# Cosmological constant problem and Hubble tension in scale-dependent cosmology

Cristóbal Laporte with <u>P. Alvarez , B. Koch & A. Rincón</u>

Based on:

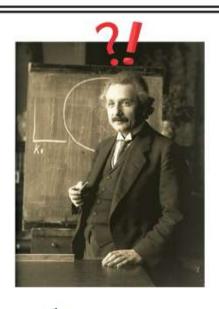
JCAP no 1, 21, 2020, arXiv:1812.10526 JCAP no 6, 019, 2021, arXiv: 2009.02311

arXiv:2205.05592

#### Content

- Cosmological constant problem: status
- Scale-dependent framework and evolving Universe
- Possible solution: Deflation during inflation
- Link to Asymptotic Safety
- # Hubble tension under the light of SD gravity
- Conclusions

Radboud U. '21







$$R_{\mu\nu} - \frac{1}{2}Rg_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu} \qquad R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} + Q_{\mu\nu} = 8\pi G T_{\mu\nu} \qquad \qquad R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

 $\Rightarrow \dot{a} > 0$  not static

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3p) + \frac{\Lambda}{3} \equiv 0 \qquad \qquad \frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3p) + \frac{\Lambda}{3} \equiv 0$$

$$\Rightarrow \dot{a} \equiv 0 \quad \text{static} \qquad \Rightarrow \dot{a} > 0 \quad \text{not static}$$

$$\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3p) + \frac{\Lambda}{3} \equiv 0$$
$$\Lambda > 0 \Rightarrow \ddot{a} > 0$$

1917

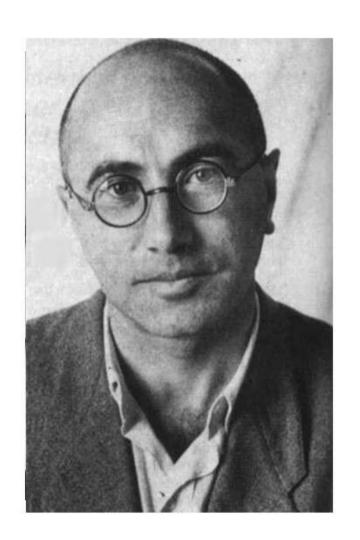
1929

1998

Yakov Zeldovich

Quantum fluctuation predict value of  $\Lambda$ 

1967



Steven Weinberg

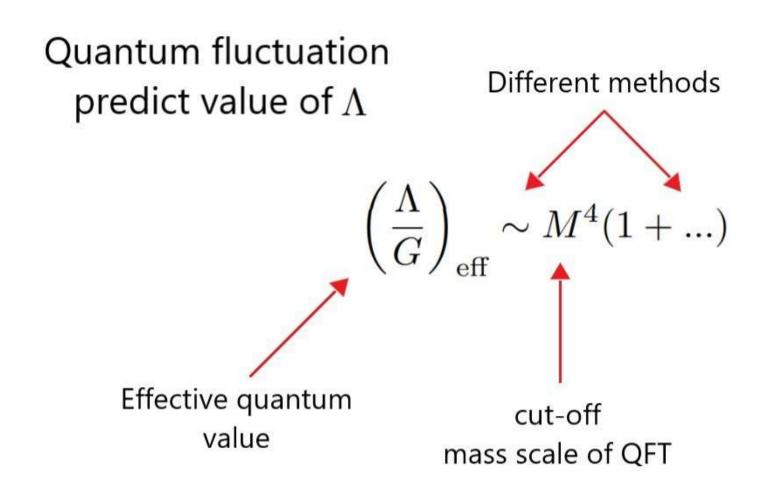
Quantum fluctuation predict value of  $\Lambda$ 

Problem since '98

ref[3]: S. Wienberg '98

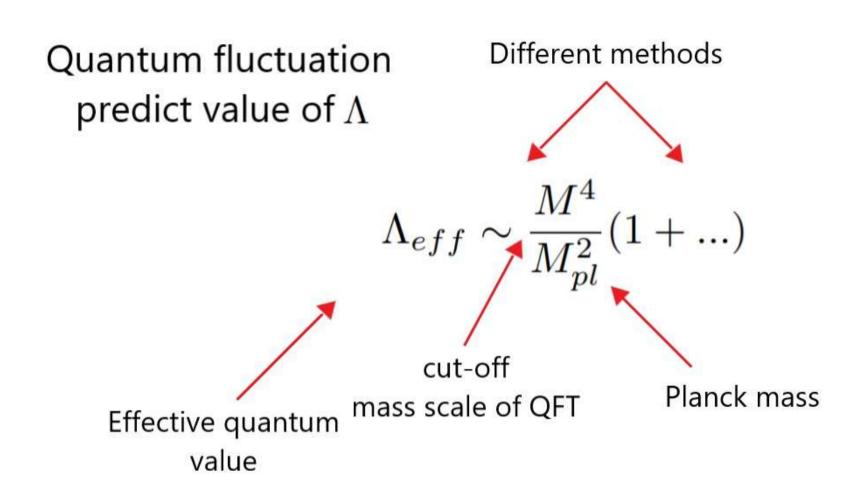


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Use

$$G \sim rac{1}{M_{pl}^2}$$



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Quantum fluctuation predict value of  $\Lambda$ 

Highest physical

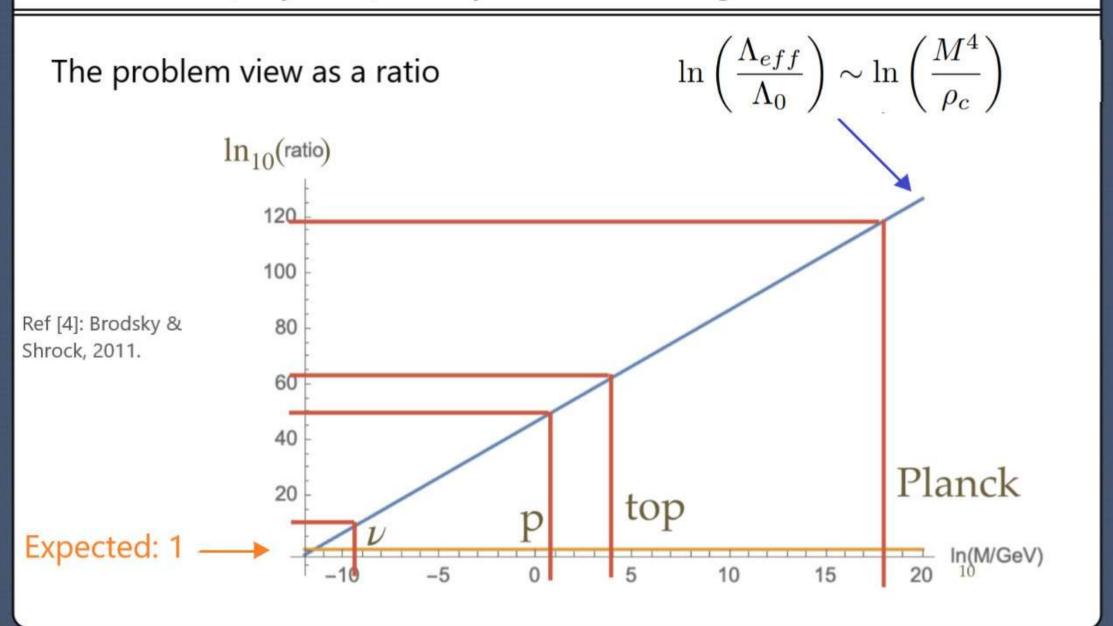
mass scale

$$\Lambda_{eff} \sim \frac{M^4}{M_{pl}^2} (1 + \dots)$$

Observed value

$$\Lambda_0 = \frac{\rho_c}{M_{pl}^2} \approx \frac{10^{-47} \text{GeV}^4}{M_{pl}^2}$$

 $\rho_c$  = observed critical energy density



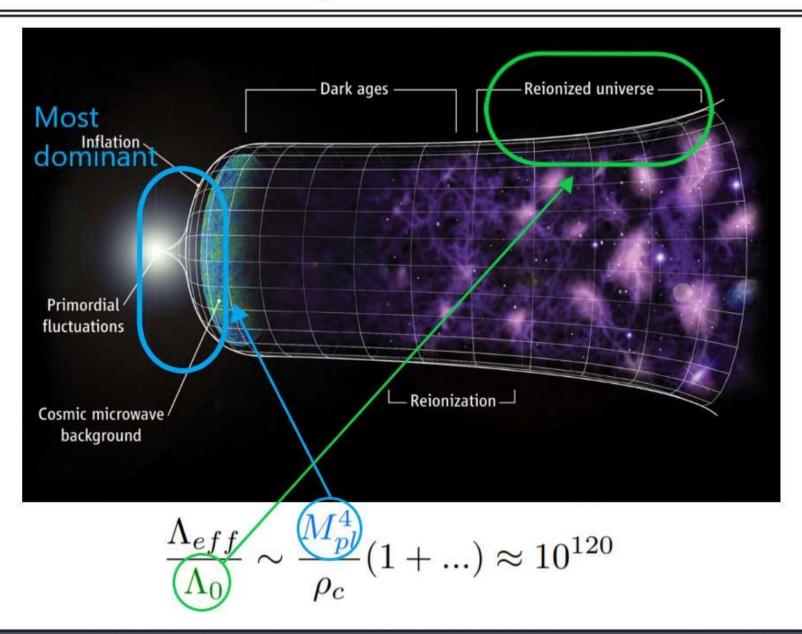
The problem view as a ratio

$$\frac{\Lambda_{eff}}{\Lambda_0} \sim \frac{1}{G_N \Lambda_0} \sim \frac{M_{pl}^4}{\rho_c} pprox 10^{120}$$

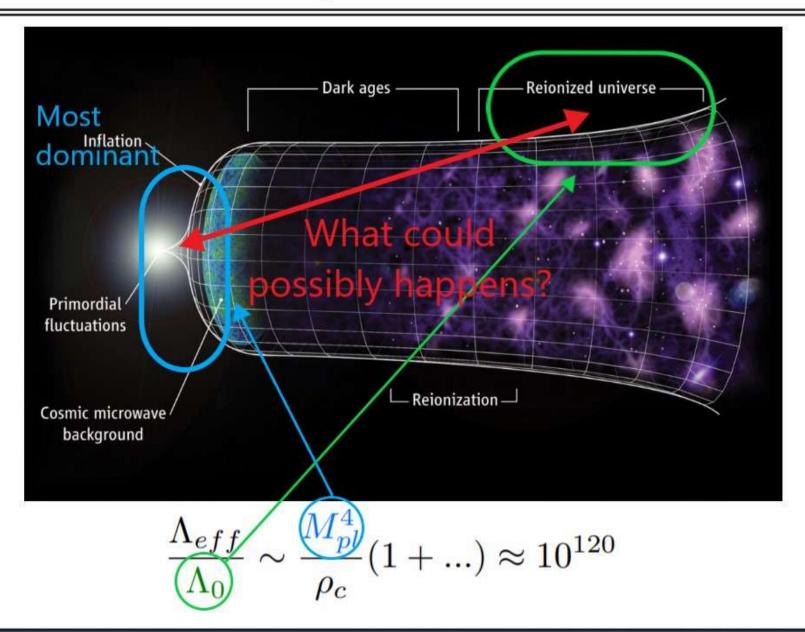
We try to address this problem

Assuming there are quantum fluctuations of gravity associated to the Planck scale

# **Evolving Universe Issue**



# **Evolving Universe Issue**



# Scale-dependent Framework

Gravity as classical theory

$$S = \frac{1}{16\pi} \int d^4x \sqrt{-g} \left( \frac{R}{G_N} - 2\frac{\Lambda_0}{G_N} \right)$$

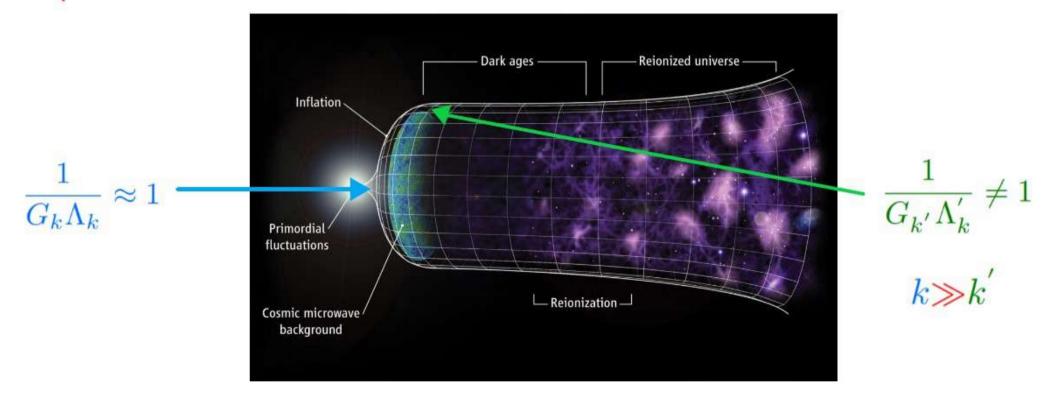
Gravity as an effective QFT

$$\Gamma_{\mathbf{k}} = \frac{1}{16\pi} \int d^4x \sqrt{-g} \left( \frac{R}{G_{\mathbf{k}}} - 2 \frac{\Lambda_{\mathbf{k}}}{G_{\mathbf{k}}} \right) + \dots$$

# Scale Dependent Framework

$$\Gamma_{\mathbf{k}} = \frac{1}{16\pi} \int d^4x \sqrt{-g} \left( \frac{R}{G_{\mathbf{k}}} - 2 \frac{\Lambda_{\mathbf{k}}}{G_{\mathbf{k}}} \right) + \dots$$

### Implications for CCP?



$$\Gamma_{\mathbf{k}} = \frac{1}{16\pi} \int d^4x \sqrt{-g} \left( \frac{R}{G_{\mathbf{k}}} - 2 \frac{\Lambda_{\mathbf{k}}}{G_{\mathbf{k}}} \right)$$

We need to solve the gap equations

$$G_{\mu\nu} = -\Lambda_k g_{\mu\nu} - \Delta t_{\mu\nu}$$

with

$$\Delta t_{\mu\nu} = G_k \left( g_{\nu} \nabla^{\alpha} \nabla_{\alpha} - \nabla_{\mu} \nabla_{\nu} \right) G_k^{-1}$$

Gap equations

$$\left(\frac{\dot{a}}{a}\right)^2 + \frac{\kappa}{a^2} - \frac{\Lambda_k}{3} = \underbrace{\frac{\rho_{SD}}{3}}_{\text{Scale-dependent}}$$
 
$$2\frac{\ddot{a}}{a} + \left(\frac{\dot{a}}{a}\right)^2 + \frac{\kappa}{a^2} - \Lambda_k = -\underbrace{\rho_{SD}}_{\text{SD}}$$
 quantities

Since 
$$k = k(t) \Rightarrow G_k = G(t)$$
 &  $\Lambda_k = \Lambda(t)$ 

$$\frac{1}{3}\rho_{SD} = \left(\frac{\dot{a}}{a}\right) \left(\frac{\dot{G}}{G}\right) \quad \& \quad -\left(p_{SD}\right) = -2\left(\frac{\dot{G}}{G}\right)^2 + \frac{\ddot{G}}{G} + 2\left(\frac{\dot{G}}{G}\right) \left(\frac{\dot{a}}{a}\right)$$

Gap equations

$$\left(\frac{\dot{a}}{a}\right)^2 + \frac{\kappa}{a^2} - \frac{\Lambda(t)}{3} = \left(\frac{\dot{G}}{G}\right) \left(\frac{\dot{a}}{a}\right)$$

$$2\frac{\ddot{a}}{a} + \left(\frac{\dot{a}}{a}\right)^2 + \frac{\kappa}{a^2} - \Lambda(t) = -2\left(\frac{\dot{G}}{G}\right)^2 + \frac{\ddot{G}}{G} + 2\left(\frac{\dot{G}}{G}\right)\left(\frac{\dot{a}}{a}\right)$$

Problem: 2 Equations

3 Unknown functions  $a(t) G(t) \& \Lambda(t)$ 

Solution: Impose an energy condition

Null Energy Condition (NEC)

$$\Delta t_{\mu\nu}\ell^{\mu}\ell^{\nu} = 0$$

where

$$\frac{\mathrm{d}\ell^{\mu}}{\mathrm{d}t} + \Gamma^{\mu}_{\alpha\beta} = 0 \Rightarrow \ell^{\mu} = \frac{c_0}{a} \left( 1, \frac{(1 - \kappa r^2)^{-\frac{1}{2}}}{a}, 0, 0 \right)$$

thus

$$-2\left(\frac{\dot{G}}{G}\right)^{2} + \left(\frac{\dot{G}}{G}\right) - \left(\frac{\dot{G}}{G}\right)\left(\frac{\dot{a}}{a}\right) = 0$$

Gap equations

$$\left(\frac{\dot{a}}{a}\right)^2 + \frac{\kappa}{a^2} - \frac{\Lambda(t)}{3} = \left(\frac{\dot{G}}{G}\right) \left(\frac{\dot{a}}{a}\right)$$

$$2\frac{\ddot{a}}{a} + \left(\frac{\dot{a}}{a}\right)^2 + \frac{\kappa}{a^2} - \Lambda(t) = -2\left(\frac{\dot{G}}{G}\right)^2 + \frac{\ddot{G}}{G} + 2\left(\frac{\dot{G}}{G}\right)\left(\frac{\dot{a}}{a}\right)$$

**NEC** 

$$-2\left(\frac{\dot{G}}{G}\right)^{2} + \left(\frac{\dot{G}}{G}\right) - \left(\frac{\dot{G}}{G}\right)\left(\frac{\dot{a}}{a}\right) = 0$$

3 equations, 3 unknowns



Solution

$$a(t) = a_i \exp\left(\frac{t}{\sqrt{\Lambda_0/3}}\right)$$

$$G(t) = \frac{G_0}{1 + \xi a(t)}$$

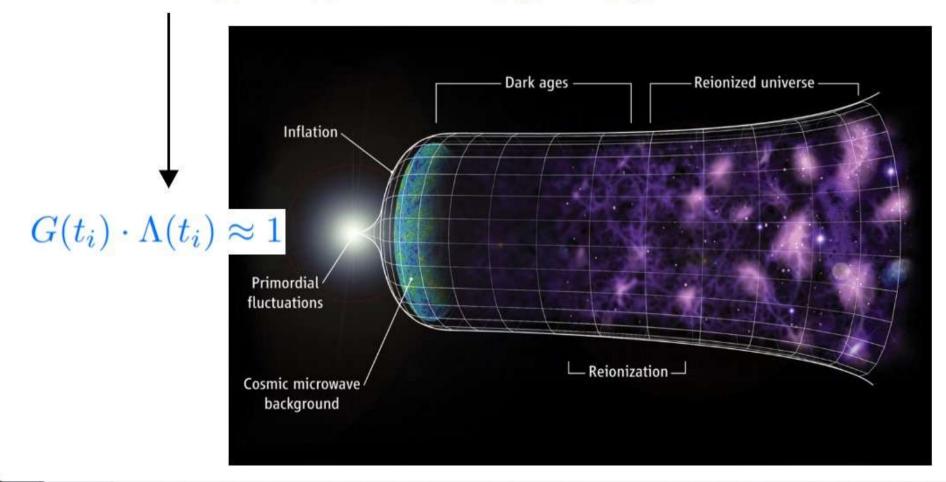
$$\Lambda(t) = \Lambda_0 \left(\frac{1 + 2\xi a(t)}{1 + \xi a(t)}\right)$$

Three integration constants:  $G_0, \Lambda_0(\xi)$ 

What does it mean for the CCP?

#### What does this means for the CCP?

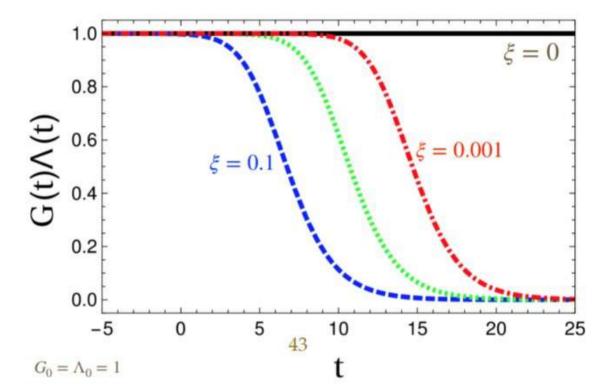
$$G_k \cdot \Lambda_k = G(t) \cdot \Lambda(t) \longrightarrow G(t_f) \cdot \Lambda(t_f) \ll 1$$



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#### What does this means for the CCP?

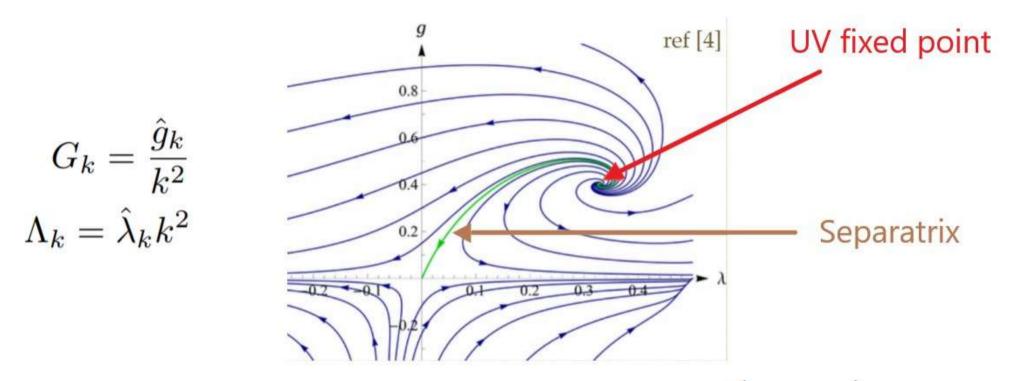
$$G(t) \cdot \Lambda(t) = \frac{G_0}{1 + \xi a(t)} \cdot \Lambda_0 \left( \frac{1 + 2\xi a(t)}{1 + \xi a(t)} \right)$$



Nice, link to AS?

### Link to AS?

#### Remember:

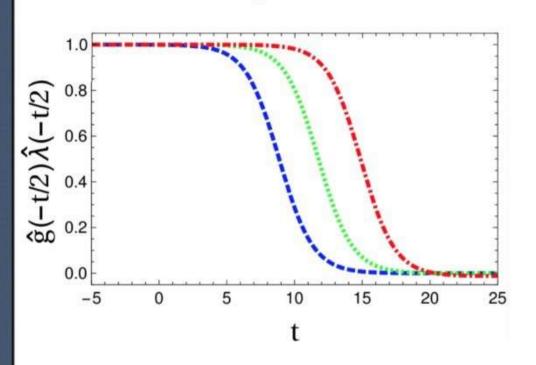


$$\hat{g}(\hat{t}) = \frac{g_0 e^{2\hat{t}}}{1 + g_0(e^{2\hat{t}})/g^*} \quad \hat{\lambda}(\hat{t}) = \frac{g^* \lambda_0 + e^{-2\hat{t}} \left(e^{4\hat{t}} - 1\right) g_0 \lambda^*}{1 + g_0 \left(e^{2\hat{t}} - 1\right)/g^*}$$

### Link to AS?

AS renormalization flow

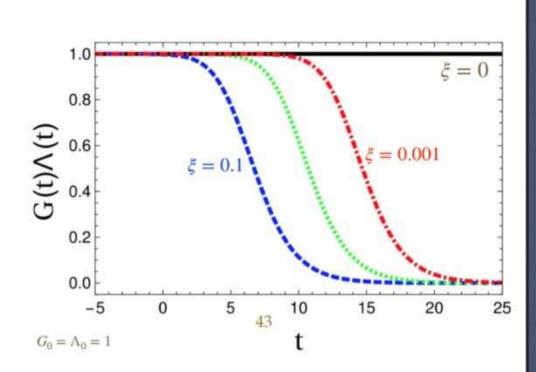
$$G_k \cdot \Lambda_k = \frac{\hat{g}_k}{k^2} k^2 \hat{\lambda}_k = \hat{g}_k \cdot \hat{\lambda}_k$$



VS

SD & NEC

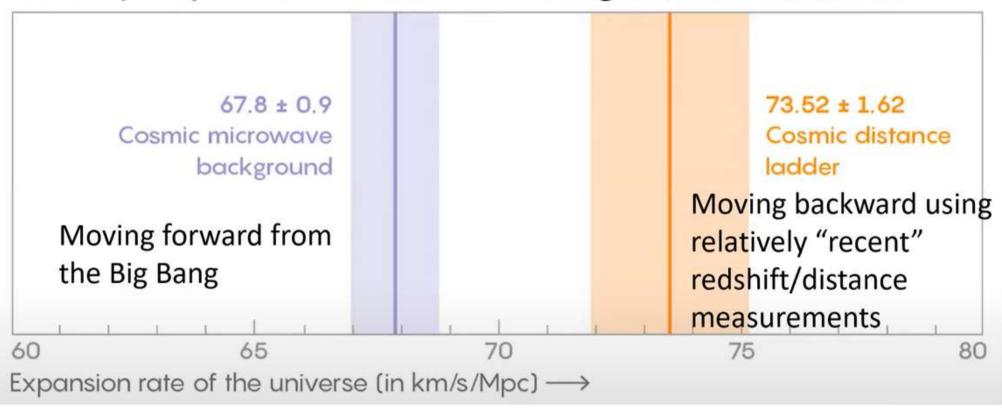
$$G(t) \cdot \Lambda(t)$$



Looks familiar!

### The Hubble constant problem

#### A Discrepancy in the Hubble Constant Using two different Methods



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# Can SD cosmology alleviate the H0 tension?

Time



Redshift

Tension:

$$\Delta(z) \equiv \frac{H(z)}{H^{(\Lambda {
m CDM})}(z)} - 1,$$

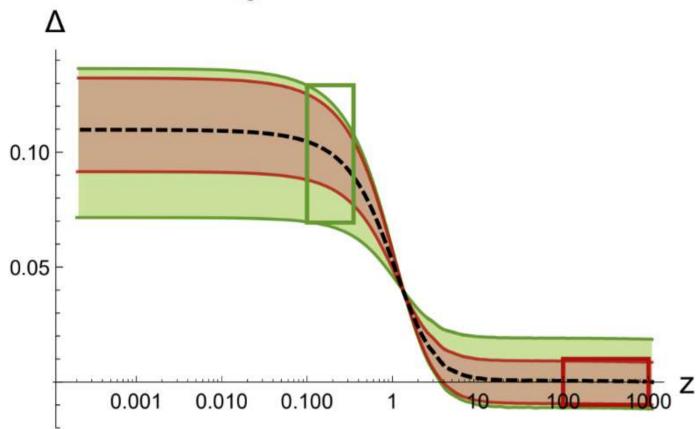
Conflict between CMB and SN Ia

$$\Delta = \frac{H_0^{\text{(late)}}}{H_0^{\text{(early)}}} - 1 \approx 0.09 \pm 0.02$$
.

# Can SD cosmology alleviate the H0 tension?

#### Conflict between CMB and SN Ia

$$\Delta = \frac{H_0^{({
m late})}}{H_0^{({
m early})}} - 1 \approx 0.09 \pm 0.02$$
.



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# Concluding comments

- CCP: watch out SD during inflation
- Showed: with SD CCP is alleviated
- Beautiful matching between AS & SD
- SD scenario of gravity could offer an alleviation to the conflict between early and late-time measurements

### Backup

$$\hat{g}(\hat{t}) \cdot \hat{\lambda}(\hat{t}) = \frac{g_0 e^{2\hat{t}}}{1 + g_0(e^{2\hat{t}})/g^*} \cdot \frac{g^* \lambda_0 + e^{-2\hat{t}} \left(e^{4\hat{t}} - 1\right) g_0 \lambda^*}{1 + g_0 \left(e^{2\hat{t}} - 1\right)/g^*}$$

Approximate to UV FP and separatrix

$$\hat{g}(\hat{t}) \cdot \hat{\lambda}(\hat{t}) = g^* \lambda^* \left( \frac{g^* \lambda_0}{g_0 \lambda^*} + e^{2\hat{t}} \right) \left( e^{2\hat{t}} + \frac{g^*}{g_0} \right)^{-2} \equiv G(t) \cdot \Lambda(t)$$

For

$$g^* \lambda^* \to G_0 \Lambda_0$$
  
 $g_0 \to \frac{G_0}{a_i \xi}$   
 $\hat{t} \to -\frac{t}{2\tau}$ 

### Backup

### Modified Friedmann equations

$$\begin{split} \frac{1}{H_0^2}\left(H^2-H\frac{\dot{g}}{g}\right) &= \Omega_\Lambda\lambda(t) + \frac{\Omega_r}{a^4}g + \frac{\Omega_m}{a^3}g \\ \frac{1}{H_0^2}\left(2\dot{H}+3H^2-2H\frac{\dot{g}}{g}+2\frac{\dot{g}^2}{g^2}-\frac{\ddot{g}}{g}\right) &= 3\Omega_\Lambda\lambda(t) - \frac{\Omega_r}{a^4}g \\ \frac{\ddot{g}}{g}-H\frac{\dot{g}}{g}-2\frac{\dot{g}^2}{g^2} &= 0\,. \end{split}$$

Initial conditions

$$a(t_0) = 1$$
,  $\dot{a}(t_0) = H_0$ ,  
 $g(t_0) = 1$ ,  $\dot{g}(t_0) \sim 0.1$ ,



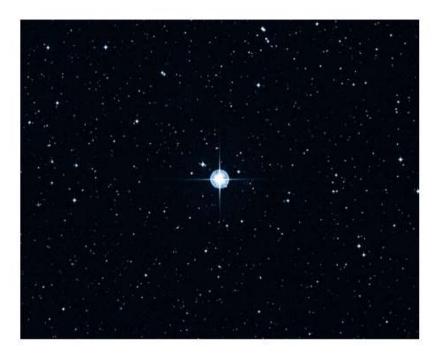
- Expansion rate of the Universe at present as implied by late time observations



**\*\*** Numerical constraints

- Expansion rate at  $z\sim 10^3$  as it is inferred from CMB anisotropies





- Age of the Universe coming from stellar evolution

Bond et all, ApJ letters '13



Dimensionless function

$$x = \frac{a(t)}{a_0},$$

System of 2 FE + NEC

$$\frac{x'(\tau)^2}{x(\tau)^2} - \frac{g'(\tau)x'(\tau)}{g(\tau)x(\tau)} = \Omega_m g(\tau)x(\tau)^{-3} + \Omega_r g(\tau)x(\tau)^{-4} + \Omega_\Lambda \lambda(\tau),$$

$$\frac{x'(\tau)^2}{x(\tau)^2} - 2\frac{g'(\tau)x'(\tau)}{g(\tau)x(\tau)} + 2\frac{g'(\tau)^2}{g(\tau)^2} + 2\frac{x''(\tau)}{x(\tau)} - \frac{g''(\tau)}{g(\tau)} = -\Omega_r g(\tau)x(\tau)^{-4} + 3\Omega_\Lambda \lambda(\tau),$$

$$\frac{g''(\tau)}{g(\tau)} - \frac{g'(\tau)x'(\tau)}{g(\tau)x(\tau)} - 2\frac{g'(\tau)^2}{g(\tau)^2} = 0.$$

 Expansion rate of the Universe at present as implied by late time observations

$$\Sigma_{(i)} = \frac{\text{abs}\left[\frac{\dot{x}/x}{\dot{x}^{(\Lambda\text{CDM})}/x^{(\Lambda\text{CDM})}}\Big|_{z\sim 0} - \frac{H_0^{(\text{late})}}{H_0^{(\text{early})}}\right]}{\frac{H_0^{(\text{late})}}{H_0^{(\text{early})}}\left(\left(\frac{\Delta H_0^{(\text{late})}}{H_0^{(\text{late})}}\right)^2 + \left(\frac{\Delta H_0^{(\text{early})}}{H_0^{(\text{early})}}\right)^2\right)^{1/2}}$$

- Expansion rate at  $z\sim 10^3$  as it is inferred from CMB anisotropies

$$\Sigma_{(ii)} = \frac{\text{abs}\left[\frac{\dot{x}/x}{\dot{x}^{(\Lambda\text{CDM})}/x^{(\Lambda\text{CDM})}}\Big|_{z \sim 10^3} - 1\right]}{\Delta H_0^{(\text{early})}/H_0^{(\text{early})}}$$

- Age of the Universe coming from stellar evolution

$$\Sigma_{(iii)} = \frac{\text{abs}\left[t_{\text{age}} - t_{\text{age}}^*\right]}{\Delta t_{\text{age}}^*}.$$

