## A Reduced Model for Breit-Wheeler Pair Production by a Gaussian or Laguerre-Gaussian Laser Beam of Arbitrary Polarization and a Gamma flash

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# Reduced model for laser spatio-temporal shape



A model for pair production in head-on collision of a laser pulse and gamma photons (electrons)



Parametric study at fixed laser energy and focus on the soft shower regime of pair production (low secondary pairs)

- Time dependence and 'building block'
- Role of the laser pulse's spatial shape (Laguerre-Gauss beams)







# Reminder: Nonlinear Breit-Wheeler

## Gamma photon propagating in a laser field

Rate of pair production by photon decay (LCFA)  $W_0 = \frac{2\alpha mc^2}{2\hbar} = 3.8 \times 10^{18} \mathrm{s}^{-1}$  $\gamma_{\gamma} = \frac{\hbar \omega_{\gamma}}{m_o c^2}, \quad a_0 = \frac{eE_0}{mc\omega}, \quad \chi_{\gamma} \propto \gamma_{\gamma} \frac{E_0}{E_S} \propto \gamma_{\gamma} a_0$ 

$$W_{\rm BW} = W_0 \frac{b_0(\chi_{\gamma})}{\gamma_{\gamma}}$$

In the soft shower regime (no secondary gamma photons)

$$\frac{d}{dt}N_{\gamma} = -W_{\rm BW}(t)N_{\gamma}$$

The probability for a photon to decay in  $\Delta t$ 

 $P(\Delta t) = 1 - \exp\left[-\int_{t_0}^{t_0 + \Delta t} W_{\text{BW}}(t')dt'\right]$ 

Key element



# Pair creation in a plane wave

#### Head-on collision with a plane wave

**Field**: 
$$\mathbf{E}(z, t) = \hat{\frac{E_0}{\sqrt{1+\epsilon^2}}} \left[ \sin(\omega t - kz) \mathbf{\hat{x}} + \epsilon \cos(\omega t - kz) \mathbf{\hat{y}} \right]$$
 (normalisation constant energy)

Photon quantum parameter :  $\chi_{\gamma}(t) = 2\gamma_{\gamma} \frac{E_0}{E_{\gamma}}$ 

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## Building block : the rate integral over $\tau/4$

$$\int_{0}^{\tau/4} W_{\rm BW}(t') dt' = \frac{W_{0}}{2\omega\gamma_{\gamma}} \underbrace{\int_{0}^{\pi} b_{0}(\chi_{0}\Psi_{\varepsilon}(\varphi)) d\varphi}_{\mathcal{J}_{\varepsilon}(\chi_{0})}$$

$$\int_{0} \frac{\sin^2(2\omega t) + \varepsilon^2 \cos^2(2\omega t)}{1 + \varepsilon^2}$$

**Examples :**   $\chi_{\gamma}(t) = \chi_0 |\sin(2\omega t)|$  for LF  $\chi_{\gamma} = \chi_0 / \sqrt{2}$  for CP

 $\Psi_{\varepsilon}(2\omega t)$ 

Time dependent and  $\tau/4$  periodic (half period crossing, e.g. LP ->)





# Probability in a LP plane wave

Example LP (
$$\varepsilon = 0$$
) :  $\mathcal{I}_0(\chi_0) = \int_0^{\pi} b_0$   
analytica

$$\mathcal{I}_{0}(\chi_{0}) \simeq \pi b_{0}(\chi_{0}) \min \left\{ F\left(\sqrt{\frac{2b_{0}(\chi_{0})}{3\chi_{0}b_{0}'(\chi_{0})}}\right), f \right\}$$
$$\left( \begin{array}{c} F(s) = \sqrt{2/\pi}s \, \mathrm{erf}\left(\pi\sqrt{2}/(4s)\right) \\ f = \frac{1}{\pi} \int_{0}^{\pi} \mathrm{sin}^{2/3}(\varphi) \mathrm{d}\varphi \end{array} \right)$$

The probability to create a pair after a half period crossing  $P(\tau/4) = 1 - \exp\left[-\frac{W_0}{2\omega\gamma\gamma}\mathcal{I}_0(\chi_0)\right]$ **N.B.**  $P(\tau/4, a_0, \gamma_{\gamma})$ 

- $v_0(\chi_0 \sin(\varphi)) d\varphi$  can be solved numerically or
- al approximation :



# Probability after half period crossing (LP)



Numerical solution of the equicontours of constant probability : straight lines at low  $a_0, \gamma_{\gamma}$  but cross values of constant  $\chi_0$ minimum  $a_0$  for a given probability, along the line  $\chi_0 \sim 16.5$ 



# Analytical approximation



 $\begin{cases} F(s) = \sqrt{2/\pi}s \operatorname{erf}(\pi\sqrt{2}/(4s)) \\ f = \frac{1}{\pi} \int_0^{\pi} \sin^{2/3}(\varphi) d\varphi \end{cases}$ 



# Cross benchmark with PIC/MC

## Excellent match with PIC/MC **Smilei**)\* simulations $P(\tau/4)$



Generalization to more complex space and time dependences

\*https://github.com/SmileiPIC/Smilei





# Finite pulse duration and parameter space



# Laguerre-Gauss beams and space dependence

Complex envelope  $u_{p\ell}(\rho, \phi, z) = C_{p\ell} \frac{w_0}{w(z)} \left( \frac{\sqrt{2}\rho}{w(z)} \right)$ 

Solutions of the paraxial equation with constant total energy

## The beam's transverse size increases with $\ell$ , the amplitude decreases



$$\left(\frac{\sqrt{2}\rho}{w(z)}\right)^{|\ell|} L_p^{|\ell|} \left(\frac{2\rho^2}{w^2(z)}\right) \exp\left[-\frac{\rho^2}{w^2(z)}\right] \exp\left[-i\psi_{pl}(z) + i\ell\phi + i\frac{z\rho^2}{w^2(z)}\right]$$
  
Generalised phase  
Polynomial factor

#### Maximum amplitude

#### Photon quantum parameter $a_0 = 2000, \gamma_{\gamma} = 400$











# Total cross section of pair creation

Compute the time integrated probability for a finite time duration with fields' spatial distribution given by the LG fi

Example 
$$\ell = 5$$
  
 $a_0 = 2000, \gamma_{\gamma} = 400$   
 $\tau_{\rm FWHM} = 5\tau$ 

Prediction







held 
$$\frac{1}{\sqrt{\ell!}} \left(\frac{\sqrt{2}\rho}{w_0}\right)^{|\ell|} \exp\left[-\frac{\rho^2}{w_0^2}\right]$$



# Pair creation in LG beams

#### First set of PIC/MC **Smilei**) simulations $a_0 = 2000, \gamma_{\gamma} = 400, \tau_{\text{FWHM}} = 5\tau, \chi_0 = 4.85$





# Pair creation in less focused Gaussian beams

Comparison with Extended Gaussian (EG)  $a_0 = 2000, \gamma_{\gamma} = 400, \tau_{\rm FWHM} = 5\tau, \chi_0 = 4.85$ 



Extended Gaussian (EG) are constructed in order to have the same total energy and the same peak value as the LG of the corresponding order

 $\ell \neq 0, p=0$ 



#### Number of produced pairs in EG higher than LG for $\ell < 4$ (bigger interaction area)







# LG beams vs EG for lower intensities

#### Importance of $a_0$ versus $\gamma_{\nu}$



High probability isocontours (P>0.63)



**Start from lower** a<sub>o</sub> **: increasing** the size is not enough to compensate for the field's amplitude diminution, no improvement on pair production.



# Guidelines for upcoming facilities



Dotted red lines = constant  $\chi_0$  (0.1;1;16.5)



## **Strategy for optimisation**

**Facilities in black**: increase focal spot size (no tight focus) as the maximum  $10^{23}$ probability is already  $\sim 1$ 

**Facilities in white** : need to increase laser amplitude (tight focus) in order to have significant pair creation







# Thank you

More details : A. Mercuri-Baron et al 2021 New J. Phys. 23 085006