

# PION-PION SCATTERING FROM NUCLEON-MESON FLUCTUATIONS

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## **Relevant publications:**

JE and J.-P. Blaizot, *Phys. Rev. D* **105** (2022), 074031;

N. Cichutek, F. Divotgey, and JE, *Phys. Rev. D* **102** (2020), 034030;

F. Divotgey, JE, and M. Mitter, *Phys. Rev. D* **99** (2019), 054023;

JE, F. Divotgey, M. Mitter, and D.H. Rischke, *Phys. Rev. D* **98** (2018), 014024.



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# Nucleon-meson model with parity doubling

► **Euclidean effective action of investigated model:**

$$\begin{aligned} \Gamma_k [\varphi, \bar{\psi}_1, \psi_1, \bar{\psi}_2, \psi_2] = \int_x \left\{ \frac{1}{2} Z_k (\partial_\mu \varphi) \cdot \partial_\mu \varphi + V_k(\varphi^2) - h\sigma \right. \\ \left. + \bar{\psi}_1 \left[ Z_k^\psi \gamma_\mu \partial_\mu + y_{1,k} (\sigma + i\gamma_5 \vec{\pi} \cdot \vec{\tau}) \right] \psi_1 \right. \\ \left. + \bar{\psi}_2 \left[ Z_k^\psi \gamma_\mu \partial_\mu - y_{2,k} (\sigma - i\gamma_5 \vec{\pi} \cdot \vec{\tau}) \right] \psi_2 \right. \\ \left. + m_{0,k} (\bar{\psi}_1 \psi_2 + \bar{\psi}_2 \psi_1) \right\} \end{aligned}$$

► O(4)-vector  $\varphi = (\vec{\pi}, \sigma)$ , wave-function renormalizations  $Z_k, Z_k^\psi$

► Mirror-assigned fermion doublet  $(\psi_1, \psi_2)$

⇒ Chiral-invariant fermion mass:  $m_0 (\bar{\psi}_1 \psi_2 + \bar{\psi}_2 \psi_1)$

[Detar, Kunihiro ('89); Hatsuda, Prakash ('89); Jido, Nemoto, Oka, Hosaka ('00); Jido, Oka, Hosaka ('01)]

# Nucleon and chiral partner

- ▶ Fermion doublet  $(\psi_1, \psi_2)$  defined to be parity-even

- ▶ **Transformation into physical parity-opposite basis:**

$$\begin{pmatrix} \text{nucleon} \\ \text{chiral partner} \end{pmatrix} = \begin{pmatrix} N_+ \\ N_- \end{pmatrix} = \begin{pmatrix} \mathbb{1} & 0 \\ 0 & \gamma_5 \end{pmatrix} \begin{pmatrix} \cos \omega & \sin \omega \\ -\sin \omega & \cos \omega \end{pmatrix} \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix}$$

- ▶ Rotation angle  $\omega$  for the parity doublet:

$$\omega = \frac{1}{2} \arctan \left[ \frac{2m_0}{(y_1 + y_2)\sigma_0} \right], \quad \lim_{\sigma_0 \rightarrow 0} \omega = \frac{\pi}{4}, \quad \lim_{m_0 \rightarrow 0} \omega = 0$$

- ▶ Order parameter for chiral symmetry breaking: isoscalar condensate  $\sigma_0$
- ▶ Interpretation of chiral partner: lightest resonance  $N(1535)$
- ▶ Frequently discussed, also at nonzero densities

[Gallas, Giacosa, Rischke ('10); Weyrich, Strodthoff, von Smekal ('15); Marczenko, Blaschke, Redlich, Sasaki ('18); Tripolt, Jung, von Smekal, Wambach ('21); Minamikawa, Kojo, Harada ('21)]

## Compute S-wave pion-pion scattering lengths:

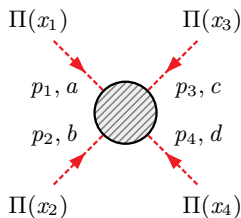
- ▶ Present a “fresh” view on the topic
- ▶ Include loop corrections in an efficient manner, i.e., by means of the **functional renormalization group (FRG)**
- ▶ Test and exemplify the class of **linear sigma models (LSMs)**

## Draw connection between the nucleon-meson model/LSM and the corresponding nonlinear sigma model (NLSM):

- ▶ Use stereographic coordinates
- ▶ Verify convergence of the employed low-energy expansion
- ▶ Elucidate subtleties and intricacies concerning the truncation of higher-derivative pion interactions in the effective action

# Pion-pion scattering

► Pion-pion scattering process:



► Formulation in terms of **stereographic projections** and elimination of the radial  $\theta$ -field through the **equation of motion (EOM)**

$$\Pi^{\bar{a}} = 2f_\pi \frac{\pi^{\bar{a}}}{\theta + \sigma}, \quad \bar{a} = 1, 2, 3, \quad \theta = |\varphi| = \sqrt{\vec{\pi}^2 + \sigma^2}$$

$$\text{Metric: } g_{\bar{a}\bar{b}} = \frac{16f_\pi^2 \delta_{\bar{a}\bar{b}}}{(4f_\pi^2 + \Pi^2)^2}, \quad \text{EOM: } \frac{\delta\Gamma}{\delta\theta} = 0$$

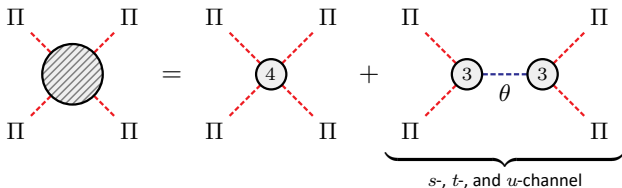
## Stereographic action

⇒ **Stereographic form of the (bosonic) effective action:**

$$\Gamma[\Pi] = \int_x \left[ \frac{\theta^2}{2} g_{\bar{a}\bar{b}} (\partial_\mu \Pi^{\bar{a}}) \partial_\mu \Pi^{\bar{b}} + \frac{1}{2} (\partial_\mu \theta) \partial_\mu \theta + V(\theta^2) - h\theta \frac{4f_\pi^2 - \Pi^2}{4f_\pi^2 + \Pi^2} \right],$$

$$\theta = f_\pi + \epsilon \xi_1(\Pi) + \epsilon^2 \xi_2(\Pi) + \dots, \quad \epsilon = \frac{M_\pi^2}{M_\sigma^2}$$

- ▶ Formulation constitutes **low-energy expansion** in  $\epsilon \simeq 0.0844$
- ▶ To lowest order, NLSM on three-sphere recovered:  $\theta = f_\pi$
- ▶ Employing EOM equivalent to summation of all tree diagrams (based on  $\Gamma$ ):



# Scattering amplitude

⇒ Computation of the scattering amplitude reduces to a **single Feynman diagram**

$$\text{EOM: } \frac{\delta\Gamma}{\delta\theta} = 0 \quad \Rightarrow \quad \mathcal{M}_{\pi\pi} \sim \begin{array}{c} \Pi \qquad \qquad \Pi \\ \diagdown \qquad \diagup \\ \text{---} \text{---} \text{---} \text{---} \\ \diagup \qquad \diagdown \\ \Pi \qquad \qquad \Pi \end{array} \equiv i\Gamma^{(4)}[\Pi]$$

► General structure of scattering amplitude:

$$\mathcal{M}_{\pi\pi}^{abcd}(s, t, u) = i\mathfrak{A}(s, t, u) \delta^{ab} \delta^{cd} + i\mathfrak{A}(t, u, s) \delta^{ac} \delta^{bd} + i\mathfrak{A}(u, s, t) \delta^{ad} \delta^{bc}$$

► **S-wave isospin-zero and isospin-two scattering lengths:**

$$a_0^0 = \frac{1}{32\pi} [3\mathfrak{A}(s, t, u) + \mathfrak{A}(t, u, s) + \mathfrak{A}(u, s, t)] \Big|_{s=4M_\pi^2, t=u=0},$$

$$a_0^2 = \frac{1}{32\pi} [\mathfrak{A}(t, u, s) + \mathfrak{A}(u, s, t)] \Big|_{s=4M_\pi^2, t=u=0}$$

# Scattering lengths

⇒ Findings up to  $\mathcal{O}(p^{12})$ :

$$a_0^0 = \frac{M_\pi^2}{32\pi f_\pi^2} (7 + 29\epsilon + 108\epsilon^2 + 432\epsilon^3 + 1728\epsilon^4 + 6912\epsilon^5 + 15360\epsilon^6 + 12288\epsilon^7) \longrightarrow \frac{M_\pi^2}{32\pi f_\pi^2} (1 - \epsilon) \left[ 7 + 9 \sum_{n=1}^{\infty} (4\epsilon)^n \right],$$

$$a_0^2 = -\frac{M_\pi^2}{16\pi f_\pi^2} (1 - \epsilon), \quad \text{expressions correct up to } \epsilon^5 = \left( \frac{M_\pi^2}{M_\sigma^2} \right)^5 \sim 10^{-6}$$

- ▶ Geometric series convergent for  $\epsilon < 1/4$ , yields “radius of convergence”
- ▶ Reference values: Experiment/Chiral perturbation theory

$$a_0^0 = 0.2198 \pm 0.0126, \quad a_0^2 = -0.0445 \pm 0.0023$$

[Weinberg ('66); Gasser, Leutwyler ('84, '85); Leutwyler ('94, '12); Bijmans et al. ('96, '97); Colangelo, Gasser, Leutwyler ('00, '01); Ananthanarayan, Colangelo, Gasser, Leutwyler ('01); Gasser ('09)]



## General procedure

1. Compute dressed vertices from **FRG integration** within nucleon-meson model

$$\Gamma_{UV} = S \text{  } \Gamma_{IR} = \Gamma$$

$$\partial_k \Gamma_k = \frac{1}{2} \text{tr} \left[ (\partial_k \mathcal{R}_k) \left( \Gamma_k^{(2)} + \mathcal{R}_k \right)^{-1} \right]$$

⇒ Constraint: chiral symmetry breaking at  $\Lambda_\chi = 1.2 \text{ GeV} \simeq 4\pi f_\pi$

2. Transform to nonlinear **stereographic action**

3. Perform **low-energy expansion** up to  $\mathcal{O}(p^{12})$ , accurate up to  $\epsilon^5$

4. Compute **S-wave isospin-zero and isospin-two scattering lengths**

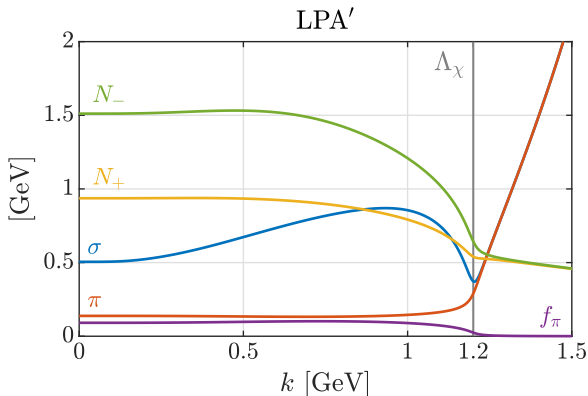
⇒ Compare to mean-field (MF) and one-loop calculations (“backwards”)

⇒ Comment on model parameters:  $\sigma$ -mass and chiral-invariant nucleon mass  $m_0$

[Wetterich ('93); Ellwanger ('94); Morris ('94); Pawłowski ('07); Dupuis et al. ('21)]

# FRG integration

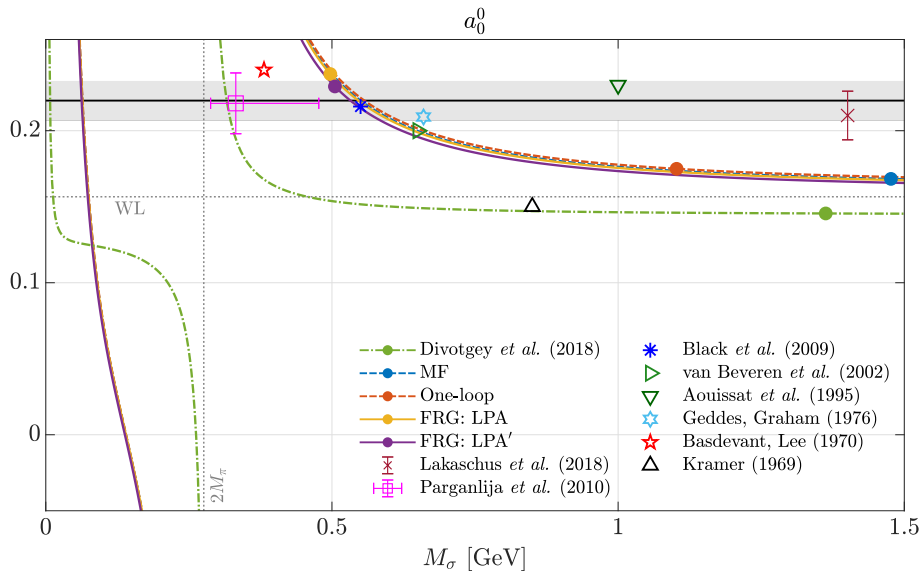
- ▶ FRG flow in LPA'-truncation (found to reproduce scattering lengths):



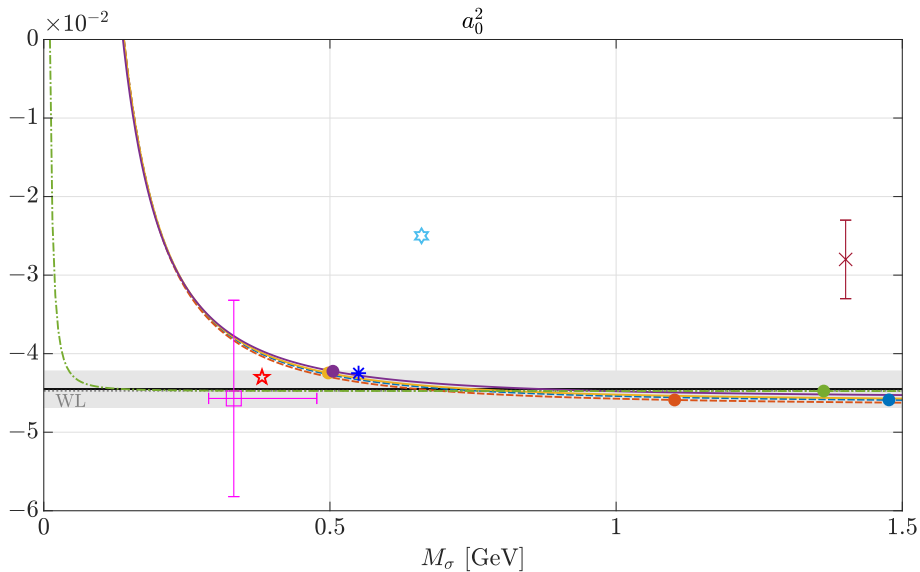
⇒  $\sigma$ -mass in FRG (with given conditions) typically around  $M_\sigma \simeq 500$  MeV

⇒ Chiral-invariant nucleon mass amounts to  $m_0 = 824.5$  MeV in the IR

# Isospin-zero scattering length

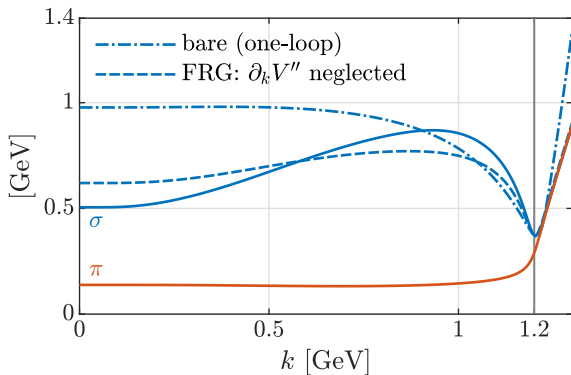


# Isospin-two scattering length



## Isoscalar mass in FRG

- Scale dependence of  $M_\sigma$  for different approximations:

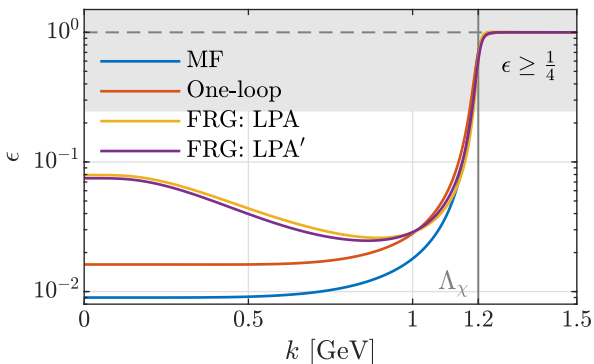


⇒ “Down-bending” of  $M_\sigma$  with “advancing” the approximation

⇒  $M_\sigma$  successively shrinks, leading to improved results on the scattering lengths

## Radius of convergence

- Dynamic ratio  $\epsilon = M_\pi^2/M_\sigma^2$ :

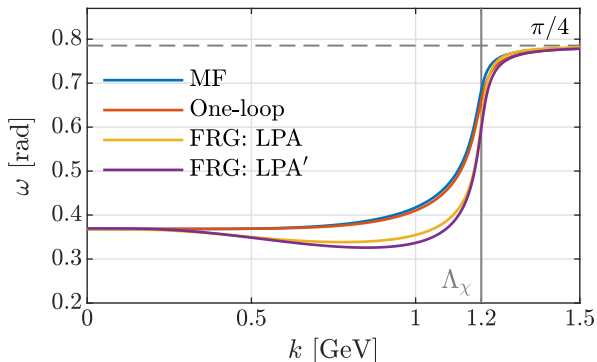


- Geometric series convergent for  $\epsilon < 1/4$ , divergent for  $\epsilon \geq 1/4$

⇒ Low-energy expansion valid for  $k \lesssim \Lambda_\chi$  (“radius of convergence” w.r.t.  $k$ -scale)

# Rotation angle

- Dynamic rotation angle  $\omega$ :



- $\omega$  parametrizes transformation to physical basis

⇒ Angles coincide in the IR, reflecting the common fermion parameters

## Conclusion

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- ▶ Pion-pion scattering as **interesting application** of parity-doublet model
- ⇒ Scattering lengths brought into **simultaneous agreement** with experiment, with the requirement of chiral symmetry breaking roughly at  $4\pi f_\pi$

$$\text{FRG: } a_0^0 = 0.2291, \quad a_0^2 = -0.0422;$$

$$\text{experiment: } a_0^0 = 0.2198 \pm 0.0126, \quad a_0^2 = -0.0445 \pm 0.0023$$

- ▶ Embedding of the topic into the **dynamical context of the FRG**
- ▶ **“Backwards”-determination of  $\sigma$ -mass**: insisting on pertinent scales and employing common parameters when “reducing” the approximation
- ⇒  $\sigma$ -mass **tends to shrink** with “advancing” the approximation, from  $M_\sigma > 1 \text{ GeV}$  (MF, one-loop) towards  $M_\sigma \simeq 500 \text{ MeV}$  (FRG)
- ⇒ Model study suggests a **chiral-invariant mass** around  $m_0 = 824.5 \text{ MeV}$



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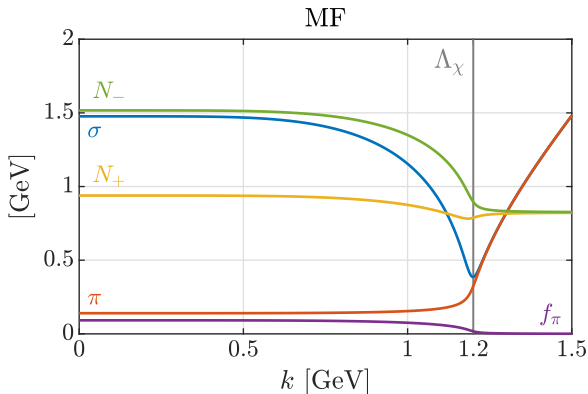
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# MF approximation

- ▶ MF integration (with chiral-symmetry breaking scale  $\Lambda_\chi$ ):

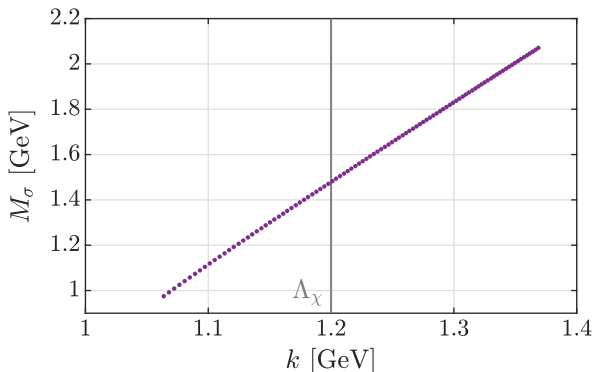


$\Rightarrow$   $\sigma$ -mass in MF (with given conditions):  $M_\sigma > 1$  GeV

- ▶ Chiral-invariant nucleon mass fixed to LPA'-value to achieve comparability

## Isoscalar-mass prediction in MF

- $M_\sigma$  in the IR, depending on the scale of chiral symmetry breaking:



$\Rightarrow M_\sigma \simeq 1.5$  GeV in the IR, with the required breaking scale of  $\Lambda_\chi \simeq 4\pi f_\pi$

$\Rightarrow M_\sigma$  uniquely determined, once fermionic model parameters fixed