$ and $(\tilde{y}, \tilde{\eta})$ lie on the same null bicharacteristic strip}

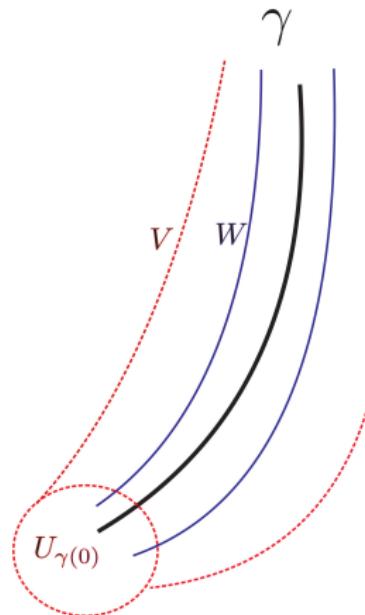
where $\tilde{x} = (t, x), \tilde{\xi} = (\xi_0, \xi), \tilde{y} = (s, y), \tilde{\eta} = (\eta_0, \eta)$.

Propagation of singularities. (Taylor)

Let $P_{KG}(x, \xi) \in C^1 S^2_{1,0}$, γ a null bicharacteristic strip.
If $v \in D'(X)$ solves

$$P_{KG}v = f \iff P_{KG}^\# v = g := f - P_{KG}^b v,$$

$g \in H_{mcl}^\sigma(V)$ and $v \in H_{mcl}^{1+\sigma}(U_{\gamma(0)})$ then $v \in H_{mcl}^{1+\sigma}(W)$.

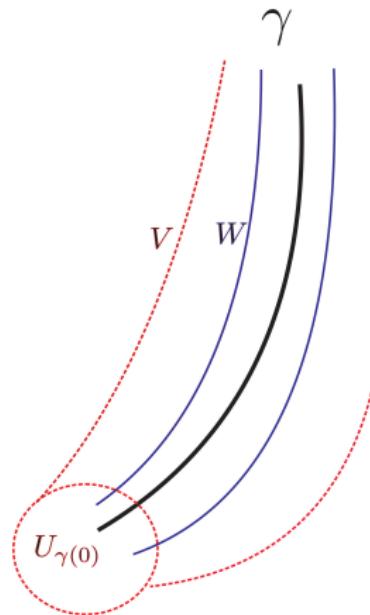


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If $g \in H_{mcl}^\sigma(V)$, $\gamma(0) \in WF^{\sigma+1}(v)$ then $\gamma \in WF^{\sigma+1}(v)$.

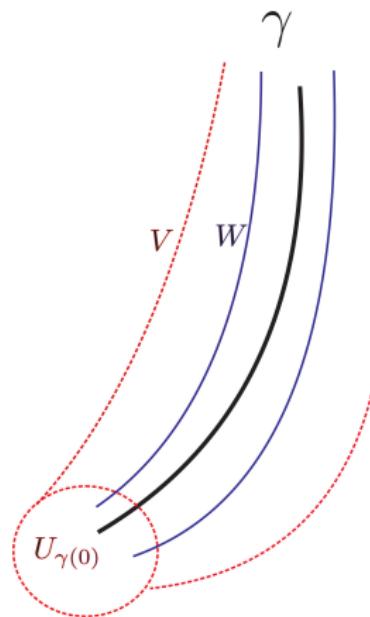


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Adiabatic Ground States

A quasifree state ω_N is an **adiabatic state** of order N if its two-point function Λ_{2N} satisfies for all $s \leq N + \frac{3}{2}$

$$WF'^s(\Lambda_{2N}) \subset C^+$$

$$C^+ = \left\{ (\tilde{x}, \tilde{\xi}, \tilde{y}, \tilde{\eta}) \in C; \tilde{\xi}^0 \geq 0, \tilde{\eta}^0 \geq 0 \right\},$$

where

$$WF'(\Lambda_{2N}) := \{(\tilde{x}, \tilde{\xi}; \tilde{y}, -\tilde{\eta}) \in T^*(M \times M); (\tilde{x}, \tilde{\xi}; \tilde{y}, \tilde{\eta}) \in WF(\Lambda_{2N})\}$$

Ground State

Ultradynamic setting: $M = \mathbb{R} \times \Sigma$, Σ compact

$$ds^2 = dt^2 - h_{ij}(x)dx^i dx^j$$

Klein-Gordon operator: $\partial_{tt} - \Delta_h + m^2$

The **ground state**, ω_g , is completely determined by its two-point function

$$\omega_g^{(2)}(t, x; s, y) = \sum_{l \in \mathbb{N}} \frac{e^{i\lambda_l(t-s)}}{\lambda_l} \phi_l(x)\phi_l(y)$$

The eigenvalues of $-\Delta_h \phi + m^2$ are $\{\lambda_j^2\}_{j \in \mathbb{N}}$ and the set of eigenvectors $\{\phi_l\}$.

Theorem

(Schrohe, - '22) Let (M, g) be a C^τ ultrastatic spacetime with $\tau > 2$, $\dim M = 4$ and $\omega_g^{(2)}$ the two-point function of the ground state. Then

$$WF'^{-\frac{3}{2}+\tau-\tilde{\epsilon}}(\omega_g^{(2)}) \subset C^+ \text{ for every } \tilde{\epsilon} > 0$$

If the spacetime is smooth, then the ground states are adiabatic states of infinite order (Hadamard states).

Lemma

$$\omega_g^{(2)} \in H_{loc}^{-\frac{1}{2}-\epsilon}(M \times M) \text{ for } \epsilon > 0$$

For

$$u(t, s, x, y) = \sum_{j,k} u_{jk}(t, s) \phi_j(x) \phi_k(y) \quad \text{with } u_{jk} = \langle u, \phi_j \otimes \phi_k \rangle$$

$$u \in H^s(\mathbb{R}^2 \times \Sigma^2)$$

$$\iff \sum_{j,k} \int_{\mathbb{R}^2} (\xi_0^2 + \eta_0^2 + \lambda_j^2 + \lambda_k^2)^s |(\mathcal{F} u_{jk})(\xi_0, \eta_0)|^2 d\xi_0 d\eta_0 < \infty \}.$$

$$\omega_g^{(2)}(t, x; s, y) = \sum_{I \in \mathbb{N}} \frac{e^{i\lambda_I(t-s)} \phi_I(x) \phi_I(y)}{\lambda_I}$$

$$\|\omega_g^{(2)}\|_{H_{loc}^{-\frac{1}{2}-\epsilon}(M \times M)}^2 \leq \sum_I \frac{C}{\lambda_I^{3+\frac{2}{3}\epsilon}} \underbrace{\leq}_{\text{Weyl's law}} \sum_I \frac{C'}{I^{1+\epsilon}} < \infty$$

Lemma

For any $\tilde{\epsilon} > 0$

$$WF^{-\frac{1}{2}-\tilde{\epsilon}+\tau}(\omega_G^{(2)}) \subset \text{Char}(P_{KG}) \times \text{Char}(P_{KG}).$$

$$\begin{aligned} P_{(t,x)}(\tilde{x}, \tilde{\xi}, \tilde{y}, \tilde{\eta}) &= (-\xi^0{}^2 + h^{ij}(x)\xi_i\xi_j) + i\frac{1}{\sqrt{h}}\partial_{x^i}(h^{ij}\sqrt{h}(x))\xi_j + m^2 \\ P_{(t,x)}\omega_G^{(2)} = 0 &\iff P_{(t,x)}^\# \omega_G^{(2)} = -P_{(t,x)}^b \omega_G^{(2)} \end{aligned}$$

Lemma

For any $\tilde{\epsilon} > 0$

$$WF^{-\frac{1}{2}-\tilde{\epsilon}+\tau}(\omega_G^{(2)}) \subset \text{Char}(P) \times \text{Char}(P).$$

$$P_{(t,x)}(\tilde{x}, \tilde{\xi}, \tilde{y}, \tilde{\eta}) = (-\xi^0{}^2 + h^{ij}(x)\xi_i\xi_j) + i\frac{1}{\sqrt{h}}\partial_{x^i}(h^{ij}\sqrt{h}(x))\xi_j + m^2$$

$$P_{(t,x)}\omega_G^{(2)} = 0 \iff P_{(t,x)}^\# \omega_G^{(2)} = -P_{(t,x)}^b \omega_G^{(2)}$$

$$P_{(t,x)}^b : \underbrace{H^{s+2-\tau\delta \sim s+2-\tau+\tilde{\epsilon}}}_{H^{-\frac{1}{2}-\epsilon}} \rightarrow H^s$$

$$\delta \sim 1, 0^- < s < \tau, s = -\frac{5}{2} + \tau - \tilde{\epsilon}$$

$$P_{(t,x)}^\# \omega_G^{(2)} = -P_{(t,x)}^b \omega_G^{(2)} =: g \in H^{s=-\frac{5}{2}+\tau-\tilde{\epsilon}}$$

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$$P_{(t,x)}^\# \omega_G^{(2)} = -P_{(t,x)}^b \omega_G^{(2)} =: g \in H^{s=-\frac{5}{2}+\tau-\tilde{\epsilon}}$$

$$WF^{s+2=-\frac{1}{2}+\tau-\tilde{\epsilon}}(\omega_G^{(2)}) \subset \text{Char}(P_{(t,x)}^\#) \cup \underbrace{WF^s(g)}_{\emptyset}$$

Lemma

For any $\tilde{\epsilon} > 0$

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$$WF^{s+2=-\frac{1}{2}+\tau-\tilde{\epsilon}}(\omega_G^{(2)}) \subset \text{Char}(P_{(t,x)}^\#) \cup \underbrace{WF^s(g)}_{\emptyset} = \text{Char}(P_{(t,x)})$$

$$\text{Char}(P_{(t,x)}) = (\text{Char}(P_{KG}) \times T^*M) \cup \{(\tilde{x}, \tilde{y}, \tilde{\xi}, \tilde{\eta}); \tilde{\xi} = 0, \tilde{\eta} \neq 0\}$$

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For any $\tilde{\epsilon} > 0$

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$$WF^{-\frac{1}{2}+\tau-\tilde{\epsilon}}(\omega_G^{(2)}) \subset \text{Char}(P_{(t,x)}) \cap \text{Char}(P_{(s,y)})$$

$$\text{Char}(P_{(t,x)}) \cap \text{Char}(P_{(s,y)}) = \text{Char}(P_{KG}) \times \text{Char}(P_{KG})$$

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$$(\partial_t + \partial_s)\omega_G^{(2)} = 0 \Rightarrow WF^I(\omega_G^{(2)}) \subset \text{Char}(\partial_t + \partial_s) =$$

$$\{(\tilde{x}, \xi_0, \xi, \tilde{y}, \eta_0, \eta) \in T^*(M \times M) \setminus \{0\}; \xi_0 + \eta_0 = 0\}$$

$$\tilde{\xi} = 0 \implies \xi_0 = 0 \implies \eta_0 = 0 \implies \eta_0^2 = h^{ij}(y)\eta_i\eta_j = 0 \implies$$

$$\tilde{\eta} = 0$$

Lemma

(Positivity) For all $s \in \mathbb{R}$,

$$WF^s(\omega_g^{(2)}) \subset \{(\tilde{x}, \tilde{\xi}, \tilde{y}, \tilde{\eta}) \in T^*(M \times M); \tilde{\xi}^0 > 0\}$$

We define $F : \mathbb{R} + i[0, \epsilon_1] \subset \mathbb{C} \rightarrow \mathcal{D}'(\Sigma \times M)$ for $\epsilon_1 > 0$ by

$$F(z) := F(t, \epsilon) = \sum_j e^{i(t+i\epsilon)\lambda_j} e^{-is\lambda_j} \phi_j(x) \phi_j(y)$$

Then, F is holomorphic and $\lim_{\epsilon \rightarrow 0} F = \omega_g^{(2)}$.

Lemma

Let $(\tilde{x}, \tilde{y}) \in M \times M$ be such that \tilde{x} and \tilde{y} are not causally related, i.e. $\tilde{x} \notin J(\tilde{y})$. Then $(\tilde{x}, \tilde{\xi}, \tilde{y}, \tilde{\eta}) \notin WF^{-\frac{1}{2}-\epsilon+\tau}(\omega_g^{(2)})$.

$$WF^{-\frac{1}{2}-\epsilon+\tau}(\omega_G^{(2)}) \subset N_+ \times N_-,$$

$$N_{\pm} := \{(t, x, \xi_0, \xi) \in \text{Char}(P_{KG}); \pm \xi_0 > 0\}$$

$\omega_G^{(2)}|_{\mathcal{Q}}$, \mathcal{Q} the set of pairs of causally separated points.

$\omega_G^{(2)} = \omega^+ + iK_G$. Then $\omega_G^{(2)}|_{\mathcal{Q}} = \omega^+|_{\mathcal{Q}}$.

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"flip" map $\rho(\tilde{x}, \tilde{y}) = (\tilde{y}, \tilde{x})$, $\rho^* \omega_G^{(2)}|_{\mathcal{Q}} = \omega_G^{(2)}|_{\mathcal{Q}}$.

$$WF^{-\frac{1}{2}-\epsilon+\tau}(\omega_G^{(2)}|_{\mathcal{Q}}) = WF^{-\frac{1}{2}-\epsilon+\tau}(\rho^* \omega_G^{(2)}|_{\mathcal{Q}})$$

$$\subset \rho^* WF^{-\frac{1}{2}-\epsilon+\tau}(\omega_G^{(2)}|_{\mathcal{Q}}) \subset \rho^*(N_+ \times N_-) = N_- \times N_+$$

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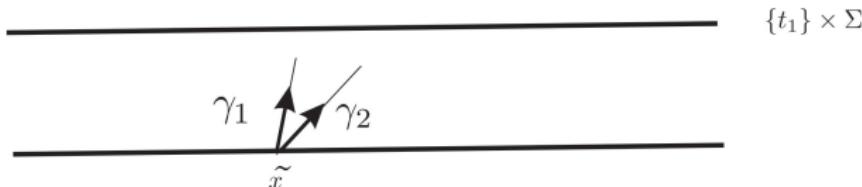
$$WF^{-\frac{1}{2}-\epsilon+\tau}(\omega_G^{(2)}|_{\mathcal{Q}}) \subset (N_+ \times N_-) \cap (N_- \times N_+) = \emptyset.$$

Lemma

If $(\tilde{x}, \tilde{\xi}, \tilde{x}, \tilde{\eta}) \in WF^{-\frac{3}{2}+\tau-\tilde{\epsilon}}(\omega_g^{(2)})$ for $\tilde{\epsilon} > 0$. Then $\tilde{\eta} = -\tilde{\xi}$.

Let $(\tilde{x}, \tilde{\xi}, \tilde{x}, \tilde{\eta}) \in WF^{-\frac{3}{2}-\epsilon+\tau}(\omega_G^{(2)})$ with $\tilde{\eta} \neq \lambda \tilde{\xi}$ then

$(\gamma(\tilde{x}, \tilde{\xi}), \gamma(\tilde{x}, \tilde{\eta})) \in WF^{-\frac{3}{2}-\epsilon+\tau}(\omega_G^{(2)})$ ($P^b \omega_G^{(2)} \in H^{-\frac{5}{2}-\epsilon+\tau=\sigma}$).

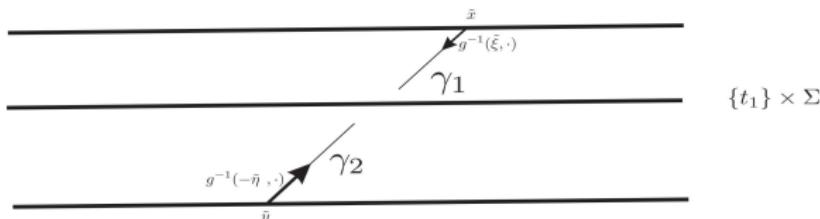


Exists $\{t_1\} \times \Sigma$ where $\Pi \gamma(\tilde{x}, \tilde{\xi}) = (t_1, w_1)$, $\Pi \gamma(\tilde{x}, \tilde{\eta}) = (t_1, w_2)$ are causally separated !

Then $\tilde{\eta} = \lambda \tilde{\xi}$, $\lambda \in \mathbb{R}$.

From $(\partial_t + \partial_s) \omega_G^{(2)} = 0$ we have $\lambda = -1$.

Proof of the Theorem: Let $(\tilde{x}, \tilde{\xi}, \tilde{y}, -\tilde{\eta}) \in WF^{-\frac{3}{2}-\epsilon+\tau}(\omega_G^{(2)})$ then
 $(\gamma(\tilde{x}, \tilde{\xi}), \gamma(\tilde{y}, -\tilde{\eta})) \in WF^{-\frac{3}{2}-\epsilon+\tau}(\omega_G^{(2)})$ and
 $(t_1, w_1, \chi, t_1, w_1, -\chi) \in WF^{-\frac{3}{2}-\epsilon+\tau}(\omega_G^{(2)})$



Define

$$\tilde{\gamma}(\lambda) = \begin{cases} \gamma_1(\lambda) := \Pi(\gamma(\tilde{x}, \tilde{\xi}))(\lambda) & \lambda \in (-\infty, t_1] \\ -\gamma_2(\lambda) := -\Pi(\gamma(\tilde{y}, -\tilde{\eta})(\lambda)) = \Pi(\gamma(\tilde{y}, \tilde{\eta})(\lambda)) & \lambda \in (t_1, \infty) \end{cases}$$

where $-\gamma_2$ denotes the curve with opposite orientation.

Then $\exists a, b \in \mathbb{R}$

$$\tilde{\gamma}(a) = \tilde{x}, \tilde{\gamma}(b) = \tilde{y}; g(\cdot, \dot{\tilde{\gamma}})|_{T_{\tilde{x}}M} = \tilde{\xi}, g(\cdot, \dot{\tilde{\gamma}})|_{T_{\tilde{y}}M} = \tilde{\eta},$$

i.e., $(\tilde{x}, \tilde{\xi}, \tilde{y}, \tilde{\eta}) \in C$.

This gives $WF'^{-\frac{3}{2}-\epsilon+\tau}(\omega_G^{(2)}) \subset C$. Combining with positivity.

$$WF'^{-\frac{3}{2}-\epsilon+\tau}(\omega_G^{(2)}) \subset C^+$$

Let $\omega_A \in \mathcal{D}'(M \times M)$ be the bidistribution given by

$$\omega_A := - \sum_j \lambda_j^{-2} e^{i\lambda_j(t-s)} \phi_j(x) \phi_j(y)$$

Then,

$$\omega_A \in H_{loc}^{\frac{1}{2}-\epsilon}(M \times M) \text{ for every } \epsilon > 0.$$

Properties: $\partial_t \omega_A = i\omega_G^{(2)}$

Idea: Estimates for ω_A and

$$WF^s(\omega_G^{(2)}) = WF^s(\partial_t \omega_A) \subset WF^{s+1}(\omega_A)$$

Let $\{\psi_j; j = 0, 1, \dots\}$ be a Littlewood-Paley partition of unity on \mathbb{R}^n , i.e., a partition of unity $1 = \sum_{j=0}^{\infty} \psi_j$, where $\psi_0 \equiv 1$ for $|\xi| \leq 1$ and $\psi_0 \equiv 0$ for $|\xi| \geq 2$ and $\psi_j(\xi) = \psi_0(2^j \xi) - \psi_0(2^{1-j} \xi)$. The support of ψ_j , $j \geq 1$, then lies in an annulus around the origin of interior radius 2^j and exterior radius 2^{1+j} .

Given $p(x, \xi) \in C^\tau S_{1,0}^m$ and $\gamma \in (0, 1)$ let

$$p^\#(x, \xi) = \sum_{j=0}^{\infty} J_{\epsilon_j} p(x, \xi) \psi_j(\xi). \quad (2)$$

Here J_ϵ is the smoothing operator given by $(J_\epsilon f)(x) = (\phi(\epsilon D)f)(x)$ with $\phi \in C_0^\infty(\mathbb{R}^n)$, $\phi(\xi) = 1$ for $|\xi| \leq 1$, and we take $\epsilon_j = 2^{-j\gamma}$.