

Multiloop flow equations for single-boson exchange fRG

Marcel Gievers

MPQ (Garching) / LMU (Munich)

$$\dot{w}_a^{(\ell+2)} = \text{wavy line} \rightarrow \dot{\gamma}_{\bar{a}}^{(\ell)} \rightarrow \text{wavy line}$$
$$\dot{\lambda}_a^{(\ell+2)} = \dot{\gamma}_{\bar{a}}^{(\ell+1)} + T_a \rightarrow \dot{\gamma}_{\bar{a}}^{(\ell)} \rightarrow T_a$$
$$\dot{M}_a^{(\ell+2)} = \dot{\gamma}_{\bar{a}}^{(\ell+1)} \rightarrow T_a + T_a \rightarrow \dot{\gamma}_{\bar{a}}^{(\ell)} \rightarrow T_a + T_a \rightarrow \dot{\gamma}_{\bar{a}}^{(\ell+1)}$$

Multiloop flow equations for single-boson exchange fRG

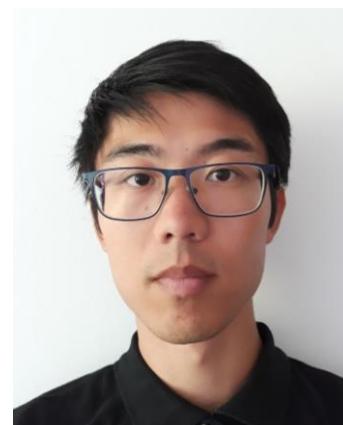
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European Physics Journal B **95** 108 (2022)



Elias Walter
LMU (Munich)



Anxiang Ge
LMU (Munich)



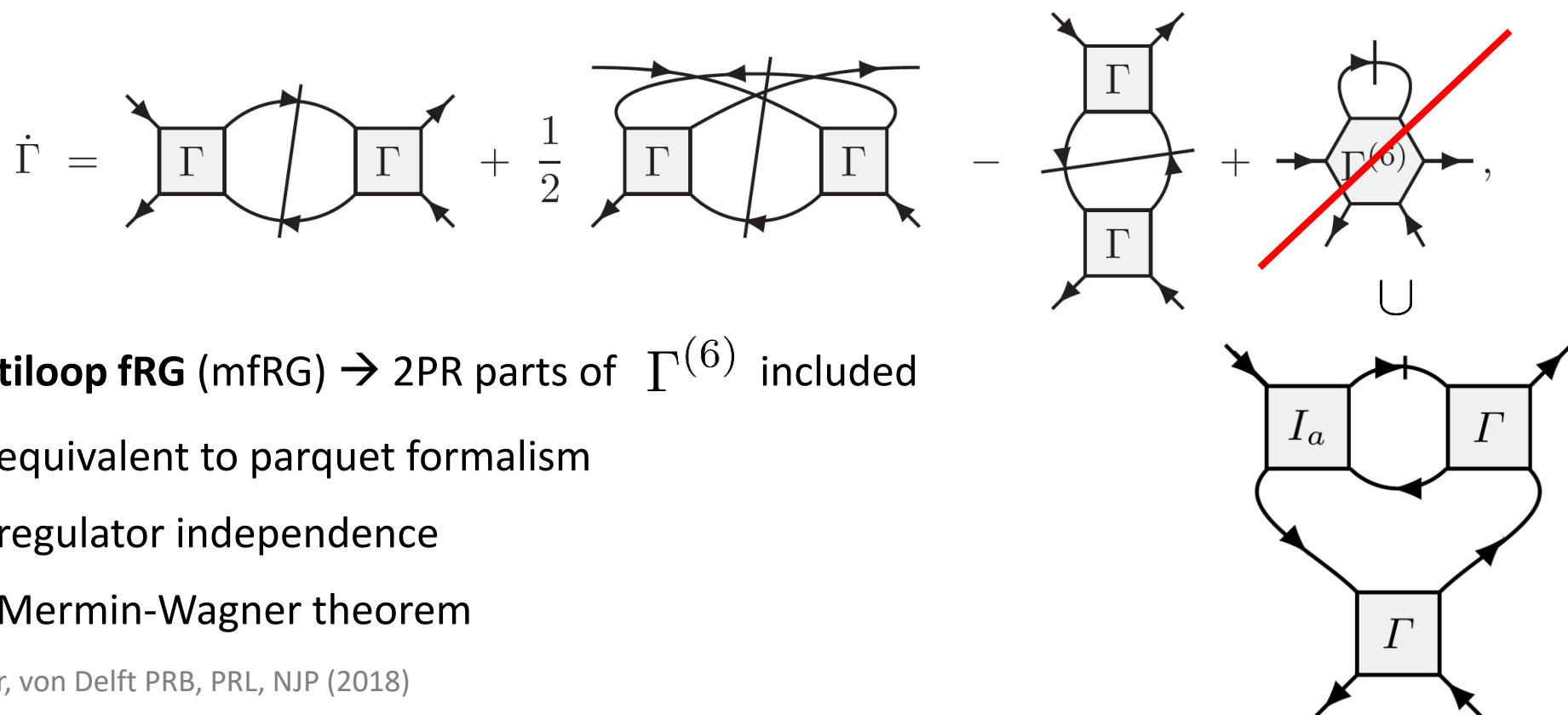
Jan von Delft
LMU (Munich)



Fabian Kugler
Rutgers U. (NJ, USA)

Background

- **Wetterich equation** → exact formulation, but vertex expansion needed
- **one-loop fRG** → 6-point vertex $\Gamma^{(6)}$ not included

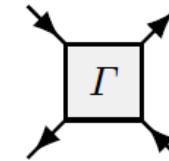


- **multiloop fRG (mfRG)** → 2PR parts of $\Gamma^{(6)}$ included
 - equivalent to parquet formalism
 - regulator independence
 - Mermin-Wagner theorem

Kugler, von Delft PRB, PRL, NJP (2018)

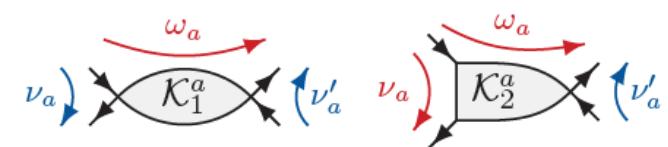
Motivation for single-boson exchange (SBE)

- problem: high-complexity of 4-point vertex $\Gamma = \Gamma_{q,k,k'}(\omega, \nu, \nu')$



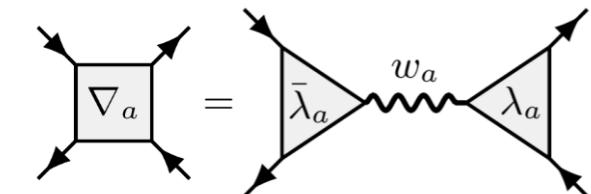
→ parametrization in terms of its **asymptotics**

Wentzell, Li, Tagliavini, Taranto, Rohringer, Held, Toschi, Andergassen
arXiv (2016), PRB (2020)



Hubbard-like models {

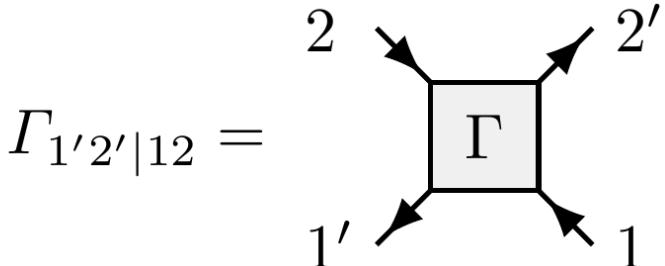
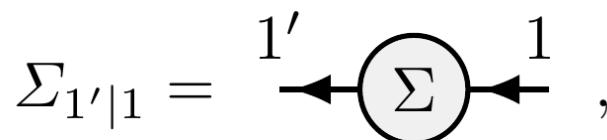
- vertex decomposition in **single-boson exchange (SBE) processes**
Krien, Valli, Capone, Lichtenstein, Rohringer, Held, Kauch, Harkov
PRBs (2019, 2020, 2021)
- **1l fRG in SBE decomposition in Hubbard model**
Bonetti, Toschi, Hille, Andergassen, Vilardi PRR (2022)



→ general formulation of SBE, **SBE mfRG** flow equations

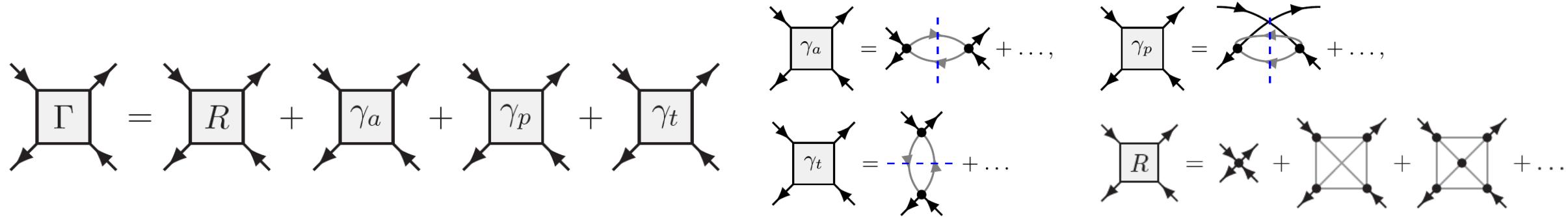
Setting: Fermionic theory

- general fermionic action: $S = -\bar{c}_{1'}[G_0^{-1}]_{1'}|_1 c_1 - \frac{1}{4}U_{1'2'}|_{12}\bar{c}_{1'}\bar{c}_{2'}c_2c_1$
- constant energy-conserving bare interaction U 
- 1- & 2-particle correlation functions: $G_{1|1'} = -\langle c_1\bar{c}_{1'} \rangle$ $G_{12|1'2'}^{(4)} = \langle c_1c_2\bar{c}_{2'}\bar{c}_{1'} \rangle$
- self-energy Σ and 4-point vertex Γ

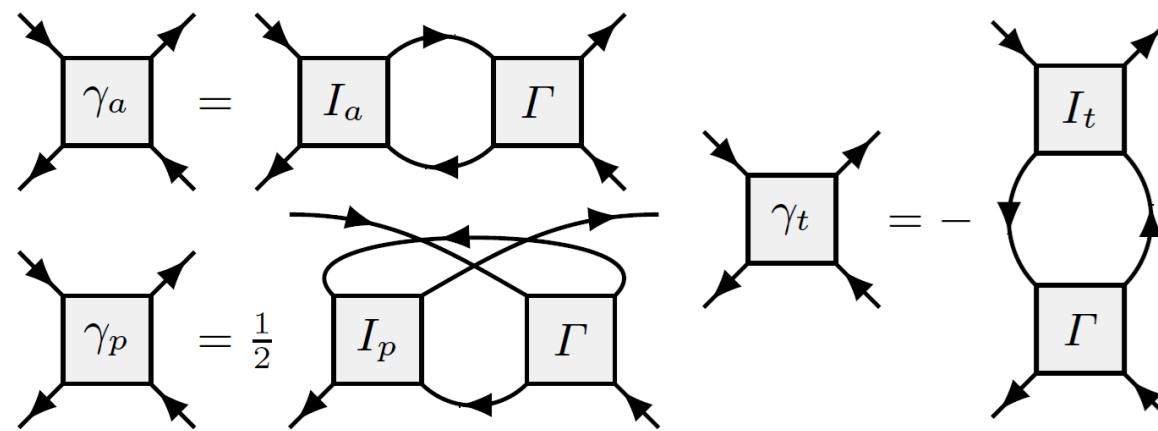


Parquet decomposition

- full 1PI vertex Γ as sum of 2PR vertices γ_r



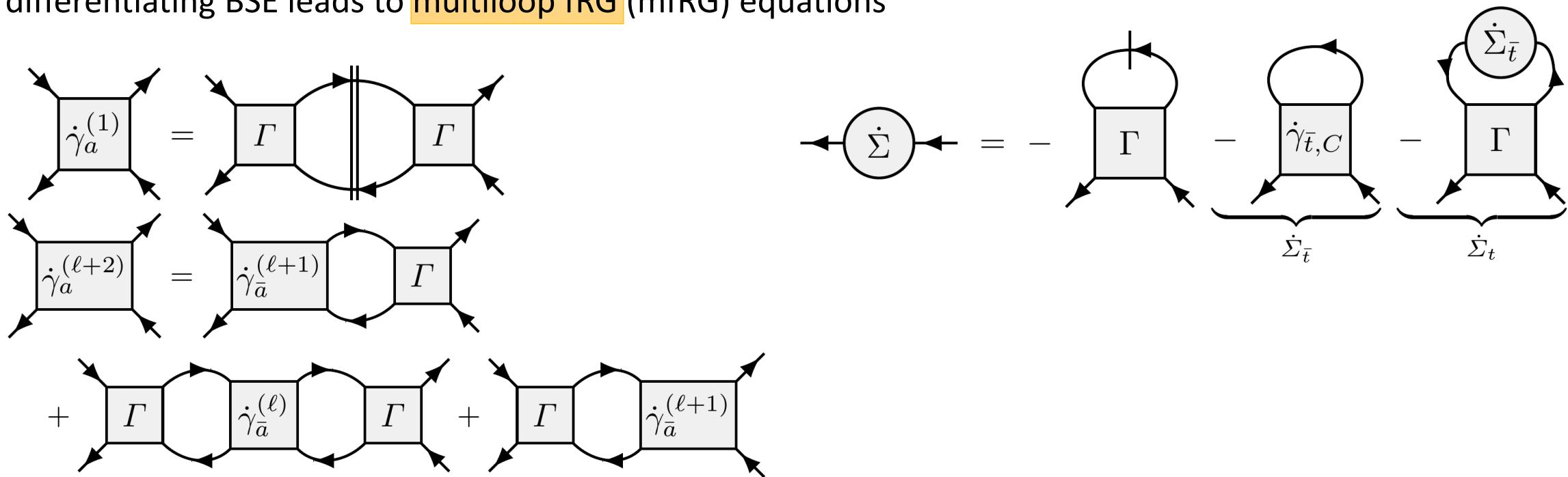
- 2PR vertices γ_r are built from Bethe-Salpeter equations (BSE) $\gamma_r = I_r \circ \Pi_r \circ \Gamma$



antiparallel ($\overline{\text{ph}}$)
parallel (pp)
transversal (ph)

Multiloop fRG

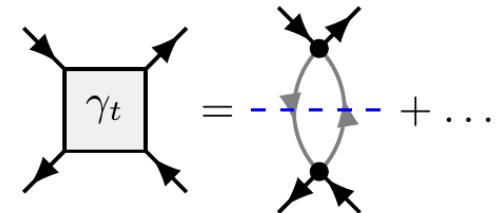
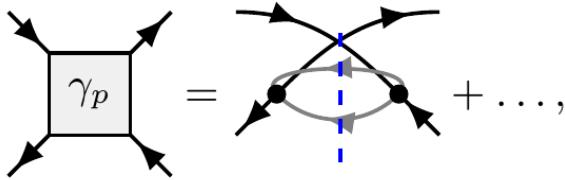
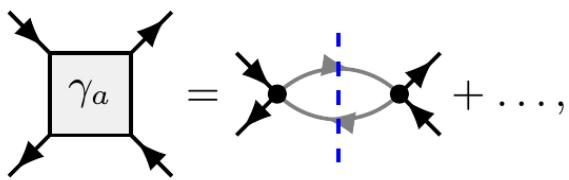
- make bare Green's function scale dependent $G_0 \rightarrow G_0(\Lambda) \Rightarrow \Sigma(\Lambda), \Gamma(\Lambda)$
- no flowing 2PI vertex $R \neq R(\Lambda)$
- differentiating BSE leads to multiloop fRG (mfRG) equations



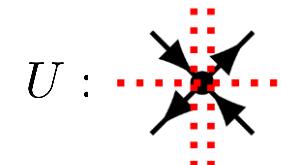
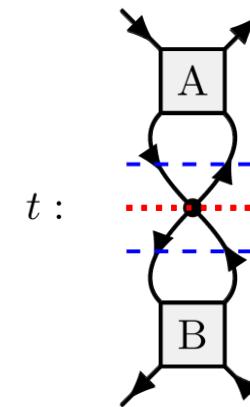
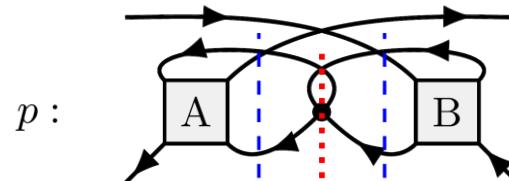
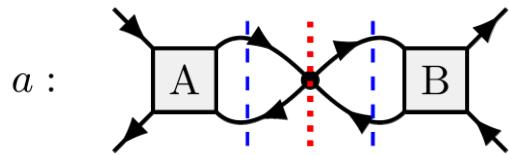
Kugler, von Delft PRB, PRL, NJP (2018)

2P reducibility vs. U reducibility

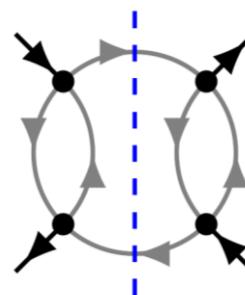
- 2P reducibility: cutting two Green's function lines, i.e., a bubble Π_r



- U reducibility: cutting out a bare interaction vertex U



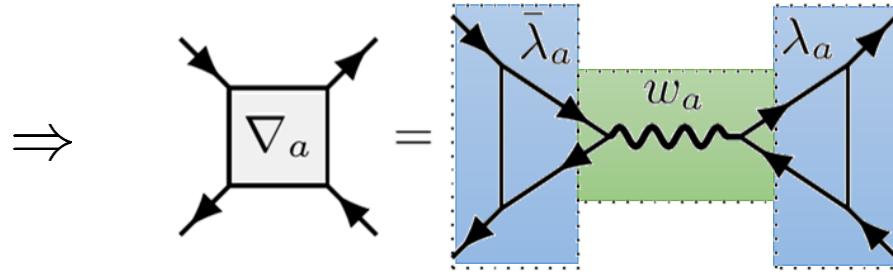
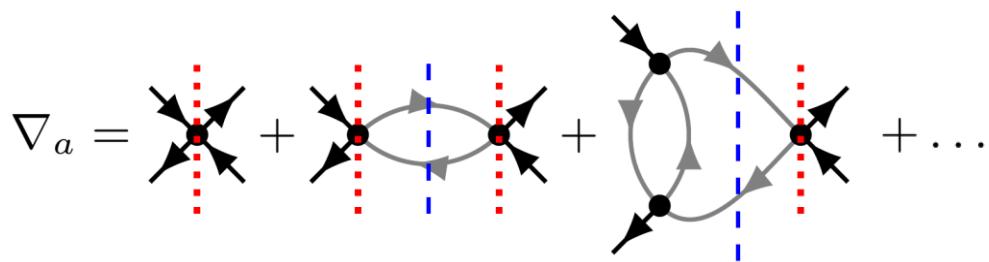
- U reducibility in $r \Rightarrow$ 2P reducibility in r
- but: 2P reducibility in $r \not\Rightarrow$ U reducibility



(except for U itself)

Single-boson exchange vertex

- U-reducible vertices $\nabla_r \rightarrow$ single-boson exchange



- effective interaction $w_r \rightarrow$ propagator of an exchange boson

$$w_r = U + U \cdot \chi_r \cdot U$$



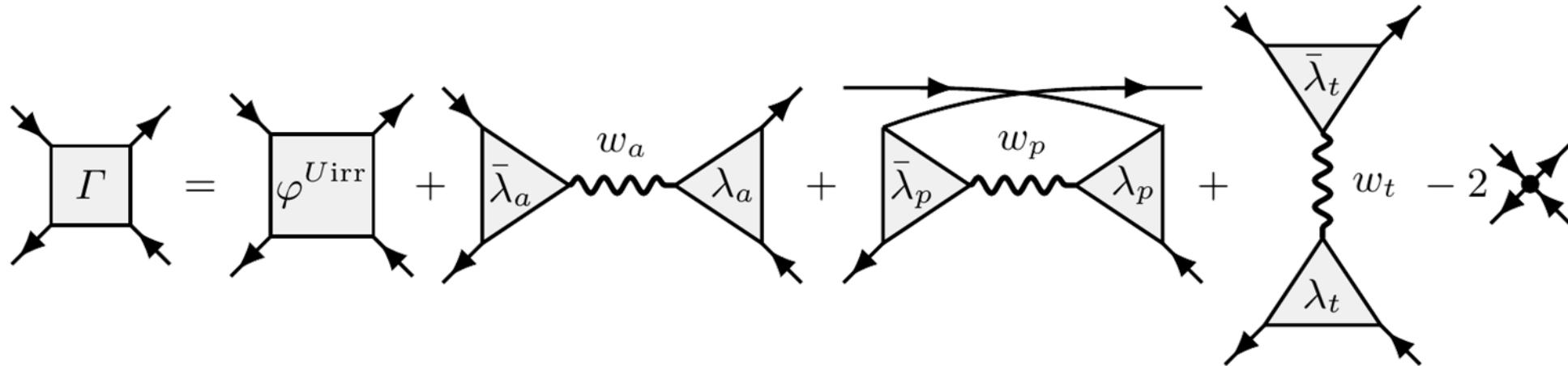
susceptibility

exchange field: $\psi = \sum_\nu \bar{c} c$

pairing field: $\phi = \sum_\nu c c$

- Hedin vertices $\lambda_r, \bar{\lambda}_r \rightarrow$ exchange process between one boson and two fermions

SBE Decomposition



Krien, Valli, Capone, Lichtenstein, Rohringer, Held, Kauch, Harkov PRBs (2019,2020,2021)

- SBE core $\varphi^{U\text{irr}}$:

$$M_a = \begin{array}{c} \text{---} \\ | \\ \text{---} \end{array} + \dots$$

multi-boson exchange

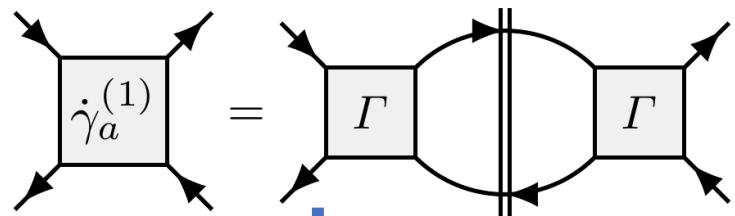
$$R - U = \begin{array}{c} \text{---} \\ | \\ \text{---} \end{array} + \dots$$

totally irreducible

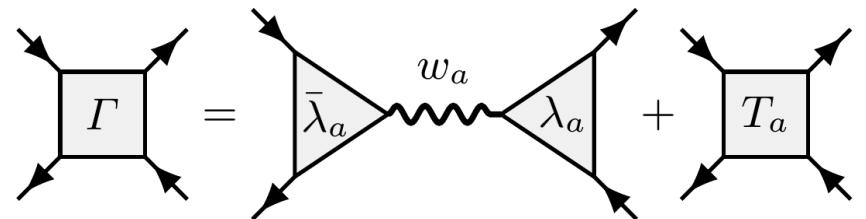
- closed system of equations: SBE equations (similar to parquet equations)
→ SBE mfRG

SBE fRG (one loop)

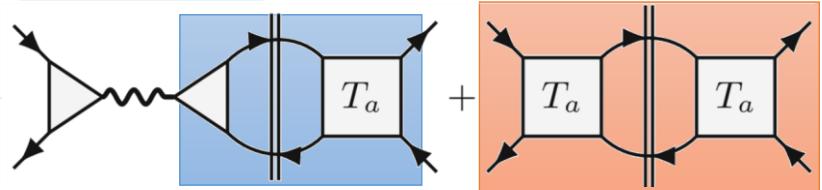
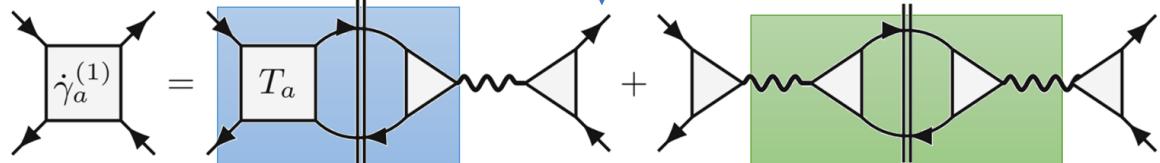
1-l fRG equation



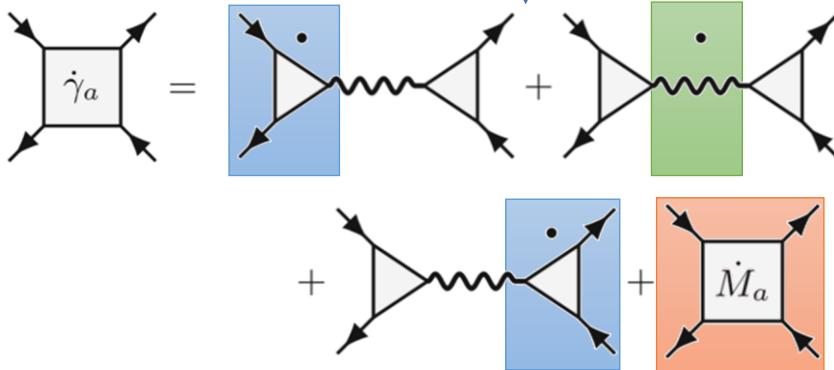
SBE decomposition



insertion



differentiation



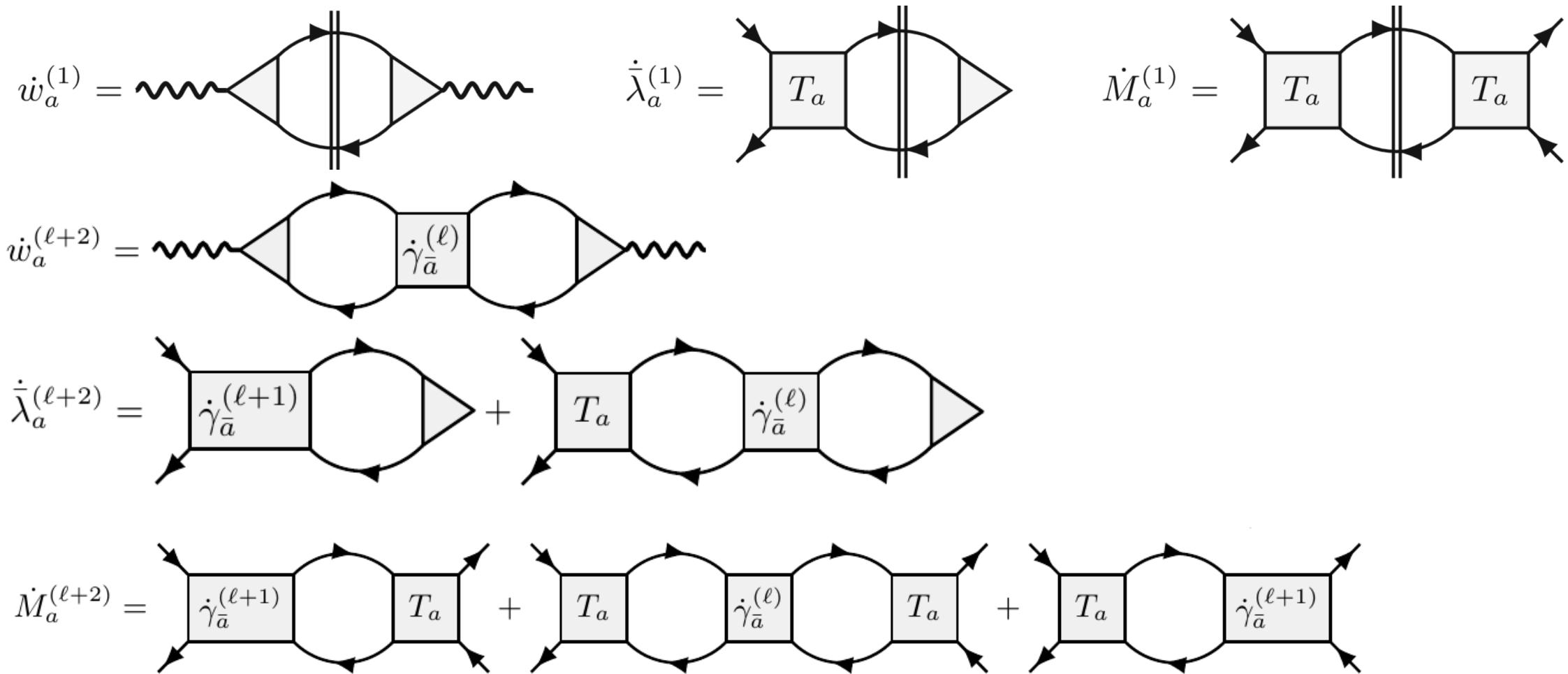
$$\Rightarrow \dot{w}_a^{(1)} = \text{[diagram: green box with wavy line and triangle]}$$

$$\dot{\lambda}_a^{(1)} = \text{[diagram: blue box with wavy line and triangle]}$$

$$\dot{M}_a^{(1)} = \text{[diagram: orange box with wavy line and triangle]}$$

Bonetti et al. PRR (2022)

mfRG for SBE ingredients



Asymptotic classes vs. SBE ingredients

- full vertex:

$$\Gamma$$

$$R + \sum_r \gamma_r$$

$$R + \sum_r [\mathcal{K}_1^r + \mathcal{K}_2^r + \mathcal{K}_{2'}^r + \mathcal{K}_3^r]$$

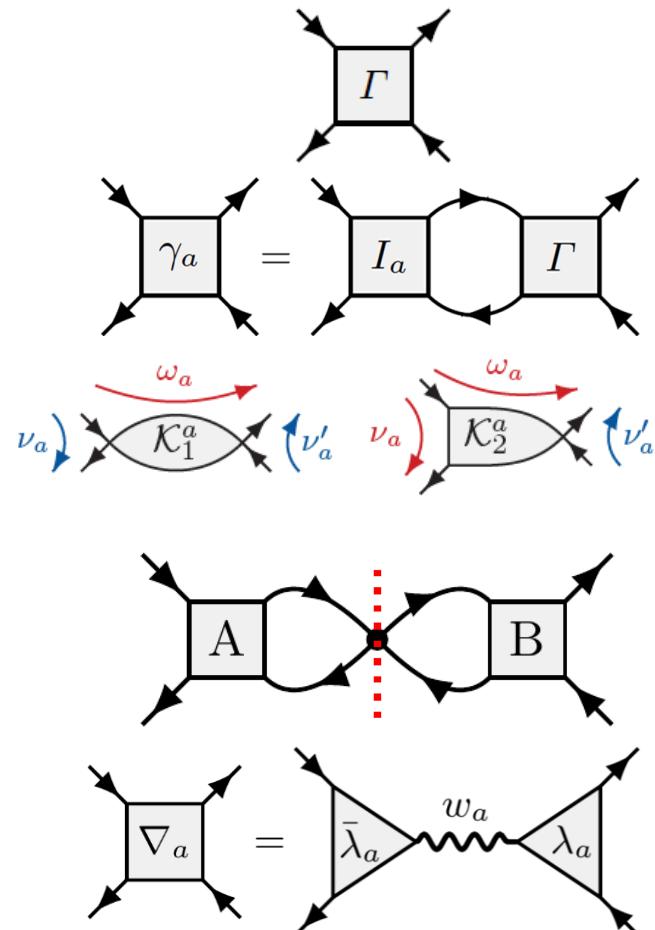
$$R - U + U + \sum_r [\bar{\lambda}_r \cdot w_r \cdot \lambda_r - U + M_r]$$

$$\varphi^{U_{\text{irr}}} - 2U + \sum_r \bar{\lambda}_r \cdot w_r \cdot \lambda_r$$

• SBE approximation: $\varphi^{U_{\text{irr}}} \rightarrow 0$

• \mathcal{K}_2 approximation: $\varphi^{\mathcal{K}_3} = R - U + \sum_r \mathcal{K}_3^r \rightarrow 0$

← includes more diagrams!



Wentzell et al. PRB (2020)

Bonetti et al. PRR (2022)

Advantages of the SBE decomposition

- SBE vertex with transparent physical meaning:

$$\mathcal{K}_3^r = (\bar{\lambda}_r - \mathbf{1}_r) \bullet w_r \bullet (\lambda_r - \mathbf{1}_r) + M_r \quad \nabla_r = \bar{\lambda}_r \bullet w_r \bullet \lambda_r$$

- at criticality: $\chi_r \sim w_r \gg 1$

$$\mathcal{K}_1^r, \mathcal{K}_2^r, \mathcal{K}_{2'}^r, \mathcal{K}_3^r \gg 1 \quad w_r \gg \bar{\lambda}_r, \lambda_r, M_r$$

- faster convergence in parquet equations
- extension of cores in frequency space

$$\varphi^{K_3}$$

Harkov, Lichtenstein, Krien PRB (2021)

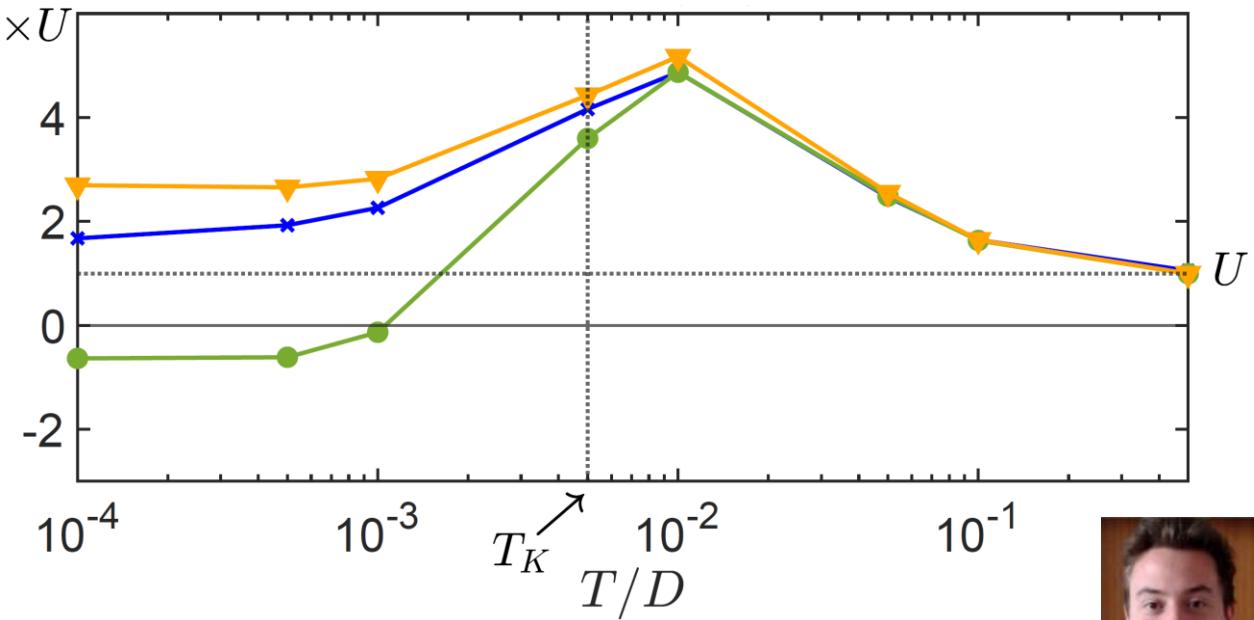
$$\varphi^{U^{\text{irr}}}$$

Fraboulet, Heinzelmann, Bonetti, Al-Eryani, Vilardi, Toschi, Andergassen arXiv (2022)

Outlook

- detailed numeric analysis left as future work
- Numerical renormalization group (NRG) study of single-impurity Anderson model (SIAM): better qualitative agreement for low temperatures & frequencies
→ promising starting point for DMF²RG
- concrete systems:
 - Keldysh mfRG of SIAM → non-eq. transport
Anxiang Ge, Elias Walter
 - Keldysh Hubbard model
Nepomuk Ritz
 - Fermi polaron problem
Marcel Gievers

- *— $\Gamma_{\uparrow\downarrow}(\pm 1, \pm 1, 0)$
- $U + \sum_{r,j \in \{1,2,2'\}} \mathcal{K}_{j,\uparrow\downarrow}^r$
- ▼— $\sum_r \nabla_{r,\uparrow\downarrow} - 2U$



Johannes Halbinger

Thank you for your attention!

$$\begin{aligned}\dot{w}_a^{(\ell+2)} &= \text{wavy line} \rightarrow \text{triangle} \rightarrow \text{square labeled } \dot{\gamma}_{\bar{a}}^{(\ell)} \rightarrow \text{triangle} \rightarrow \text{wavy line} \\ \dot{\lambda}_a^{(\ell+2)} &= \text{square labeled } \dot{\gamma}_{\bar{a}}^{(\ell+1)} \rightarrow \text{triangle} + \text{square labeled } T_a \rightarrow \text{square labeled } \dot{\gamma}_{\bar{a}}^{(\ell)} \rightarrow \text{triangle} \\ \dot{M}_a^{(\ell+2)} &= \text{square labeled } \dot{\gamma}_{\bar{a}}^{(\ell+1)} \rightarrow \text{triangle} \rightarrow \text{square labeled } T_a + \text{square labeled } T_a \rightarrow \text{square labeled } \dot{\gamma}_{\bar{a}}^{(\ell)} \rightarrow \text{triangle} \rightarrow \text{square labeled } T_a + \text{square labeled } T_a \rightarrow \text{square labeled } \dot{\gamma}_{\bar{a}}^{(\ell+1)} \rightarrow \text{triangle}\end{aligned}$$



M. Gievers, E. Walter, A. Ge, J. von Delft, F. Kugler EPJ (2022)