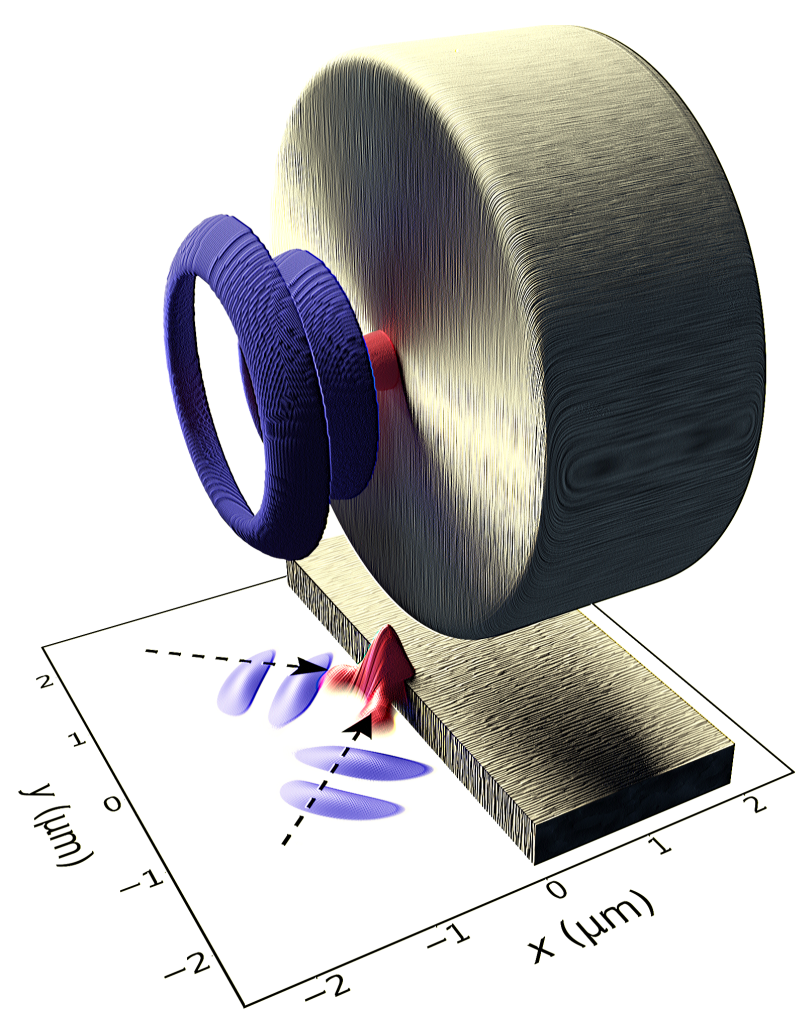


INTRODUCTION

A **tightly focused single cycle** laser results in the highest achievable intensity on the expense of the least energy for a fixed laser power, being in the so called **λ^3 regime** [1].



Simulation Setup.

Interaction of an intense λ^3 -laser with a foil results in **high** laser to **γ -photon** energy conversion efficiency, κ_γ [2].

This work performs a parametric study of the λ^3 -laser-target interaction, where the variables include the laser **polarisation** mode (radial, linear and azimuthal), the target **thickness** and electron number **density**.

The results are obtained via 3D QED PIC simulations [3].

The laser ($\lambda = 1 \mu\text{m}$) energy is 280 J, contained within **3.4 fs** FWHM duration, corresponding to an 80 PW laser. The laser is focused on a $\sim \lambda/2$ spot, at an intensity of 10^{25} Wcm^{-2} .

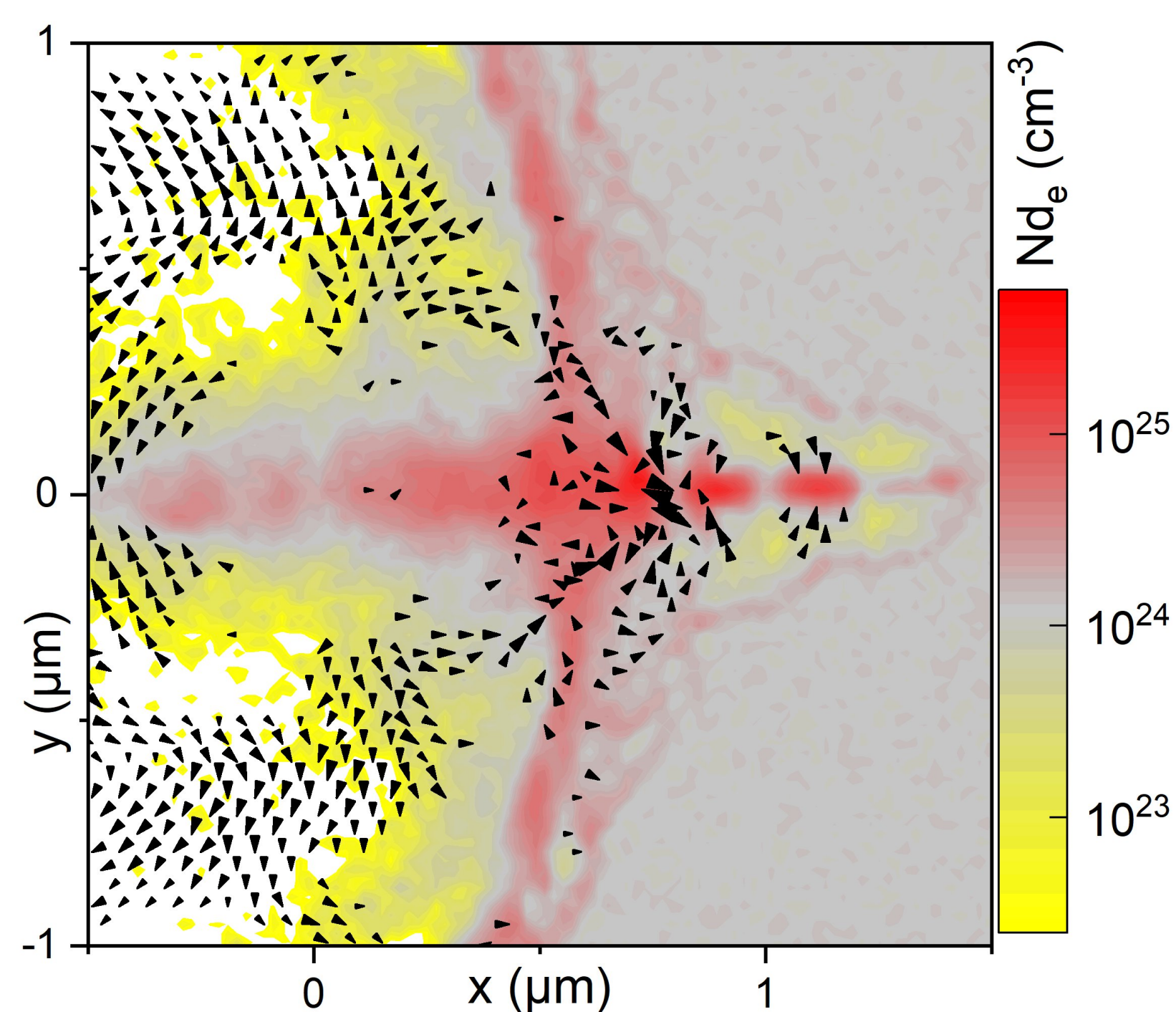
A practical advantage of our result is the elimination of the need of a preplasma in front of the target [4,5] for maximising κ_γ .

OVERVIEW OF THE RESULTS

Our multi-parametric study reveals that κ_γ is enhanced by a **radially polarised laser**, reaching $\sim 50\%$, compared to $\sim 40\%$ for linearly and $\sim 30\%$ for azimuthally polarised lasers.

For the radially polarised laser, optimum κ_γ occurs for a $2 \mu\text{m}$ thick foil having an electron number density of $1.2 \times 10^{24} \text{ cm}^{-3}$, corresponding to **titanium**.

The radially polarised λ^3 -laser is characterised by a dominant **longitudinal** electric field. This field component benefits the formation of a **conical hole** within the foil volume, enhancing the coupling of the laser field to the electrons. For a linearly polarised laser the cavity is asymmetric, suppressing its formation.



Electron density during the interaction of an 80 PW λ^3 focused laser with a titanium foil. The arrows represent the electric field vectors.

Interference of the tightly focused fields results in uneven contribution to the parameter χ_e (that governs γ -photon formation) by the electric and magnetic terms, with **non-cancelling** contribution even for electrons co-propagating the laser pulse.

The high χ_e values result in high κ_γ , which in turn creates copious pairs, with **positrons** obtaining $\sim 10\%$ of the laser energy.

REFERENCES

- [1] G. Mourou et al. *Plasma Phys. Rep.* **28** 12 (2002)
- [2] P. Hadjisolomou et al. *Phys. Rev. E* **104** 015203 (2021)
- [3] C. P. Ridgers et al. *Phys. Rev. Lett.* **108** 165006 (2012)
- [4] T. Nakamura et al. *Phys. Rev. Lett.* **108** 195001 (2012)
- [5] K. V. Leznin et al. *Phys. Plasmas* **25** 123105 (2018)
- [6] G. Battistoni et al. *Ann. Nucl. Energy* **82**, 10-18 (2015)

γ -PHOTON OPTIMUM, AND POWER SCALING

The dependency of κ_γ on the laser **power** was examined in the range 1–300 PW for a $2 \mu\text{m}$ thick titanium foil and for a radial polarisation mode.

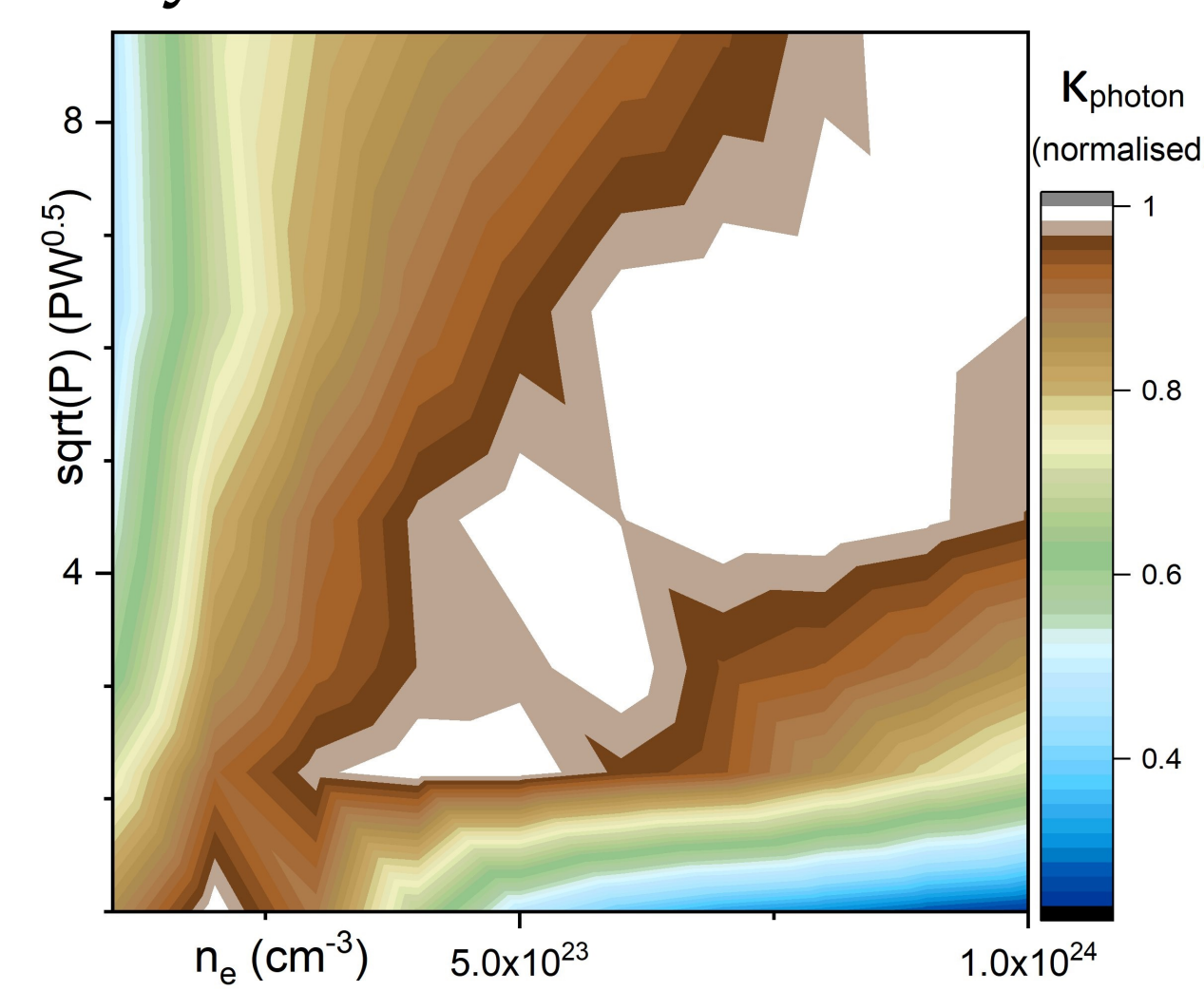
The κ_γ **increases by increasing the power**. The other particles saturate at a conversion efficiency of $\sim 10\%$, indicating that κ_γ should saturate at $\sim 60\%$.

At ~ 100 PW the electron-positron pair generation is maximised relatively to the γ -photon generation.

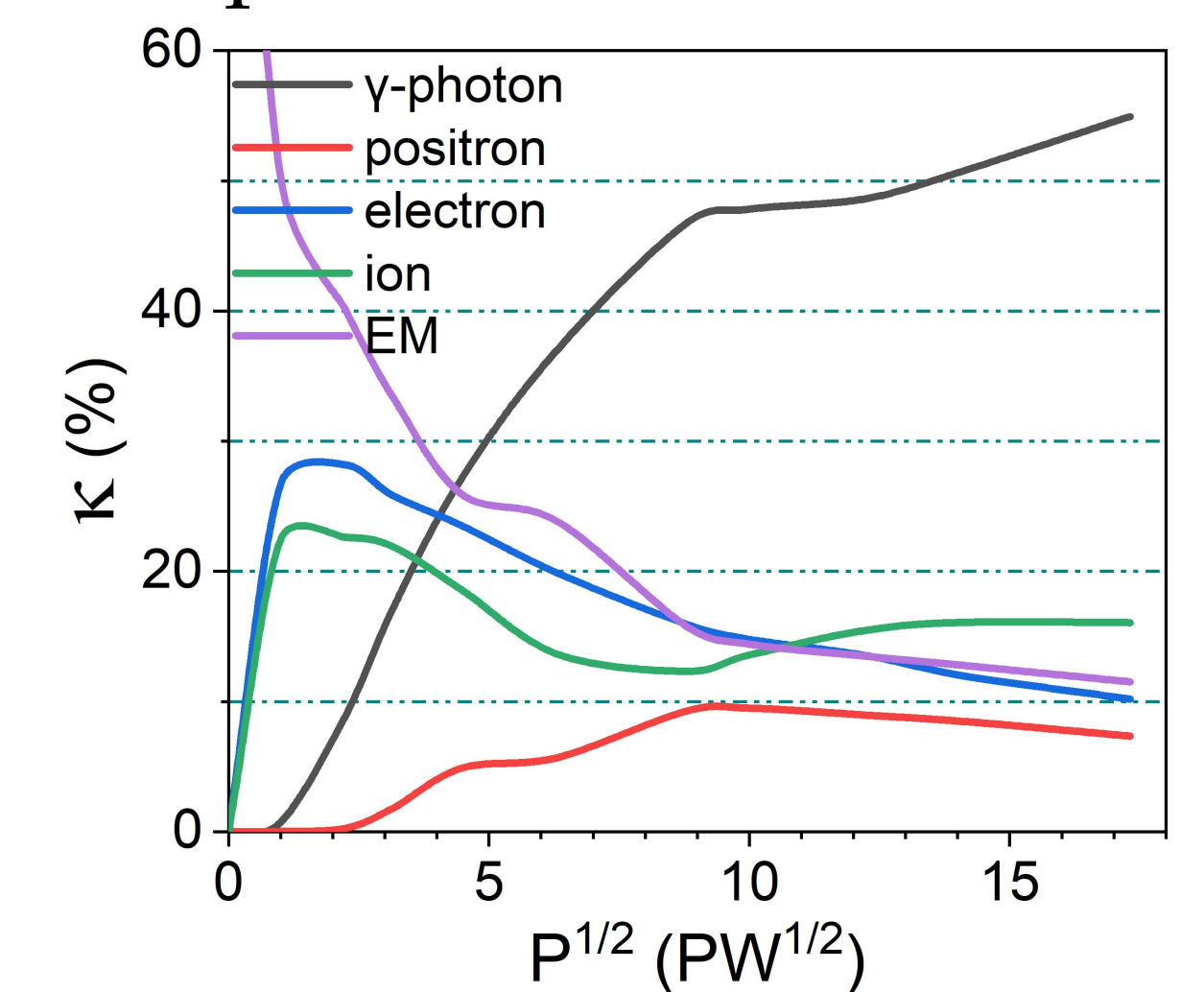
The optimum κ_γ strongly depends on the electron density at lower laser intensities. The density dependency faints as power increases.

At optimum density, a 1 PW gives a κ_γ of $\sim 3\%$.

As power increases, the same does the number of times an electron/positron emits a γ -photon, with the dependency being approximately linear with the dimensionless amplitude.



(a) The κ_γ (normalised) at each power and for different electron densities.



(b) Laser (radial) energy conversion efficiency as a function of the root of power for all particle species.

γ -FLASH INTERACTION WITH HIGH-Z TARGET

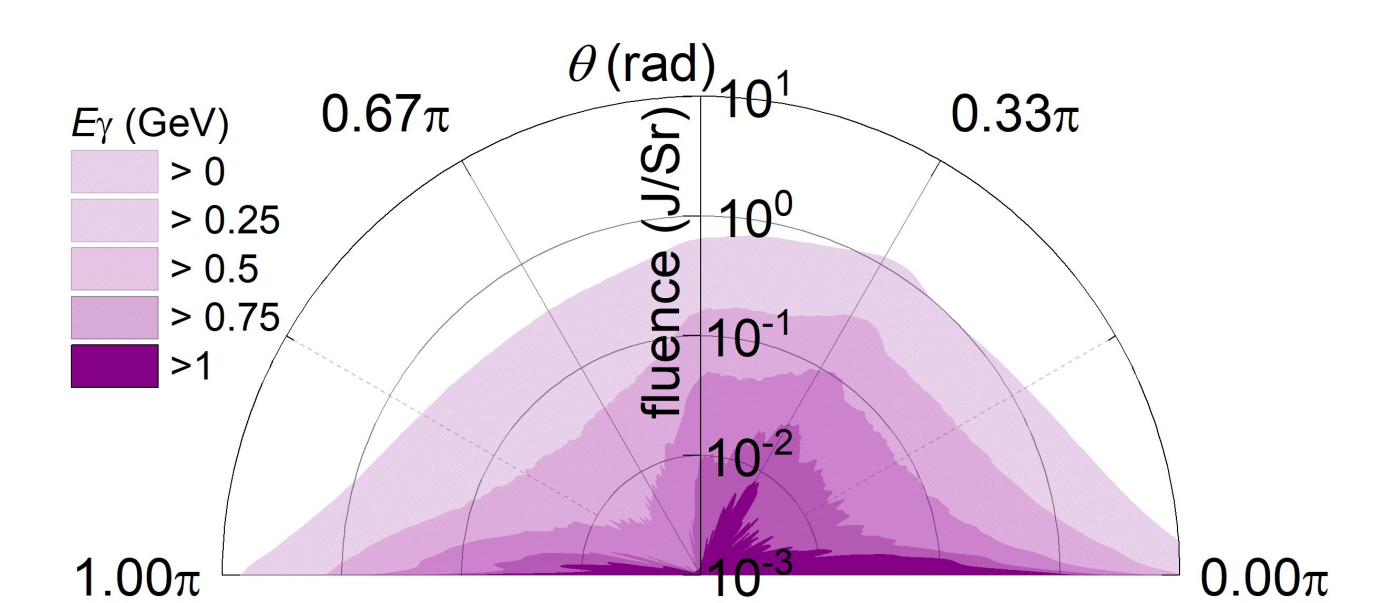
Most energetic γ -photons are generated within one laser period. During that time, γ -photon loss to pairs also occurs.

High energy γ -photons form an almost collimated beam in the **forward/backward** direction and at 60° .

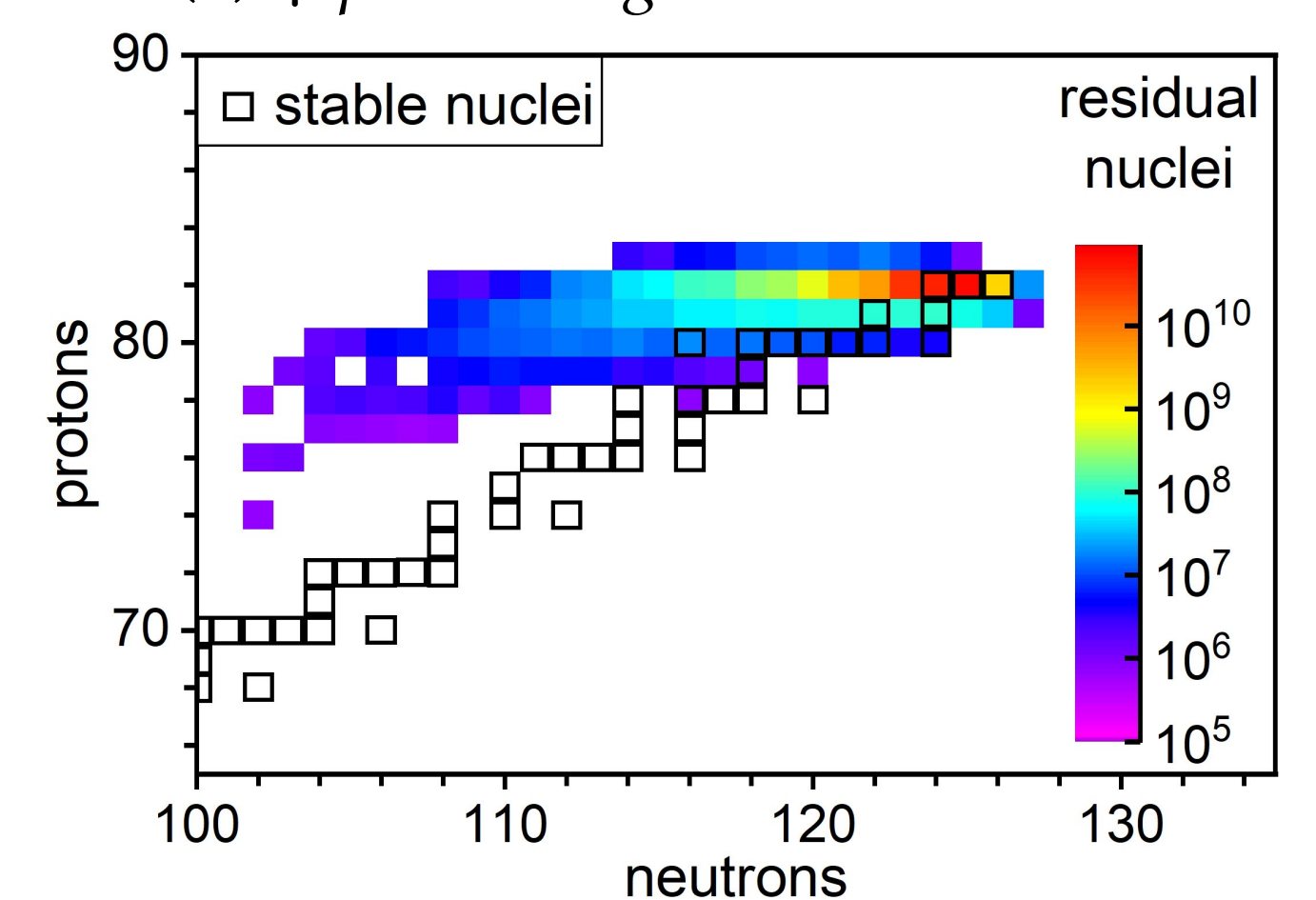
The γ -photons at 60° suffer geometrical dissipation due to their large solid angle compared to those on axis, resulting the latter having a significantly higher energy density.

The γ -photons form a ~ 30 PW **spherically expanding γ -flash**.

The PIC γ -photons are transported with **FLUKA** [6] (Monte Carlo) to simulate their interactions with matter.



(a) γ -photon angular distribution.



(b) Nuclide chart after importing PIC γ -photons in FLUKA (target 10 cm Pb).

CONCLUSIONS

A single-cycle, tightly-focused 80 PW pulse leads to the ultraintense **λ^3 regime**.

The laser **polarisation**/power and target thickness/density strongly affect the γ -photon energy conversion efficiency.

Optimum κ_γ is obtained by a radial polarisation mode, resulting in a directional γ -flash.