

w/ Y. Kaku, A. Matsumura & Y. Michimura  
based on arXiv: 2308.14552

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Aug 30<sup>th</sup> @ Leipzig

# Towards the Optimal Experiment of Gravity-induced Quantum Entanglement



# Outline

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1. Introduction
2. Previous Proposals
3. General Analysis
4. Our Proposal
5. Summary

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1. Introduction

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5. Summary

## Key Question

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Is Gravity Quantum?

# Physics celebrities said...

Ricard Feynman



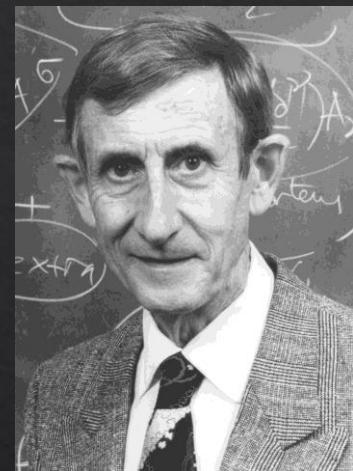
“Maybe we should **not** try to quantize gravity.”

Roger Penrose



“Quantum theory fits most **uncomfortably** with the curved space-time notion of the general relativity.”

Freeman Dyson



“Should quantum mechanics and GR be unified?  
**I don't think so.**  
Maybe, they should not be unified...”

## Key Question

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Is Gravity Quantum?

We aren't sure.

# Common sense?

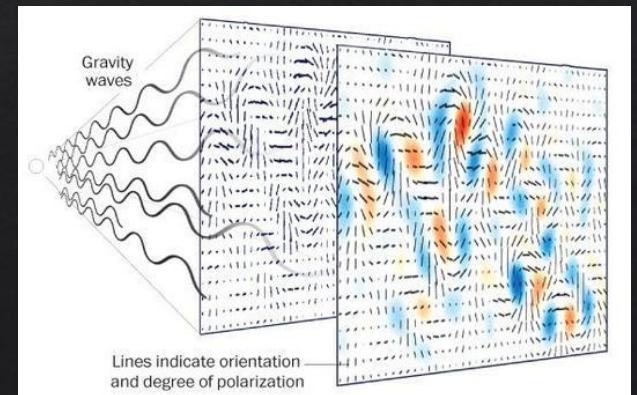
Nakayama+[0804.1827]

We often take **Quantum Gravity** for granted.

- ① Grav. fields can be in quantum superposition
- ② Graviton are quantized like QED.

(As a cosmologist, I often assume ② in my work)

Their validity has  
never been confirmed.



# That's science

---



Let's test it with experiments!

# Key Question

“Is Gravity Quantum?”



Testable

① Do weak gravitational fields  
become quantum superposition?

② graviton is (far) future step.

## Key Question

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Is Gravity Quantum?

We aren't sure.

Let's test it!

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2. Previous Proposals

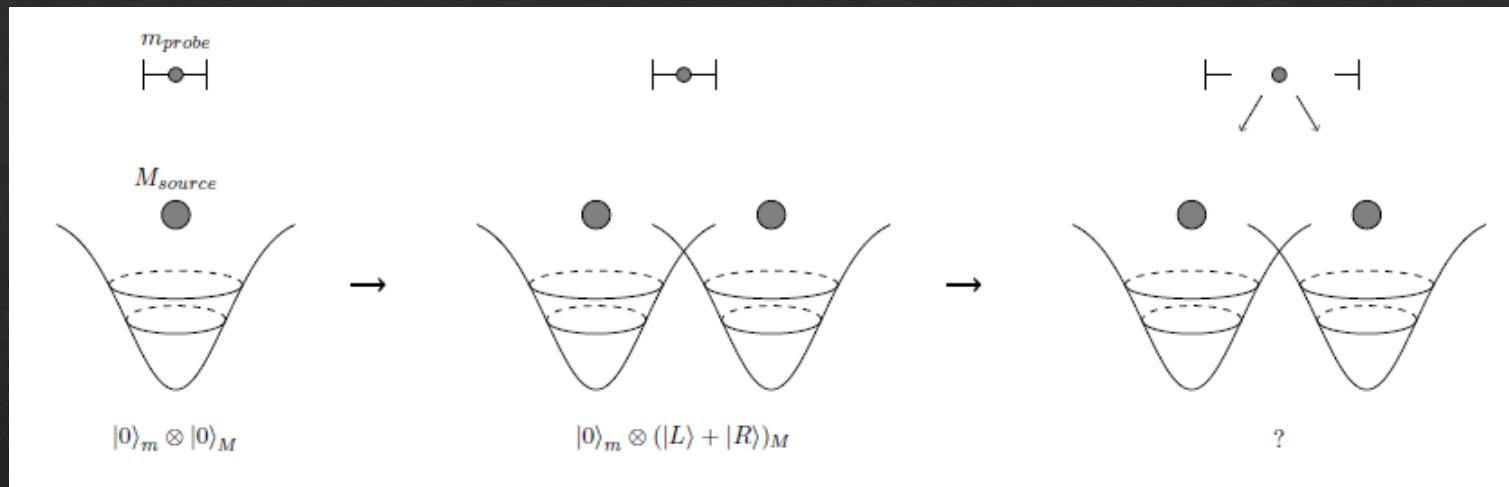
3. General Analysis

4. Our Proposal

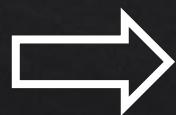
5. Summary

# Sketch of idea

Carney+[1807.11494]



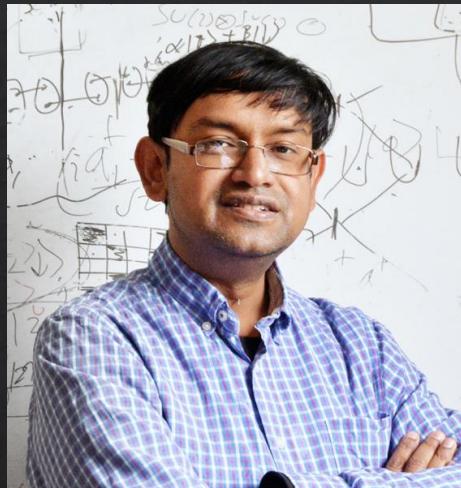
Does quantum superposition of a source mass lead to the **superposition of gravitational fields**?



We can check it with entanglement.

# Proposers

Sougato Bose et al.



Chiara Marletto

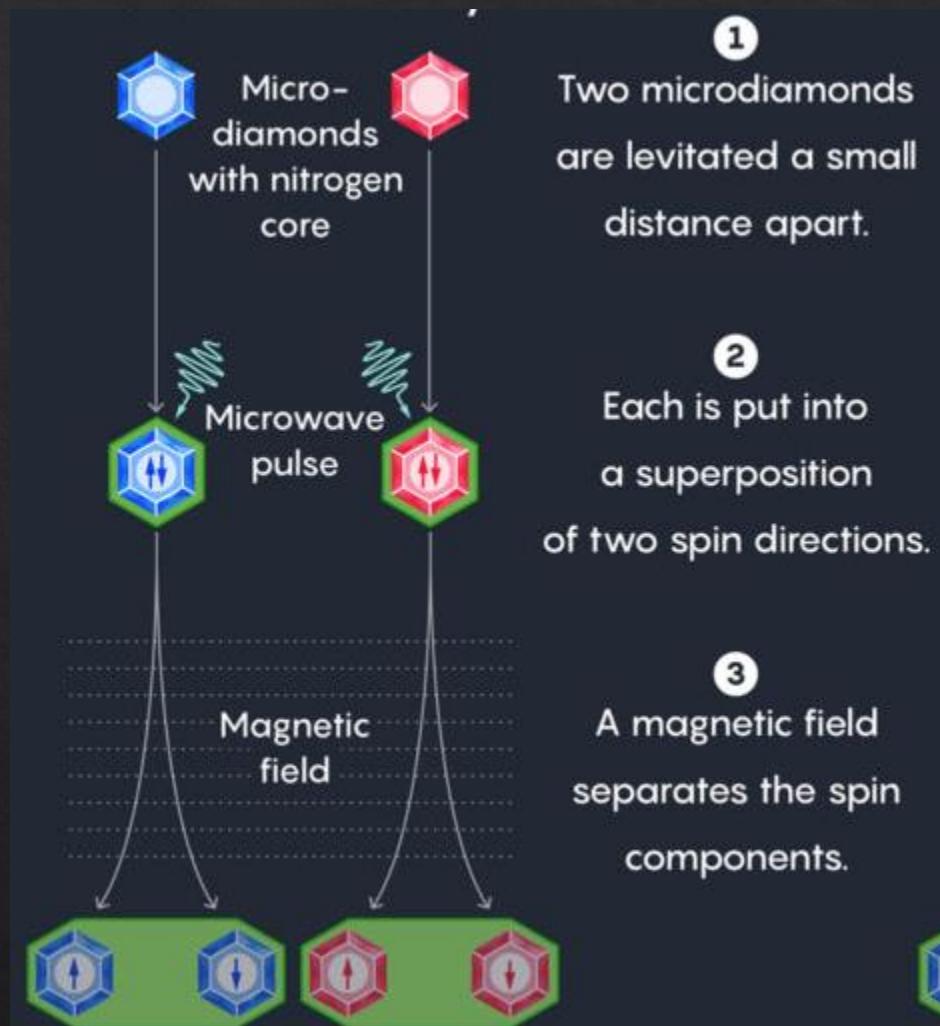


Vlatko Vedral

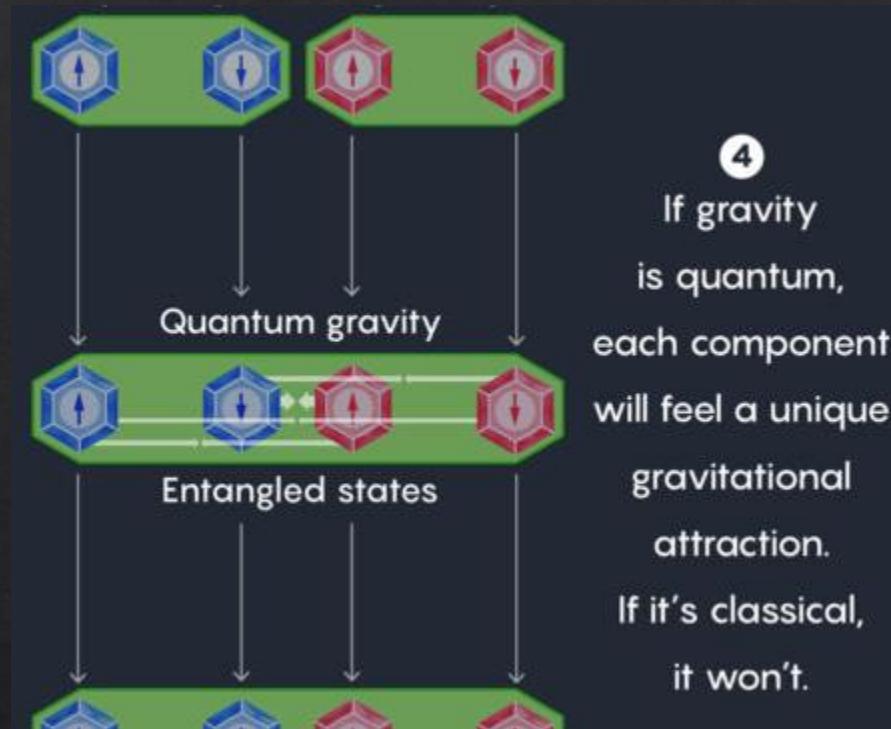


2 papers were published  
in PRL on the same day.  
= BMV proposal

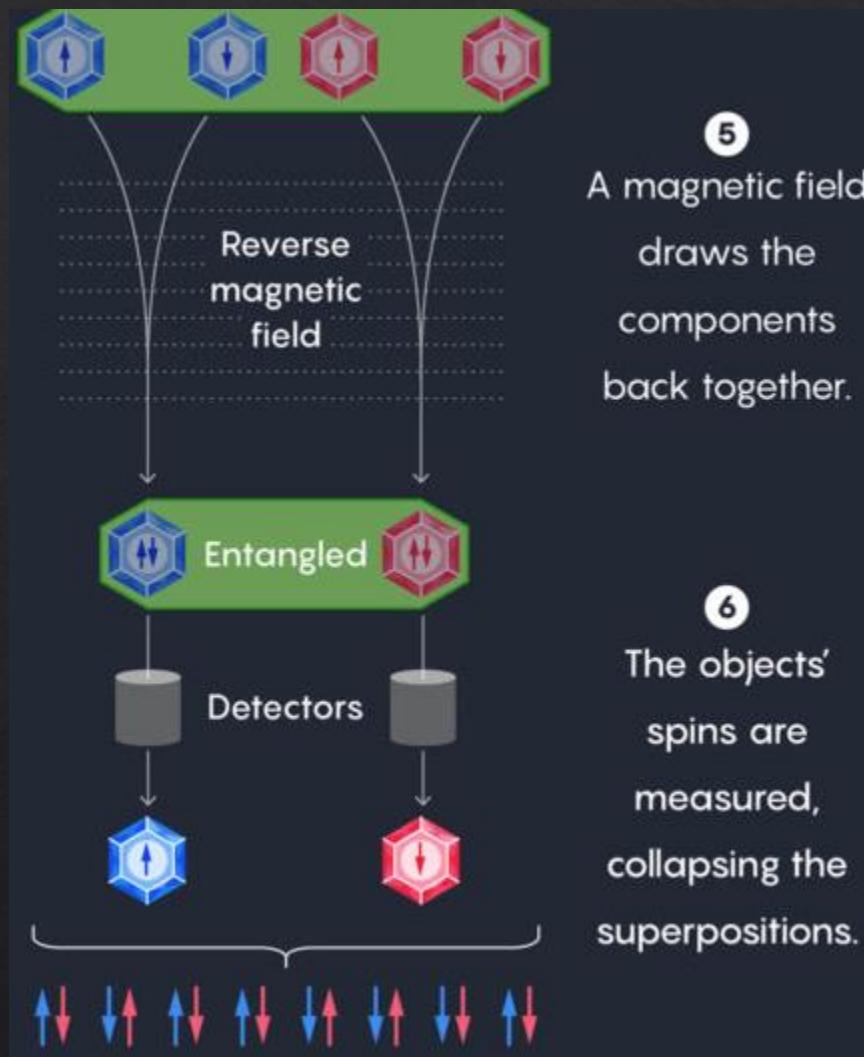
# BMV experiment



# BMV experiment



# BMV experiment



# Quantum state

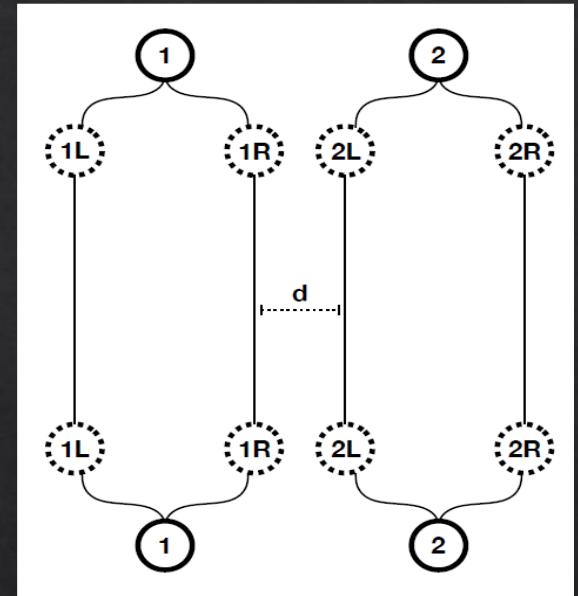
Christodoulou & Rovelli, PLB 792 (2019)[1808.05842]

The initial state is

$$\begin{aligned} |\Psi_1\rangle &= \frac{1}{2} \left( |\psi_1^L\rangle + |\psi_1^R\rangle \right) \otimes \left( |\psi_2^L\rangle + |\psi_2^R\rangle \right) \otimes |g\rangle. \\ &= \frac{1}{2} \left( |LL\rangle + |RR\rangle + |LR\rangle + |RL\rangle \right) \otimes |g\rangle. \end{aligned}$$

if GFs can be quantum superposition

$$\begin{aligned} |\Psi_2\rangle &= \frac{1}{2} \left( |LL\rangle \otimes |g_{d_{LL}}\rangle + |RR\rangle \otimes |g_{d_{RR}}\rangle \right. \\ &\quad \left. + |LR\rangle \otimes |g_{d_{LR}}\rangle + |RL\rangle \otimes |g_{d_{RL}}\rangle \right), \end{aligned}$$

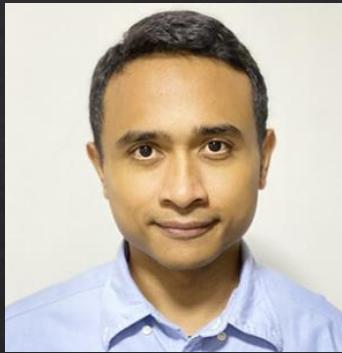


Only the nearest pair  $|RL\rangle$  gains a significant phase factor

$$\begin{aligned} |\Psi_3\rangle &= \frac{1}{2} \left( |LL g_{d_{LL}}\rangle + |RR g_{d_{RR}}\rangle \right. \\ &\quad \left. + |LR g_{d_{LR}}\rangle + e^{i \frac{Gm^2 t}{\hbar d}} |RL g_{d_{RL}}\rangle \right). \end{aligned}$$

# Another proposal

Krisnanda et al., npj Quant. Inf. 6,12 (2020)



Tanjung Krisnanda

Simple Procedure:

1. Trap two masses in a harmonic potential



# Another proposal

Krisnanda et al., npj Quant. Inf. 6,12 (2020)



Tanjung Krisnanda

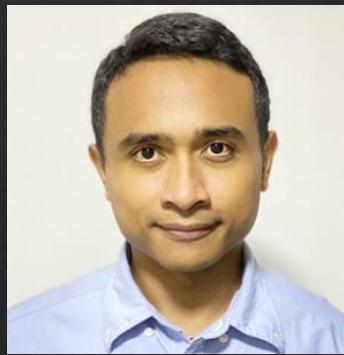
## Simple Procedure:

1. Trap two masses in a harmonic potential
2. Release and let them grav. interact



# Another proposal

Krisnanda et al., npj Quant. Inf. 6,12 (2020)

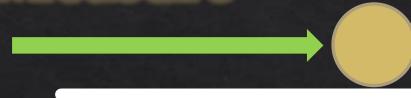


Tanjung Krisnanda

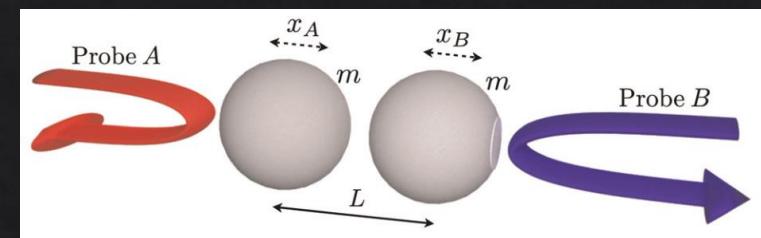
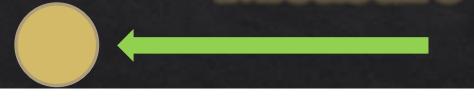
## Simple Procedure:

1. Trap two masses in a harmonic potential
2. Release and let them grav. interact
3. Measure the positions and momenta

Measure

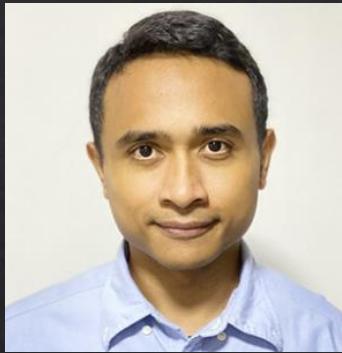


Measure



# Another proposal

Krisnanda et al., npj Quant. Inf. 6,12 (2020)



Tanjung Krisnanda

Simple Procedure:

1. Trap two masses in a harmonic potential



Wavefunction

# Another proposal

Krisnanda et al., npj Quant. Inf. 6,12 (2020)



Tanjung Krisnanda

## Simple Procedure:

1. Trap two masses in a harmonic potential
2. Release and let them grav. interact



Spread out

# Another proposal

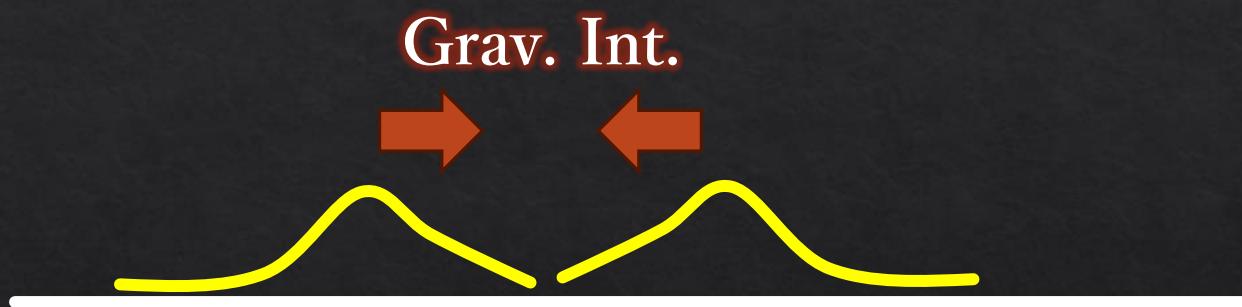
Krisnanda et al., npj Quant. Inf. 6,12 (2020)



Tanjung Krisnanda

## Simple Procedure:

1. Trap two masses in a harmonic potential
2. Release and let them grav. interact
3. Measure the positions and momenta



Entangled

# Feasibility

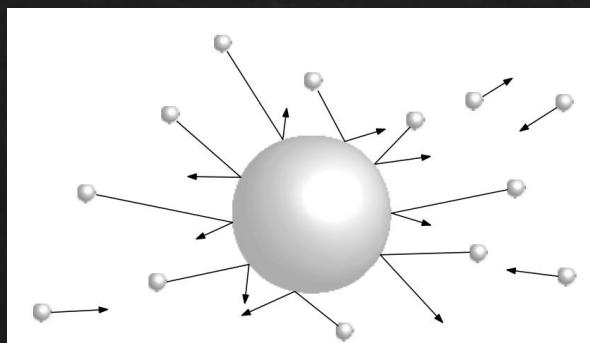
Rijavec et al., New J. Phys. 23 043040 (2021)



Simone Rijavec

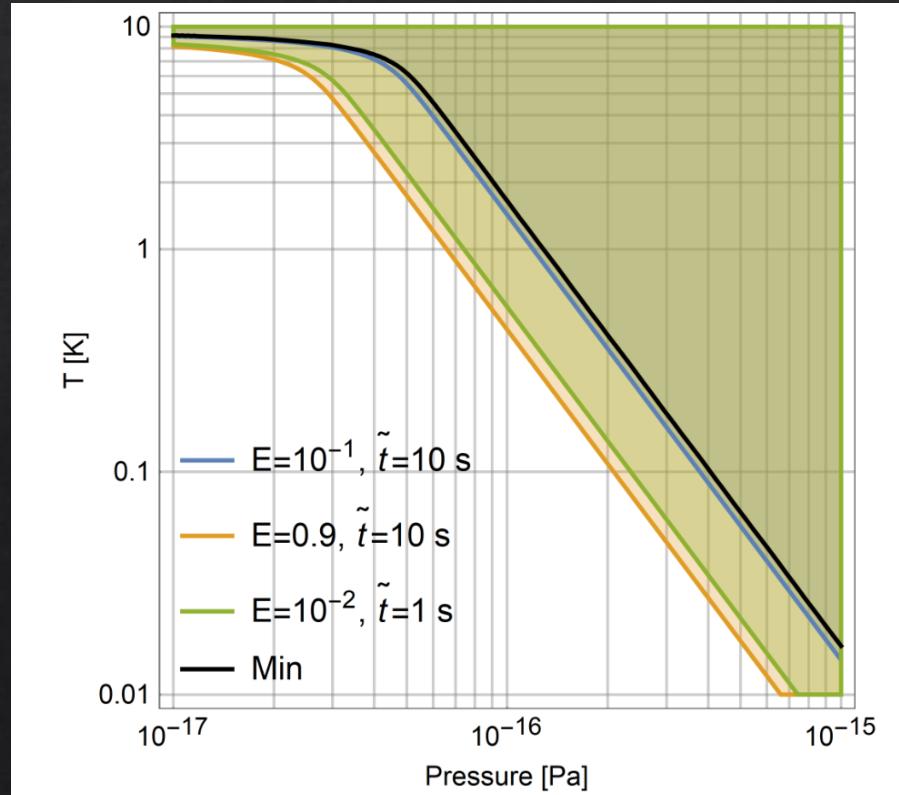


## Air molecule Scattering



For a real experiment,

1. Ultra-high vacuum to avoid decoherence



# Feasibility

Rijavec et al., New J. Phys. 23 043040 (2021)



Simone Rijavec



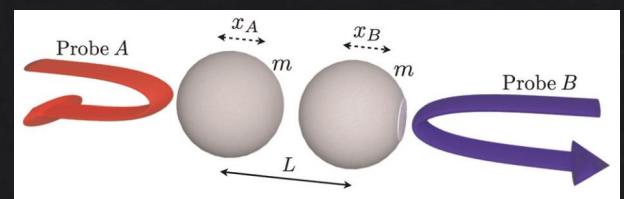
For a real experiment,

1. Ultra-high vacuum to avoid decoherence
2. Free-fall problem



Free-fall

40m down for 3 sec



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Our quadratic Hamiltonian:

$$H = \frac{p_1^2}{2m} + \frac{1}{2}k_1x_1^2 + \frac{p_2^2}{2m} + \frac{1}{2}k_2x_2^2 - \frac{Gm^2}{d^3}(x_1 - x_2)^2,$$

oscillator1

oscillator2

Grav. Int.

$(d \gg |x_1 - x_2|)$

The system is quadratic.  
→ Exactly Solvable!

Our quadratic Hamiltonian:

$$H = \frac{p_1^2}{2m} + \frac{1}{2}k_1x_1^2 + \frac{p_2^2}{2m} + \frac{1}{2}k_2x_2^2 - \frac{Gm^2}{d^3}(x_1 - x_2)^2,$$

Spring constant  $k_i$



Potential parameter:  $\lambda_i \equiv k_i/m\omega^2$

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$\lambda = 1$ : Harmonic



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Potential parameter:  $\lambda_i \equiv k_i/m\omega^2$

$\lambda = 1$ : Harmonic

$\lambda = 0$ : Free mass



Our quadratic Hamiltonian:

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Potential parameter:  $\lambda_i \equiv k_i/m\omega^2$

$\lambda = 1$ : Harmonic

$\lambda = 0$ : Free mass

$\lambda = -1$ : Inverted



# Experimental Goal

TF. et al. (2023) [2308.14552]

Good indicator of entanglement:

## Logarithmic Negativity $E_N$

- $E_N > 0 \Leftrightarrow$  Two oscillators are entangled
- Larger  $E_N$  indicates larger entanglement
- $E_N = 0.01$  is experimentally detectable.

$$E_N \equiv \max [0, -\log_2 (2\tilde{\nu}_{\min})] \quad \tilde{\nu}_{\min} \equiv \left[ \frac{1}{2} \left( \tilde{\Sigma} - \sqrt{\tilde{\Sigma}^2 - 4 \det \sigma} \right) \right]^{1/2}$$

$$u_i(t) = (X_1(t), P_1(t), X_2(t), P_2(t)),$$
$$\sigma_{ij}(t) = \frac{1}{2} \langle u_i(t) u_j(t) + u_j(t) u_i(t) \rangle.$$

$$\sigma(t) = \begin{bmatrix} \sigma_1 & \sigma_3 \\ \sigma_3^T & \sigma_2 \end{bmatrix} \quad \tilde{\Sigma} \equiv \det \sigma_1 + \det \sigma_2 - 2 \det \sigma_3$$

# Calculation

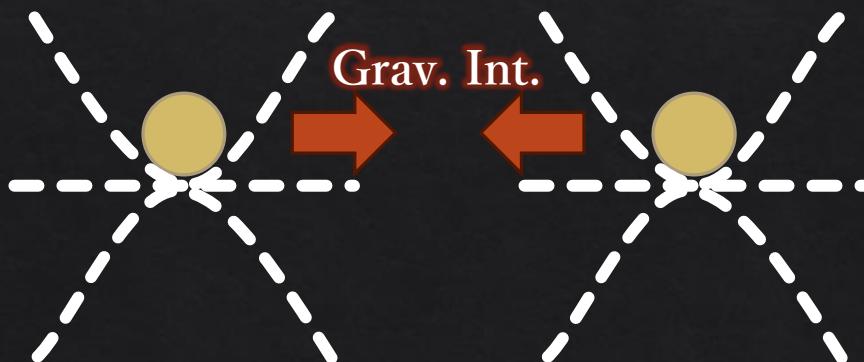
TF. et al. (2023) [2308.14552]

We compute  $E_N$  when

At  $t = 0$ ,  
they're in the  
ground state  
w/o gravity

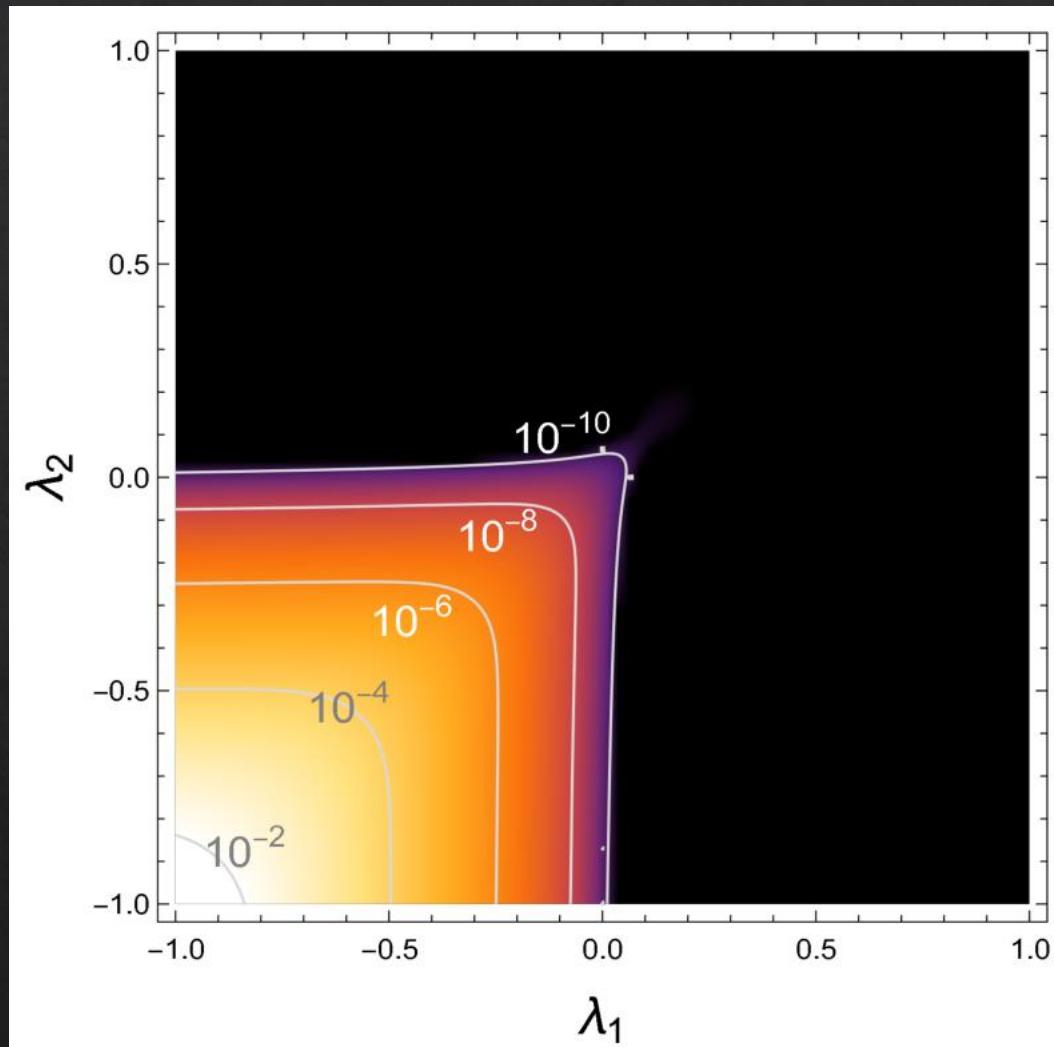


For  $t > 0$ ,  
they evolve  
in the  $\lambda_i$  potential  
w/ gravity



# Result of Entanglement

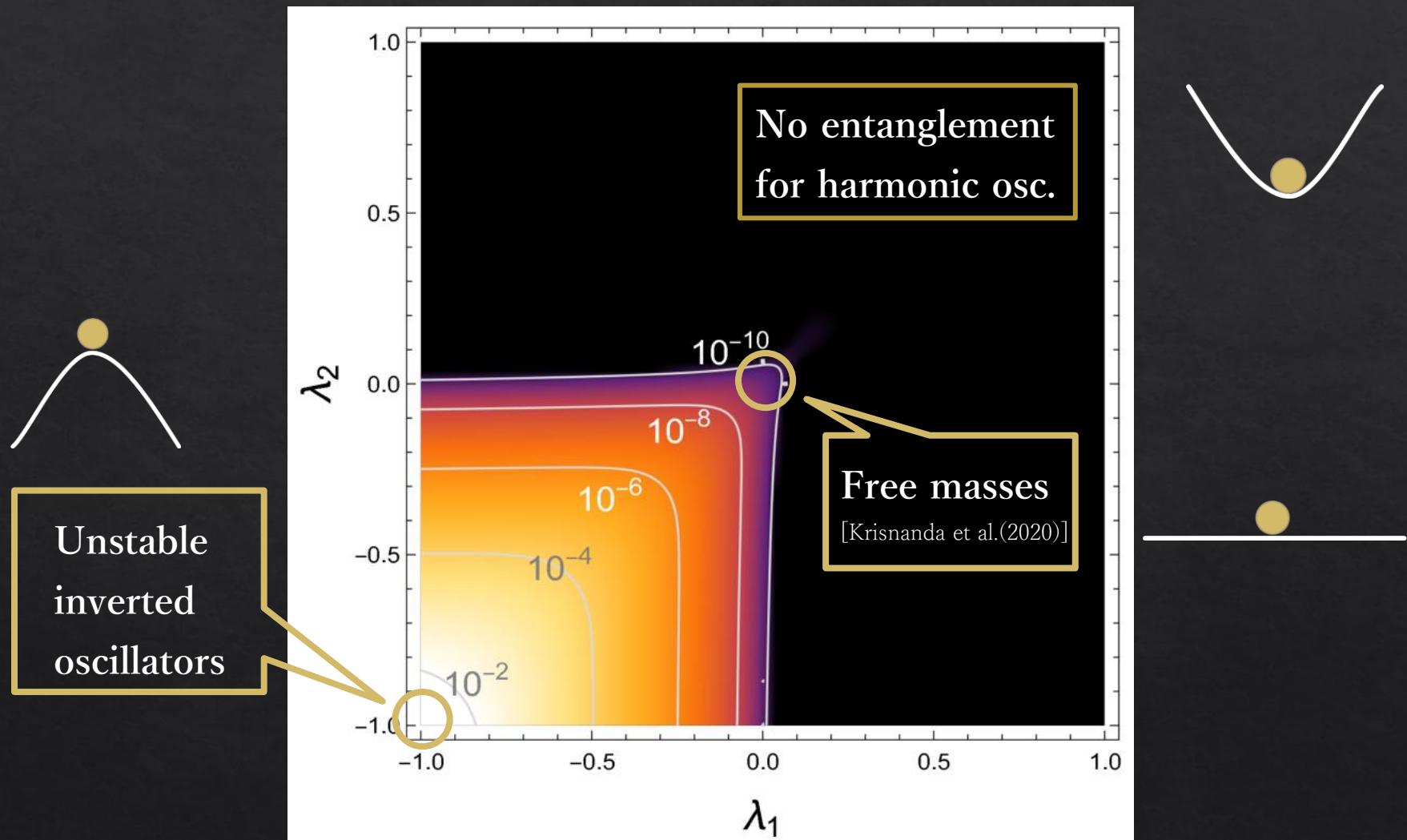
TF. et al. (2023) [2308.14552]



Contour of  $E_N$  ( $\omega t = 13$ ,  $\eta = 2\mu = 10^{-12}$  )

# Result of Entanglement

TF. et al. (2023) [2308.14552]



# Including Decoherence

TF. et al. (2023) [2308.14552]

Heisenberg-Langevin eqs:

$$\dot{X}_i = \omega P_i, \quad \dot{P}_i = -\lambda\omega X_i + \omega\eta(X_i - X_j) + \xi_i,$$

$\xi_i$ : random noise force  $\Rightarrow$  decoherence

$$\frac{1}{2} \langle \xi_i(t)\xi_j(t') + \xi_i(t')\xi_j(t) \rangle = \mu\omega\delta(t-t')\delta_{ij}.$$

$\mu$ : size of env. fluctuation



$\eta$ : grav. coupling constant

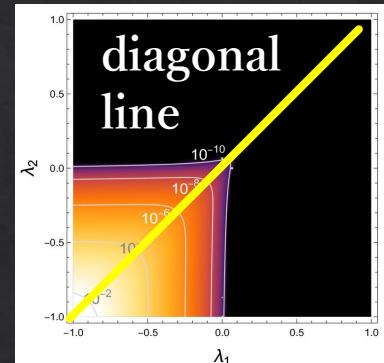
$$\eta \equiv \frac{2Gm}{\omega^2 d^3} = 2.7 \times 10^{-13} \omega_{\text{kHz}}^{-2} \left( \frac{m/d^3}{2 \text{ g/cm}^3} \right)$$

# Analytic Solution

TF. et al. (2023) [2308.14552]

For the identical oscillators ( $\lambda_1 = \lambda_2$ )  
 Logarithmic Negativity reads

$$E_N \simeq 3(\eta - \mu)f_{\text{gra}} \quad (\lambda \leq 0)$$



Grav. coupling constant

Random noise parameter

Power-law

$$f_{\text{gra}} \simeq \begin{cases} \frac{1}{2} |\sin(\omega t)| & (\lambda = 1) \\ \frac{1}{6} (\omega t)^3 & (\lambda = 0) \\ \frac{1}{8} e^{2\omega t} & (\lambda = -1) \end{cases}$$



Exponential

The time required to generate observable  $E_N = 0.01$

$$\tau_{\text{ent}} \simeq \begin{cases} \boxed{\text{w/o } \mu} & \begin{aligned} 4.2 \omega_{\text{kHz}}^{-1/3} \text{ sec} & (\lambda = 0) \\ 1.3 \times 10^{-2} \omega_{\text{kHz}}^{-1} \text{ sec} & (\lambda = -1) \end{aligned} \end{cases}$$

**300 times faster!**

Including decoherence parameter  $\mu$

$$\tau_{\text{ent}} \simeq \begin{cases} 4.2 \omega_{\text{kHz}}^{-1/3} [\eta/(\eta - \mu)]^{1/3} \text{ sec} & (\lambda = 0) \\ 1.3 \times 10^{-2} \omega_{\text{kHz}}^{-1} \text{ sec} & (\lambda = -1) \\ + \log[\eta/(\eta - \mu)]/(2\omega) & \end{cases}$$

$\simeq \mathcal{O}(10^{-3}) \omega_{\text{kHz}}^{-1} \text{ sec}$

The time required to generate observable  $E_N = 0.01$

The inverted oscillators generate  
the gravity-induced entanglement  
most quickly and are most resistant  
to decoherence.



$$\simeq \mathcal{O}(10^{-3})\omega_{\text{kHz}}^{-1} \text{sec}$$

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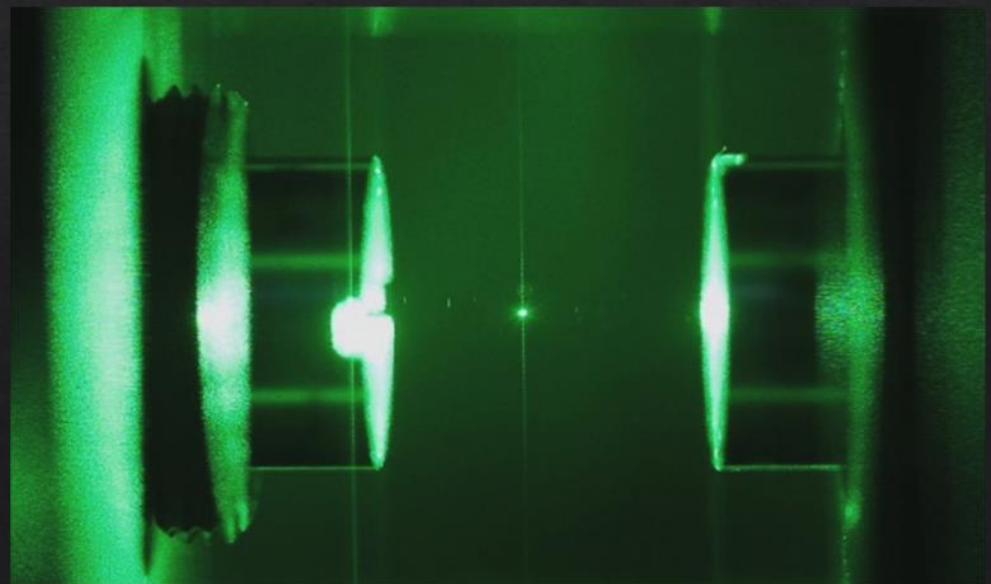
- Free fall problem



Free-fall  
40m down for 3 sec

- Optical levitation

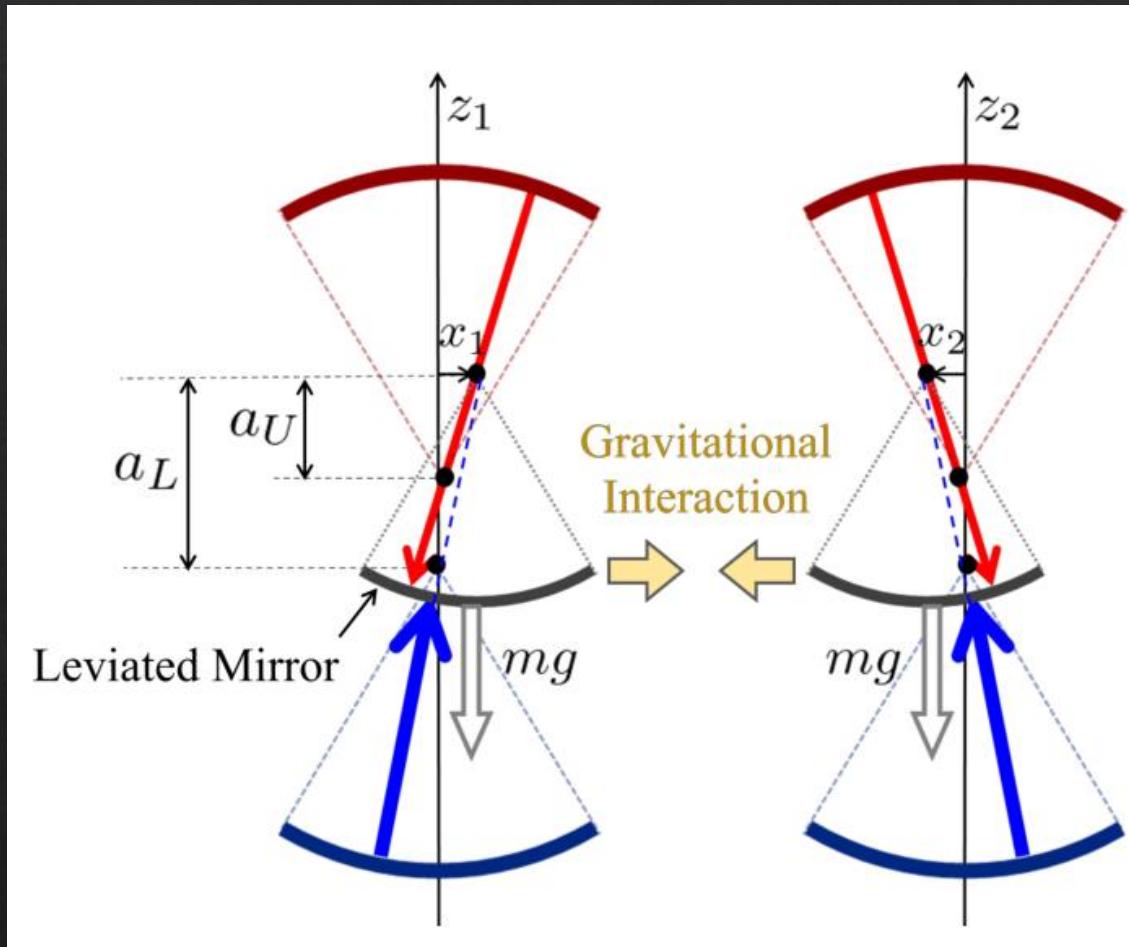
demonstrated to levitate small  
particle by laser pressure  
w/o mechanical support



# Sandwich Setup

Michimura. et al. (2017)

Inverted Oscillator  $\Leftrightarrow$  Anti-spring effect

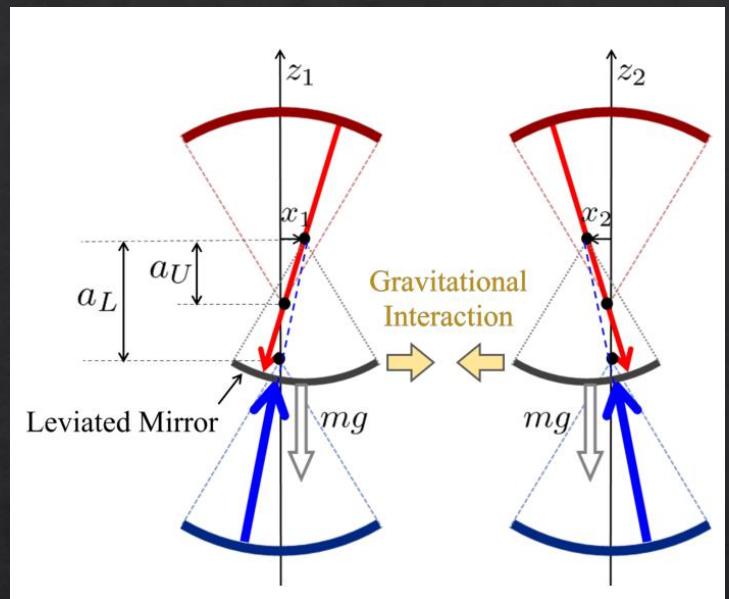


# Sandwich Setup

Michimura. et al. (2017)

We can realize high frequency inverted oscillator

$$\begin{aligned}\omega_{\text{hor}}^2 &= \frac{2}{mc} \left( \frac{P_U}{a_U} - \frac{P_L}{a_L} \right) = \frac{2(a_L - a_U)}{mc a_U a_L} P_L - \frac{g}{a_U}, \\ &\simeq -(1\text{kHz})^2 \left( \frac{m}{0.1\text{mg}} \right)^{-1} \left( \frac{P_L}{30\text{kW}} \right) \left( \frac{a_L}{2\text{mm}} \right)^{-1},\end{aligned}$$



while suppressing decoherence due to photon shot noise.

$$\begin{aligned}\mu_{\text{shot,hor}} &= \frac{16\omega_\ell P_L}{m\omega^2 c^2 T_{\text{in}}} \left( \frac{\Delta x}{a_L} \right)^2 \simeq \frac{8\omega_\ell \Delta x^2}{ca_L T_{\text{in}}}, \\ &= 2.5 \times 10^{-14} \omega_{\text{kHz}}^3 \left( \frac{a_L}{2\text{mm}} \right)^{-1} \left( \frac{m}{0.1\text{mg}} \right)^{-1} \left( \frac{\omega_{\text{in}}}{1\text{MHz}} \right)^{-2}\end{aligned}$$

# Summary

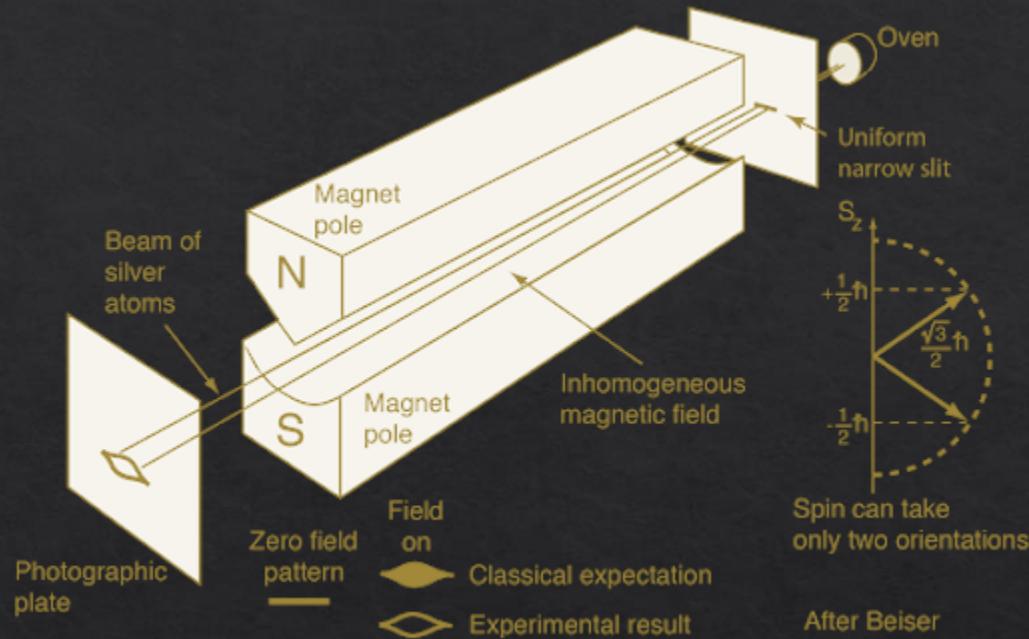
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- Lack of **experimental verification** of quantum gravity.  
Not even sure if grav. fields can be quantum superposition.
- Many proposals to test gravity-induced entanglement. “**Trap & release**” masses generates entanglement. (free-fall problem)
- We analyzed two general quadratic oscillators coupled by gravity and found **inverted oscillators** exponentially generate entanglement and resistant to decoherence.
- As an experimental implementation, we proposed **levitated mirror with anti-spring effect** in a sandwich configuration.

Thank you



# Sketch of idea



Stern-Gerlach experiment enables us to prepare the quantum superposition of a mass at two different locations.

# Superposition

Pure state:  $|\Psi\rangle = c_1|\phi_1\rangle + c_2|\phi_2\rangle$   
quantum superposition

It's undetermined whether  $\Psi = \phi_1$  or  $\phi_2$  (c.f. Schrodinger's cat)



It's pre-determined whether  $\Psi = \phi_1$  or  $\phi_2$

The probabilities of each realization  $p_1$  and  $p_2$  are known.

Its QM description = **Mixed State**

# Density matrix

$\hat{\rho}$  gives the probability and the expected value

$$p_i = \text{Tr}[\hat{P}_i \hat{\rho}] \quad \langle \hat{O} \rangle = \text{Tr}[\hat{O} \hat{\rho}]$$

quantum

$$\hat{\rho}_{\text{pure}} = |\Psi\rangle\langle\Psi| = \sum_i |c_i|^2 |\phi_i\rangle\langle\phi_i| + \sum_{i \neq j} c_i c_j^* |\phi_i\rangle\langle\phi_j|$$

classical

Interference term

$$\hat{\rho}_{\text{mix}} = \sum_i p_i |\phi_i\rangle\langle\phi_i|$$

The essential difference btw quantum and classical state appears in the interference term in the density matrix.

# Entanglement

A quantum system consists of subsystem A & B.

General state  $|\Psi\rangle = \sum_{ij} c_{ij} |\psi_i\rangle_A \otimes |\phi_j\rangle_B$

If  $|\psi\rangle_A = \sum_i a_i |\psi_i\rangle_A$  and  $|\phi\rangle_B = \sum_j b_j |\phi_j\rangle_B$  independently,

Separable state state  $|\Psi\rangle = \sum_{ij} a_i b_j |\psi_i\rangle_A \otimes |\phi_j\rangle_B$

Non separable = Entangled state

Interaction btw the subsystems can induce entanglement.

# Trace out

---

Remember  $\langle \hat{O} \rangle = \text{Tr}[\hat{O}\hat{\rho}]$

If we only consider observables of the subsystem A,  $\hat{O}_A$ ,  
we take **the trace of  $\hat{\rho}$**  over the subsystem B,

Reduced density matrix:  $\hat{\rho}_A = \text{Tr}_B[\hat{\rho}]$

This operation won't change anything in A,  
if A & B are separable.

# Decoherence

Pure entangled state

$$|\Psi\rangle = \sum_i c_i |\psi_i\rangle_A \otimes |\phi_i\rangle_B$$

Density matrix

$$\hat{\rho} = |\Psi\rangle\langle\Psi| = \sum_{ik} c_i c_k^* |\psi_i\rangle\langle\psi_k| \otimes |\phi_i\rangle\langle\phi_k|$$

Trace out:

$$\text{Tr}_B[\hat{\rho}] = \sum_l \langle \phi_l | \hat{\rho} | \phi_l \rangle = \sum_i |c_i|^2 |\psi_i\rangle\langle\psi_i|$$

When the original state is entangled,  
tracing out it into a mixed state.

**Decoherence** = interference terms (quantum-ness) vanish

# Phase from potential

Christodoulou & Rovelli, PLB 792 (2019)[1808.05842]

Schrodinger eq.:  $i\partial_t |\psi\rangle = \hat{H}|\psi\rangle$

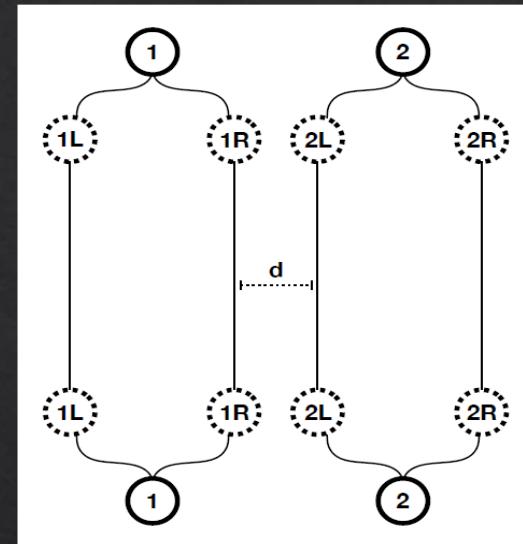
⇒ Phase from mass energy  $e^{-imt} |\psi\rangle$

GR replaces  $t$  by the proper time  $s$

Newtonian:  $ds^2 = [1 + 2\Phi]dt^2, \quad \Phi = -\frac{Gm}{d}$

The relative phase that  $|RL\rangle$  gains is

Phase:  $\phi = \frac{Gm^2}{d} t \approx 2\pi \left(\frac{t}{1\text{sec}}\right) \left(\frac{d}{1\text{mm}}\right)^{-1} \left(\frac{m}{10^{-11}\text{g}}\right)^2$



c.f.  $m_P \approx 2 \times 10^{-5}\text{g}$   
 $c * \text{sec} \approx 3 \times 10^8\text{m}$

Small mass  
compensated  
by long time

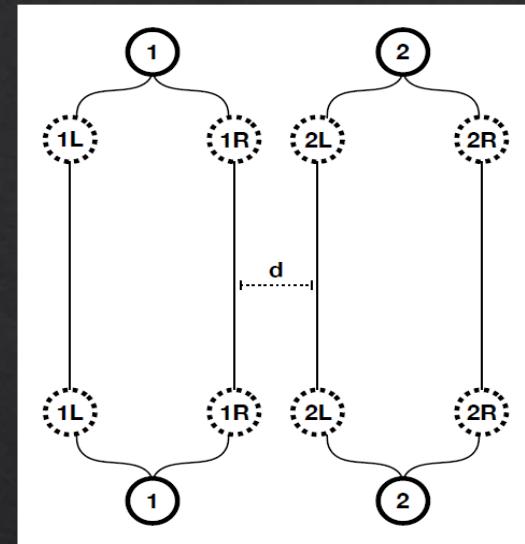
# Quantum state

Christodoulou & Rovelli, PLB 792 (2019)[1808.05842]

$$|\Psi_3\rangle = \frac{1}{2} \left( |LL g_{d_{LL}}\rangle + |RR g_{d_{RR}}\rangle + |LR g_{d_{LR}}\rangle + e^{i\frac{Gm^2 t}{\hbar d}} |RL g_{d_{RL}}\rangle \right).$$

Bring them back by inverse-SG

$$|\psi_4\rangle = \frac{1}{2} [ |\downarrow\downarrow\rangle + |\uparrow\uparrow\rangle + |\downarrow\uparrow\rangle + e^{i\phi} |\uparrow\downarrow\rangle ]$$



The entangled state is tested by Bell inequality

$$\mathcal{W} = |\langle \sigma_x^{(1)} \otimes \sigma_z^{(2)} \rangle - \langle \sigma_y^{(1)} \otimes \sigma_z^{(2)} \rangle|$$

If  $\mathcal{W} > 1$ , the state is entangled and GFs are superposed.

# General Hamiltonian

TF. et al. (2023) [2308.14552]  
]

Our quadratic Hamiltonian:

$$H = \frac{p_1^2}{2m} + \frac{1}{2}k_1x_1^2 + \frac{p_2^2}{2m} + \frac{1}{2}k_2x_2^2 - \frac{Gm^2}{d^3}(x_1 - x_2)^2,$$



dimensionless form

$$H = \frac{\omega}{2} \left[ P_1^2 + \lambda_1 X_1^2 + P_2^2 + \lambda_2 X_2^2 - \eta(X_1 - X_2)^2 \right]$$

oscillator1

oscillator2

Grav. Int.

Variable:  $P_i \equiv p_i/\sqrt{\hbar m \omega}$      $X_i \equiv \sqrt{m\omega/\hbar}x_i$

Coupling  
constant:

$$\eta \equiv \frac{2Gm}{\omega^2 d^3} = 2.7 \times 10^{-13} \omega_{\text{kHz}}^{-2} \left( \frac{m/d^3}{2 \text{ g/cm}^3} \right)$$

*Gravity  
is weak*