

# The weak-gravity bound in asymptotically safe gauge-gravity systems

ERG 2022, Berlin

July 27, 2022

---

**Marc Schiffer**, Perimeter Institute

based on

A. Eichhorn, J. Kwapisz, MS: **Phys.Rev.D 105 (2022) 106022**



- Some key questions for gravity-matter systems:

- Some key questions for gravity-matter systems:
  - ▶ Does the gravity fixed-point allow for the inclusion of SM-matter?  
[\[Doná, Eichhorn, Percacci; 2013\]](#), [\[Meibohm, Pawłowski, Reichert; 2016\]](#), [\[Biemanns, Platania, Saueressig; 2017\]](#), . . .

- Some key questions for gravity-matter systems:
  - ▶ Does the gravity fixed-point allow for the inclusion of SM-matter?  
[\[Doná, Eichhorn, Percacci; 2013\]](#), [\[Meibohm, Pawłowski, Reichert; 2016\]](#), [\[Biemanns, Platania, Saueressig; 2017\]](#), . . .
  - ▶ Does asymptotically safe quantum gravity support a UV-complete matter sector?  
Does a UV-complete matter sector pose constraints on gravity?  
[\[Harst, Reuter; 2011\]](#), [\[Eichhorn, Held, Pawłowski; 2016\]](#), [\[Christiansen, Eichhorn; 2017\]](#), . . .

- Some key questions for gravity-matter systems:
  - ▶ Does the gravity fixed-point allow for the inclusion of SM-matter?  
[Doná, Eichhorn, Percacci; 2013], [Meibohm, Pawłowski, Reichert; 2016], [Biemanns, Platania, Saueressig; 2017], . . .
  - ▶ Does asymptotically safe quantum gravity support a UV-complete matter sector?  
Does a UV-complete matter sector pose constraints on gravity?  
[Harst, Reuter; 2011], [Eichhorn, Held, Pawłowski; 2016], [Christiansen, Eichhorn; 2017], . . .
  - ▶ Is there a viable phenomenology?  
[Shaposhnikov, Wetterich; 2009], [Harst, Reuter; 2011], [Eichhorn, Held; 2017, 2018], [Eichhorn, Versteegen; 2017], . . .  
[Draper, Knorr, Ripken, Saueressig; 2020], [Knorr, Pirlo, Ripken, Saueressig; 2022]  
[Reichert, Smirnov; 2019], [Kowalska, Sessolo; 2020], [Eichhorn, Pauly; 2020], . . .
  - ▶ . . .

# The "weak gravity bound"

- Specifically in asymptotically safe gravity:
  - ▶ There exist indications that metric fluctuations must not be too strong.
  - ▶ Interacting nature of gravity induces novel interactions in the matter sector.  
[\[Eichhorn and Gies, 2011\]](#), [\[Eichhorn, 2012\]](#), [\[Meibohm and Pawłowski, 2016\]](#), [\[Eichhorn, Held and Pawłowski, 2016\]](#),  
[\[Christiansen and Eichhorn, 2017\]](#), [\[Eichhorn and Held, 2017\]](#), [\[Eichhorn, Lippoldt and Skinjar, 2017\]](#)  
[\[Eichhorn, Lippoldt and MS, 2018\]](#)
  - ▶ Beyond the weak-gravity regime, metric fluctuations can induce novel divergences in these interactions.

- Example: Abelian gauge field  $A_\mu$

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$

- ▶ From kinetic term:

$$S_{\text{kin}} = \frac{Z_A}{4} \int d^d x \sqrt{g} F_{\mu\nu} F^{\mu\nu}$$

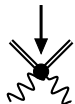
# Induced interactions

- Example: Abelian gauge field  $A_\mu$

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$

- ▶ From kinetic term:

$$S_{\text{kin}} = \frac{Z_A}{4} \int d^d x \sqrt{g} F_{\mu\nu} F^{\mu\nu}$$





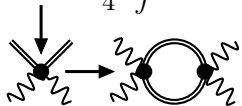
# Induced interactions

- Example: Abelian gauge field  $A_\mu$

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$

- ▶ From kinetic term:

$$S_{\text{kin}} = \frac{Z_A}{4} \int d^d x \sqrt{g} F_{\mu\nu} F^{\mu\nu}$$



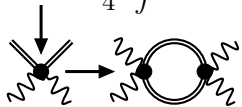
# Induced interactions

- Example: Abelian gauge field  $A_\mu$

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$

- ▶ From kinetic term:

$$S_{\text{kin}} = \frac{Z_A}{4} \int d^d x \sqrt{g} F_{\mu\nu} F^{\mu\nu}$$



$$S_{\text{int}} = \frac{w_2 k^{-d}}{8} \int d^d x \sqrt{g} (F_{\mu\nu} F^{\mu\nu})^2$$

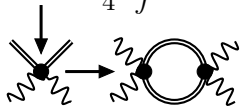
# Induced interactions

- Example: Abelian gauge field  $A_\mu$

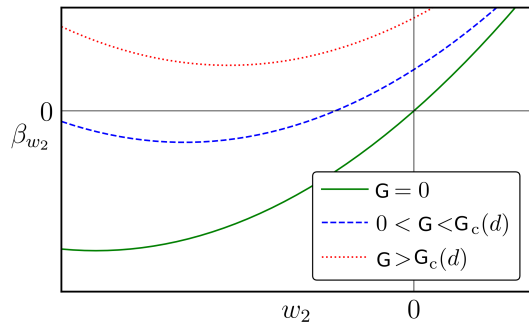
$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$

- ▶ From kinetic term:

$$S_{\text{kin}} = \frac{Z_A}{4} \int d^d x \sqrt{g} F_{\mu\nu} F^{\mu\nu}$$



$$S_{\text{int}} = \frac{w_2 k^{-d}}{8} \int d^d x \sqrt{g} (F_{\mu\nu} F^{\mu\nu})^2 + \frac{v_2 k^{-d}}{8} \int d^d x \sqrt{g} (F_{\mu\nu} \tilde{F}^{\mu\nu})^2$$



- Schematically:

$$\beta_{w_2} = B_0(G) + w_2 B_1(G) + w_2^2 B_2$$

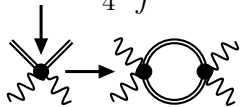
# Induced interactions

- Example: Abelian gauge field  $A_\mu$

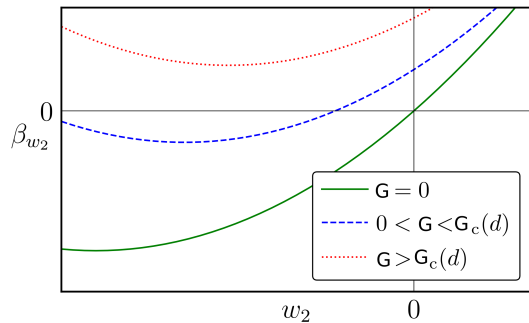
$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$

- ▶ From kinetic term:

$$S_{\text{kin}} = \frac{Z_A}{4} \int d^d x \sqrt{g} F_{\mu\nu} F^{\mu\nu}$$



$$S_{\text{int}} = \frac{w_2 k^{-d}}{8} \int d^d x \sqrt{g} (F_{\mu\nu} F^{\mu\nu})^2 + \frac{v_2 k^{-d}}{8} \int d^d x \sqrt{g} (F_{\mu\nu} \tilde{F}^{\mu\nu})^2$$

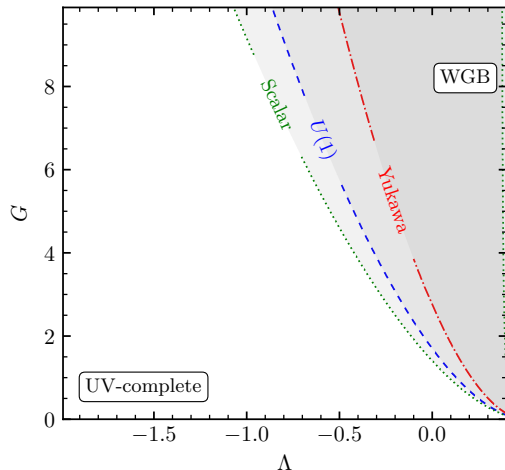


- Schematically:

$$\beta_{w_2} = B_0(G) + w_2 B_1(G) + w_2^2 B_2$$

- $\exists$  real FP only for  $B_0 \leq \frac{B_1^2(G)}{4B_2}$

# Excluded strong gravity regime: The weak-gravity bound



[MS; 2021]

- WGB explored for:

- ▶ Scalar interactions

[Eichhorn; 2012], [de Brito, Eichhorn, Robson Linos dos Santos; 2021], [Laporte, Pereira, Saueressig, Wang; 2021], [Knorr; 2022]

- ▶ Fermionic interactions

[Eichhorn, Gies; 2011], [Eichhorn, Lippoldt, MS; 2018]  
[de Brito, Eichhorn, MS; 2020]

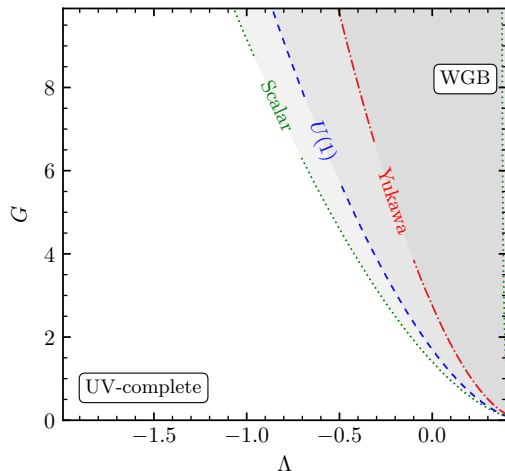
- ▶ Gauge interactions

[Christiansen, Eichhorn; 2017], [Eichhorn, MS; 2019]  
[Eichhorn, Kwapisz, MS; 2021]

- ▶ Scalar-Fermion interactions

[Eichhorn, Held; 2017]

# Excluded strong gravity regime: The weak-gravity bound



[MS; 2021]

- WGB explored for:
  - ▶ Scalar interactions  
[Eichhorn; 2012], [de Brito, Eichhorn, Robson Linos dos Santos; 2021], [Laporte, Pereira, Saueressig, Wang; 2021], [Knorr; 2022]
  - ▶ Fermionic interactions  
[Eichhorn, Gies; 2011], [Eichhorn, Lippoldt, MS; 2018]  
[de Brito, Eichhorn, MS; 2020]
  - ▶ Gauge interactions  
[Christiansen, Eichhorn; 2017], [Eichhorn, MS; 2019]  
[Eichhorn, Kwapisz, MS; 2021]
  - ▶ Scalar-Fermion interactions  
[Eichhorn, Held; 2017]
- further "pheno" consequences  
[Eichhorn, MS; 2019], [Hamada, Pawłowski, Yamada; 2020]  
[de Brito, Eichhorn, MS; 2020], . . .

## Two species of gauge fields

---

- Gravity is "blind" to internal symmetries: induced interactions for Abelian and non-Abelian gauge fields are the same (at the fixed point)

# Two species of gauge fields

- Gravity is "blind" to internal symmetries: induced interactions for Abelian and non-Abelian gauge fields are the same (at the fixed point)
- Interaction structure (involving four gauge fields and derivatives):

$$\Gamma_k^{U(1) \times U(1)} = \frac{1}{4} \int d^4x \sqrt{g} F_{\mu\nu}^a F^{\mu\nu, a} + S_{\text{gf}, A}$$



# Two species of gauge fields

- Gravity is "blind" to internal symmetries: induced interactions for Abelian and non-Abelian gauge fields are the same (at the fixed point)
- Interaction structure (involving four gauge fields and derivatives):

$$\begin{aligned}\Gamma_k^{U(1)\times U(1)} = & \frac{1}{4} \int d^4x \sqrt{g} F_{\mu\nu}^a F^{\mu\nu, a} + S_{\text{gf}, A} \\ & + \frac{k^{-4}}{16} \int d^4x \sqrt{g} \left( w_2 [F_{\mu\nu}^a F^{\mu\nu, a}]^2 + y_2 (F_{\mu\nu}^a F^{\mu\nu, b})(F_{\rho\sigma}^a F^{\rho\sigma, b}) \right. \\ & \quad \left. + x_2 [F_{\mu\nu}^2 F^{\mu\nu, 2}]^2 + z_2 (F_{\mu\nu}^1 F^{\mu\nu, 2})(F_{\rho\sigma}^1 F^{\rho\sigma, 2}) \right) \\ & + \frac{k^{-4}}{16} \int d^4x \sqrt{g} \left( v_2 [F_{\mu\nu}^a \tilde{F}^{\mu\nu, a}]^2 + t_2 (F_{\mu\nu}^a \tilde{F}^{\mu\nu, b})(F_{\rho\sigma}^a \tilde{F}^{\rho\sigma, b}) \right. \\ & \quad \left. + u_2 [F_{\mu\nu}^2 \tilde{F}^{\mu\nu, 2}]^2 + s_2 (F_{\mu\nu}^1 \tilde{F}^{\mu\nu, 2})(F_{\rho\sigma}^1 \tilde{F}^{\rho\sigma, 2}) \right)\end{aligned}$$

# Two species of gauge fields

- Gravity is "blind" to internal symmetries: induced interactions for Abelian and non-Abelian gauge fields are the same (at the fixed point)
- Interaction structure (involving four gauge fields and derivatives):

$$\begin{aligned}\Gamma_k^{U(1)\times U(1)} &= \frac{1}{4} \int d^4x \sqrt{g} F_{\mu\nu}^a F^{\mu\nu, a} + S_{\text{gf}, A} \\ &+ \frac{k^{-4}}{16} \int d^4x \sqrt{g} \left( w_2 [F_{\mu\nu}^a F^{\mu\nu, a}]^2 + y_2 (F_{\mu\nu}^a F^{\mu\nu, b})(F_{\rho\sigma}^a F^{\rho\sigma, b}) \right. \\ &\quad \left. + x_2 [F_{\mu\nu}^2 F^{\mu\nu, 2}]^2 + z_2 (F_{\mu\nu}^1 F^{\mu\nu, 2})(F_{\rho\sigma}^1 F^{\rho\sigma, 2}) \right) \\ &+ \frac{k^{-4}}{16} \int d^4x \sqrt{g} \left( v_2 [F_{\mu\nu}^a \tilde{F}^{\mu\nu, a}]^2 + t_2 (F_{\mu\nu}^a \tilde{F}^{\mu\nu, b})(F_{\rho\sigma}^a \tilde{F}^{\rho\sigma, b}) \right. \\ &\quad \left. + u_2 [F_{\mu\nu}^2 \tilde{F}^{\mu\nu, 2}]^2 + s_2 (F_{\mu\nu}^1 \tilde{F}^{\mu\nu, 2})(F_{\rho\sigma}^1 \tilde{F}^{\rho\sigma, 2}) \right)\end{aligned}$$

- Kinetic term: respects global  $O(2)$ -symmetry

# Two species of gauge fields

- Gravity is "blind" to internal symmetries: induced interactions for Abelian and non-Abelian gauge fields are the same (at the fixed point)
- Interaction structure (involving four gauge fields and derivatives):

$$\begin{aligned}\Gamma_k^{U(1)\times U(1)} &= \frac{1}{4} \int d^4x \sqrt{g} F_{\mu\nu}^a F^{\mu\nu, a} + S_{\text{gf}, A} \\ &+ \frac{k^{-4}}{16} \int d^4x \sqrt{g} \left( w_2 [F_{\mu\nu}^a F^{\mu\nu, a}]^2 + y_2 (F_{\mu\nu}^a F^{\mu\nu, b})(F_{\rho\sigma}^a F^{\rho\sigma, b}) \right. \\ &\quad \left. + x_2 [F_{\mu\nu}^2 F^{\mu\nu, 2}]^2 + z_2 (F_{\mu\nu}^1 F^{\mu\nu, 2})(F_{\rho\sigma}^1 F^{\rho\sigma, 2}) \right) \\ &+ \frac{k^{-4}}{16} \int d^4x \sqrt{g} \left( v_2 [F_{\mu\nu}^a \tilde{F}^{\mu\nu, a}]^2 + t_2 (F_{\mu\nu}^a \tilde{F}^{\mu\nu, b})(F_{\rho\sigma}^a \tilde{F}^{\rho\sigma, b}) \right. \\ &\quad \left. + u_2 [F_{\mu\nu}^2 \tilde{F}^{\mu\nu, 2}]^2 + s_2 (F_{\mu\nu}^1 \tilde{F}^{\mu\nu, 2})(F_{\rho\sigma}^1 \tilde{F}^{\rho\sigma, 2}) \right)\end{aligned}$$

- Kinetic term: respects global  $O(2)$ -symmetry

# Two species of gauge fields

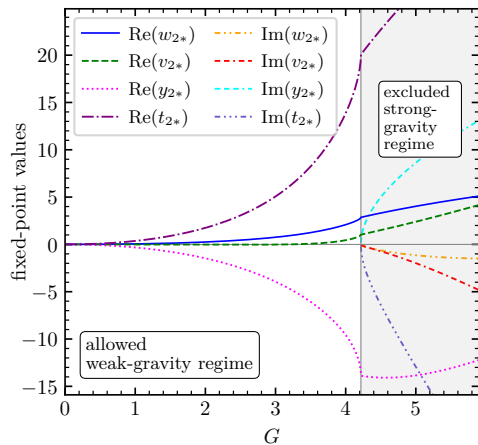
- Gravity is "blind" to internal symmetries: induced interactions for Abelian and non-Abelian gauge fields are the same (at the fixed point)
- Interaction structure (involving four gauge fields and derivatives):

$$\begin{aligned}\Gamma_k^{U(1)\times U(1)} = & \frac{1}{4} \int d^4x \sqrt{g} F_{\mu\nu}^a F^{\mu\nu, a} + S_{\text{gf}, A} \\ & + \frac{k^{-4}}{16} \int d^4x \sqrt{g} \left( w_2 [F_{\mu\nu}^a F^{\mu\nu, a}]^2 + y_2 (F_{\mu\nu}^a F^{\mu\nu, b})(F_{\rho\sigma}^a F^{\rho\sigma, b}) \right. \\ & \quad \left. + x_2 [F_{\mu\nu}^2 F^{\mu\nu, 2}]^2 + z_2 (F_{\mu\nu}^1 F^{\mu\nu, 2})(F_{\rho\sigma}^1 F^{\rho\sigma, 2}) \right) \\ & + \frac{k^{-4}}{16} \int d^4x \sqrt{g} \left( v_2 [F_{\mu\nu}^a \tilde{F}^{\mu\nu, a}]^2 + t_2 (F_{\mu\nu}^a \tilde{F}^{\mu\nu, b})(F_{\rho\sigma}^a \tilde{F}^{\rho\sigma, b}) \right. \\ & \quad \left. + u_2 [F_{\mu\nu}^2 \tilde{F}^{\mu\nu, 2}]^2 + s_2 (F_{\mu\nu}^1 \tilde{F}^{\mu\nu, 2})(F_{\rho\sigma}^1 \tilde{F}^{\rho\sigma, 2}) \right)\end{aligned}$$

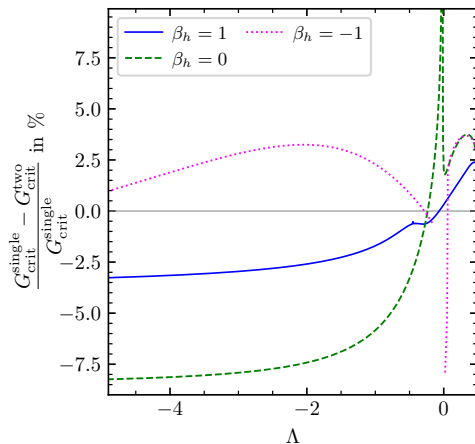
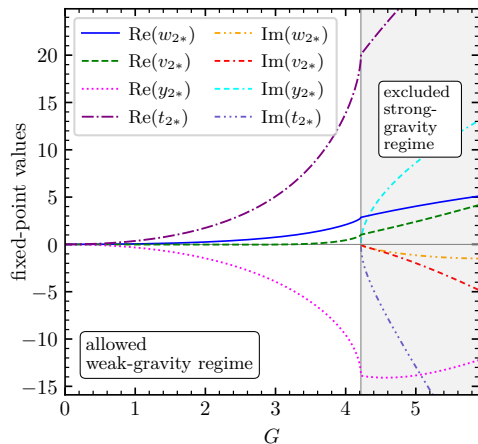
- Kinetic term: respects global  $O(2)$ -symmetry

$x_2, z_2, u_2, s_2$  break it  $\Rightarrow$  not induced see also [de Brito, Eichhorn, Robson Linos dos Santos; 2021]

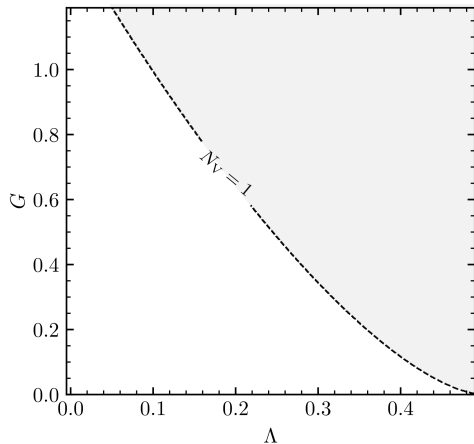
# Two gauge fields: Excluded strong-gravity regime



# Two gauge fields: Excluded strong-gravity regime

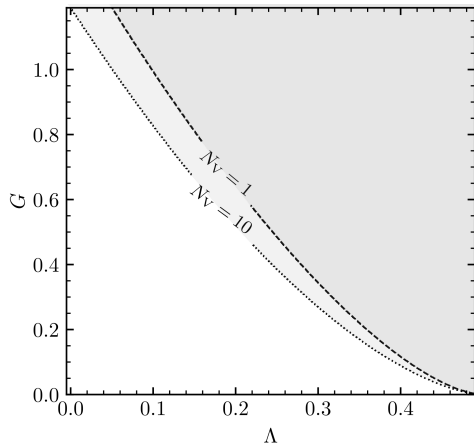


# From one to many gauge fields



- Investigate  $O(N_V)$ -symmetric system

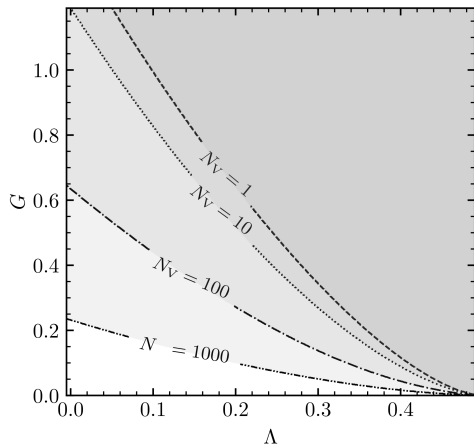
# From one to many gauge fields



- Investigate  $O(N_V)$ -symmetric system
- WGB becomes stronger for increasing  $O(N_V)$

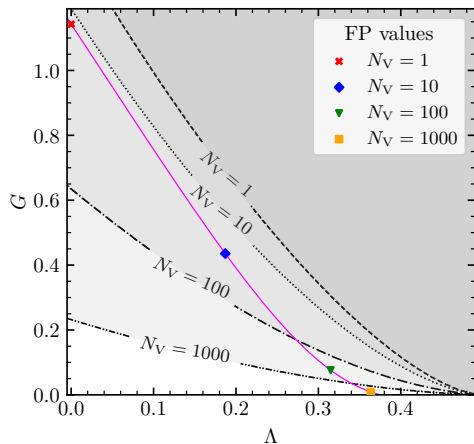


# From one to many gauge fields



- Investigate  $O(N_V)$ -symmetric system
- WGB becomes stronger for increasing  $O(N_V)$

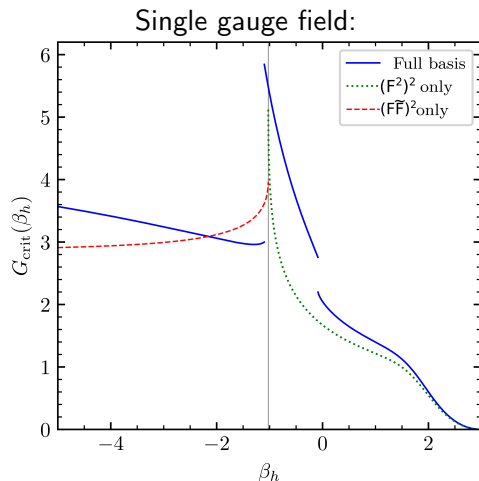
# From one to many gauge fields



- Investigate  $O(N_V)$ -symmetric system
- WGB becomes stronger for increasing  $O(N_V)$
- Compare to gravitational FP-values (background approximation):  
FP avoids WGB for all  $N_V$  (including  $N_V \rightarrow \infty$ )

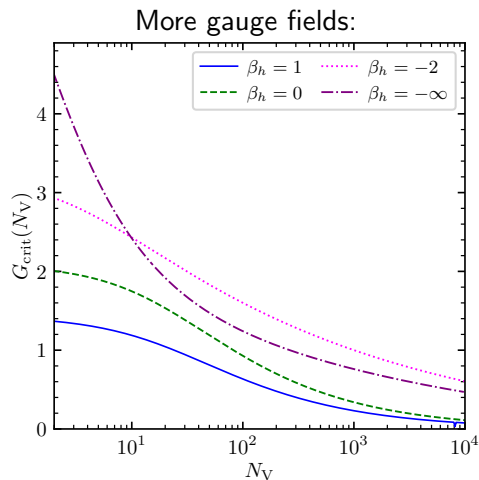
- Residual gauge dependence:  
estimate for robustness of results

# Gauge dependence of WGB



- Residual gauge dependence: estimate for robustness of results
- Qualitative behaviour is similar across different gauge-choices

# Gauge dependence of WGB



- Residual gauge dependence: estimate for robustness of results
- Qualitative behaviour is similar across different gauge-choices

# Summary + Outlook

---

# Summary + Outlook

- Summary
  - ▶ Interplay of gravity with matter:  
might put constraints on fundamental gravitational dynamics

# Summary + Outlook

- Summary
  - ▶ Interplay of gravity with matter:  
might put constraints on fundamental gravitational dynamics
  - ▶ Induced interactions:  
might limit strength of gravitational fluctuations



# Summary + Outlook

- Summary
  - ▶ Interplay of gravity with matter:  
might put constraints on fundamental gravitational dynamics
  - ▶ Induced interactions:  
might limit strength of gravitational fluctuations
  - ▶ WGB for gauge fields: becomes stronger for increasing  $N_V$ ;  
gravitational FP evades it.

# Summary + Outlook

- Summary
  - ▶ Interplay of gravity with matter:  
might put constraints on fundamental gravitational dynamics
  - ▶ Induced interactions:  
might limit strength of gravitational fluctuations
  - ▶ WGB for gauge fields: becomes stronger for increasing  $N_V$ ;  
gravitational FP evades it.
- Outlook
  - ▶ Investigate WGB for fermions up to (kinetic – term)<sup>2</sup> interactions

# Summary + Outlook

- Summary
  - ▶ Interplay of gravity with matter:  
might put constraints on fundamental gravitational dynamics
  - ▶ Induced interactions:  
might limit strength of gravitational fluctuations
  - ▶ WGB for gauge fields: becomes stronger for increasing  $N_V$ ;  
gravitational FP evades it.
- Outlook
  - ▶ Investigate WGB for fermions up to (kinetic – term)<sup>2</sup> interactions
  - ▶ Investigate WGB for scalars and gauge fields beyond (kinetic – term)<sup>2</sup> interactions

[de Brito, Knorr, MS; WIP ]

# Summary + Outlook

- Summary
  - ▶ Interplay of gravity with matter:  
might put constraints on fundamental gravitational dynamics
  - ▶ Induced interactions:  
might limit strength of gravitational fluctuations
  - ▶ WGB for gauge fields: becomes stronger for increasing  $N_V$ ;  
gravitational FP evades it.
- Outlook
  - ▶ Investigate WGB for fermions up to (kinetic – term)<sup>2</sup> interactions
  - ▶ Investigate WGB for scalars and gauge fields beyond (kinetic – term)<sup>2</sup> interactions  
[\[de Brito, Knorr, MS; WIP\]](#)
  - ▶ Investigate viability of pure-matter interacting FP [\[de Brito, Knorr, MS; WIP\]](#)  
see also [\[Laporte, Locht, Pereira, Saueressig; 2022\]](#)

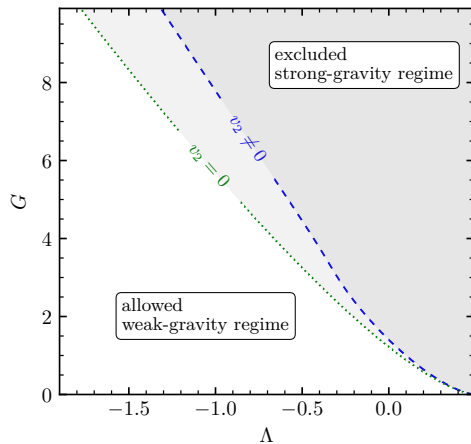
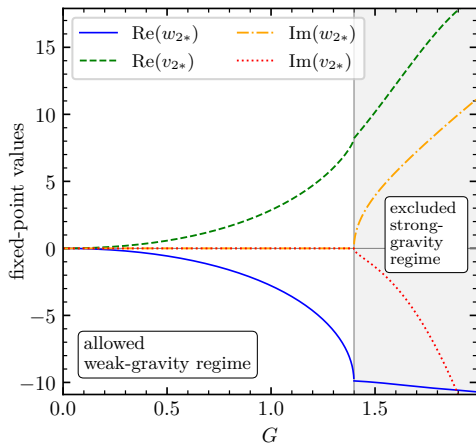
# Summary + Outlook

- Summary
  - ▶ Interplay of gravity with matter:  
might put constraints on fundamental gravitational dynamics
  - ▶ Induced interactions:  
might limit strength of gravitational fluctuations
  - ▶ WGB for gauge fields: becomes stronger for increasing  $N_V$ ;  
gravitational FP evades it.
- Outlook
  - ▶ Investigate WGB for fermions up to (kinetic – term)<sup>2</sup> interactions
  - ▶ Investigate WGB for scalars and gauge fields beyond (kinetic – term)<sup>2</sup> interactions  
[\[de Brito, Knorr, MS; WIP \]](#)
  - ▶ Investigate viability of pure-matter interacting FP [\[de Brito, Knorr, MS; WIP \]](#)  
see also [\[Laporte, Locht, Pereira, Saueressig; 2022\]](#)

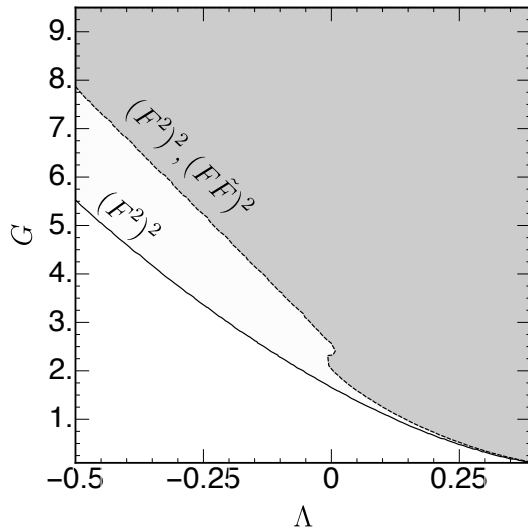
**Thank you for your attention!**

# One gauge field: Excluded strong-gravity regime

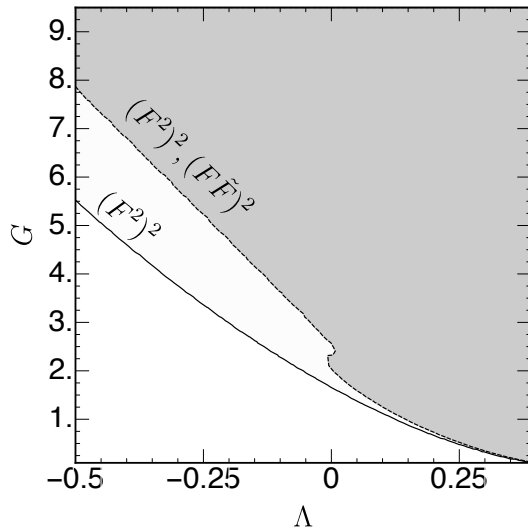
$$\Gamma_k^{U(1)} = \frac{1}{4} \int d^4x \sqrt{g} F_{\mu\nu} F^{\mu\nu} + S_{\text{gf}, A} + \frac{k^{-4}}{8} \int d^4x \sqrt{g} \left( w_2 (F_{\mu\nu} F^{\mu\nu})^2 + v_2 (F_{\mu\nu} \tilde{F}^{\mu\nu})^2 \right).$$



# Jump in $G_{\text{crit}}$



# Critical Dimension





## Solution to the triviality problem in $d > 4$

- $[\bar{g}_Y] = \frac{4-d}{2}$

$$\beta_{g_Y} = g_Y \left( \frac{d-4}{2} - f_g(d) \right) + \mathcal{O}(g_Y^3)$$

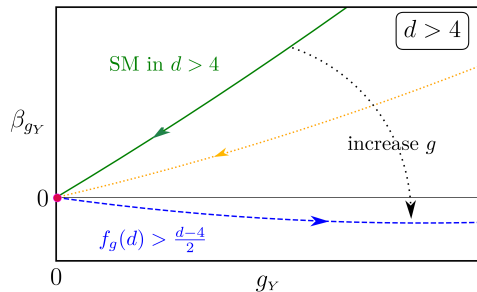
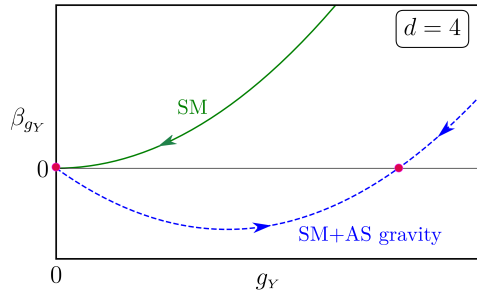
# Solution to the triviality problem in $d > 4$

- $[\bar{g}_Y] = \frac{4-d}{2}$

$$\beta_{g_Y} = g_Y \left( \frac{d-4}{2} - f_g(d) \right) + \mathcal{O}(g_Y^3)$$

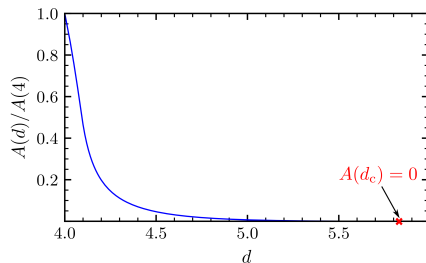
- Competition of  $f_g(d)$  with canonical mass term.
- Necessary condition for UV completion:  
Effective dimensionality below four,

$$f_g(d) > \frac{d-4}{2}$$



# UV complete matter sector beyond $d = 4$ ?

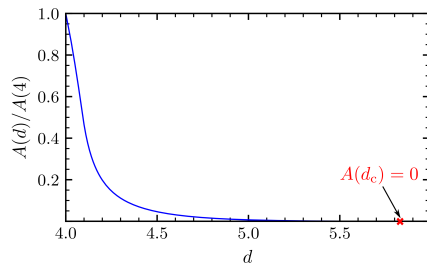
- Area of allowed region for  $G \in (0, 1000)$  and  $\Lambda \in (-1500, 0.5)$ .
- Area shrinks to zero at  $d_c \approx 5.8$ .



[Eichhorn, MS; 2019]

# UV complete matter sector beyond $d = 4$ ?

- Area of allowed region for  $G \in (0, 1000)$  and  $\Lambda \in (-1500, 0.5)$ .
- Area shrinks to zero at  $d_c \approx 5.8$ .
- Calculation leading to green and red area: Subject to systematic errors due to truncation.

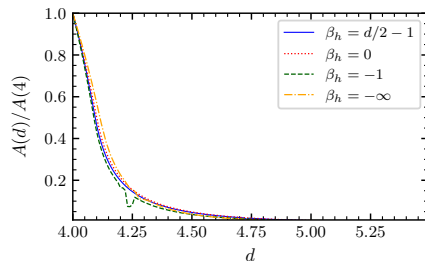


[Eichhorn, MS; 2019]

- Very large deformations necessary to make  $d \geq 6$  viable (in explored range).
- Qualitative aspects of the scenario remain unchanged.

# UV complete matter sector beyond $d = 4$ ?

- Area of allowed region for  $G \in (0, 1000)$  and  $\Lambda \in (-1500, 0.5)$ .
- Area shrinks to zero at  $d_c \approx 5.8$ .
- Calculation leading to green and red area: Subject to systematic errors due to truncation.

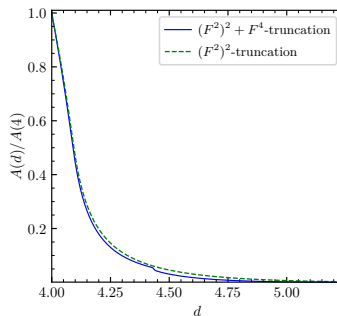


[MS; 2021]

- Very large deformations necessary to make  $d \geq 6$  viable (in explored range).
- Qualitative aspects of the scenario remain unchanged.

# UV complete matter sector beyond $d = 4$ ?

- Area of allowed region for  $G \in (0, 1000)$  and  $\Lambda \in (-1500, 0.5)$ .
- Area shrinks to zero at  $d_c \approx 5.8$ .
- Calculation leading to green and red area: Subject to systematic errors due to truncation.

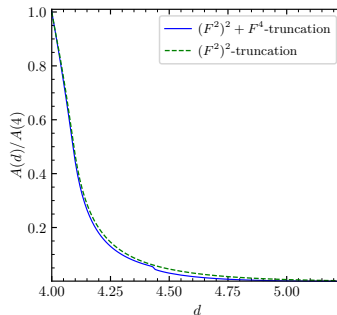


[Eichhorn, Kwapisz, MS; 2021]

- Very large deformations necessary to make  $d \geq 6$  viable (in explored range).
- Qualitative aspects of the scenario remain unchanged.

# UV complete matter sector beyond $d = 4$ ?

- Area of allowed region for  $G \in (0, 1000)$  and  $\Lambda \in (-1500, 0.5)$ .
- Area shrinks to zero at  $d_c \approx 5.8$ .
- Calculation leading to green and red area: Subject to systematic errors due to truncation.



[Eichhorn, Kwapisz, MS; 2021]

- Very large deformations necessary to make  $d \geq 6$  viable (in explored range).
- Qualitative aspects of the scenario remain unchanged.

The predictive power of the asymptotic-safety paradigm could extend to fundamental parameters of the geometry.