

Spin-foam models of Lorentzian quantum space-time

Emerging geometry at the semi-classical limit

$$Z(\Delta^*) = \sum_{j \rightarrow \mathcal{F}} \sum_{\iota \rightarrow \mathcal{E}} \int d\chi_i \Delta_{\chi_i} \left[\prod_{f \in \mathcal{F}} j_f^2 (\gamma^2 + 1) \right] \left[\prod_{e \in \mathcal{E}} \text{diagram} \right] \left[\prod_{v \in \mathcal{V}} \text{diagram} \right]$$

How to quantize gravity

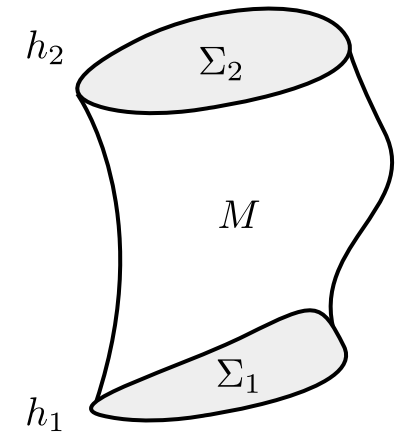
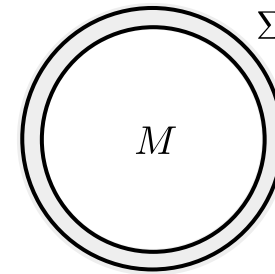
- Quantization of gravity via “sum over histories”, reminiscent of transition amplitudes.

- Would like to compute

$$Z(h) = \int_{h=g|_{\partial M}} \mathcal{D}g e^{iS_{\text{GR}}}$$

Note the i

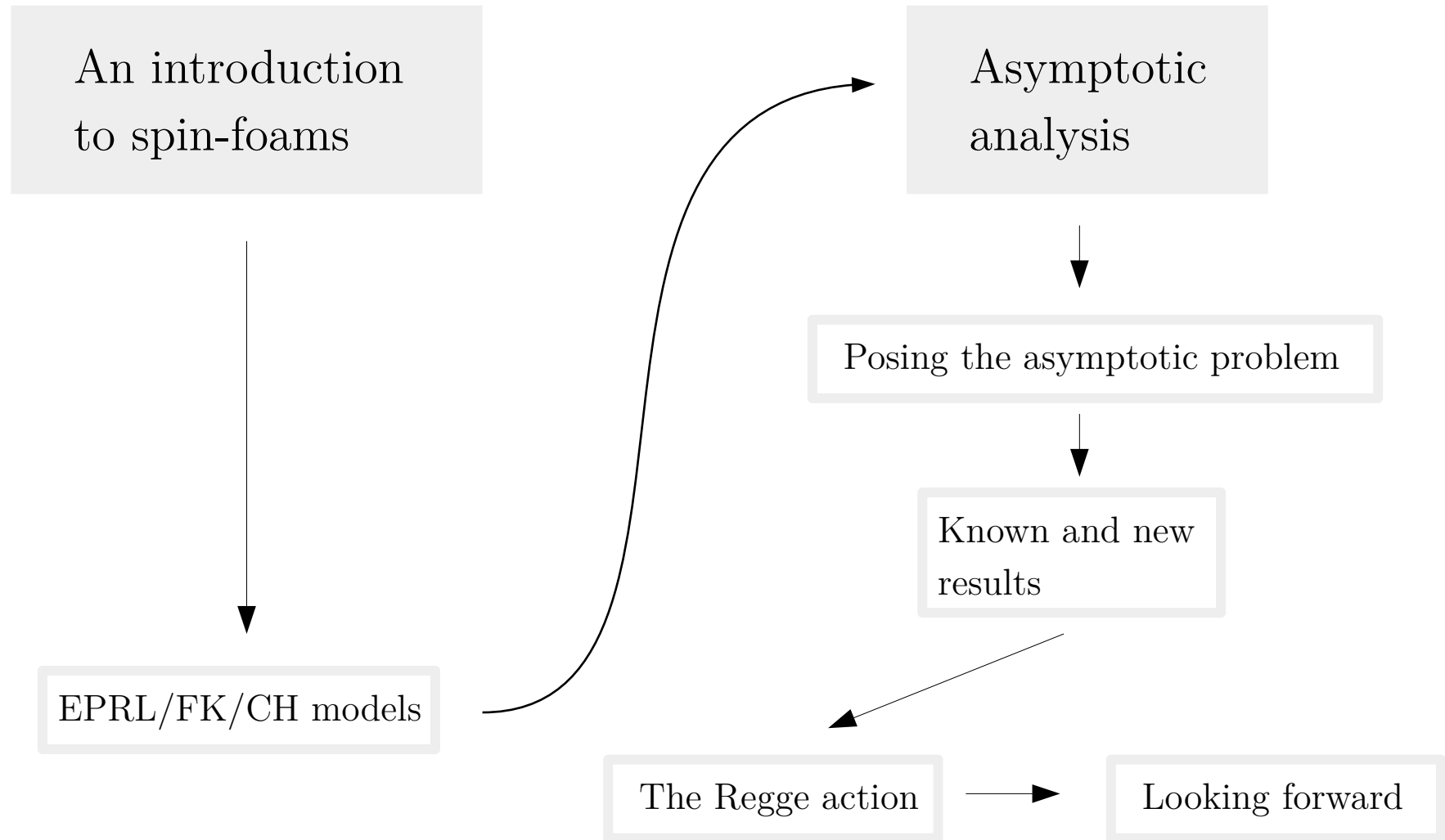
but this is technically difficult.



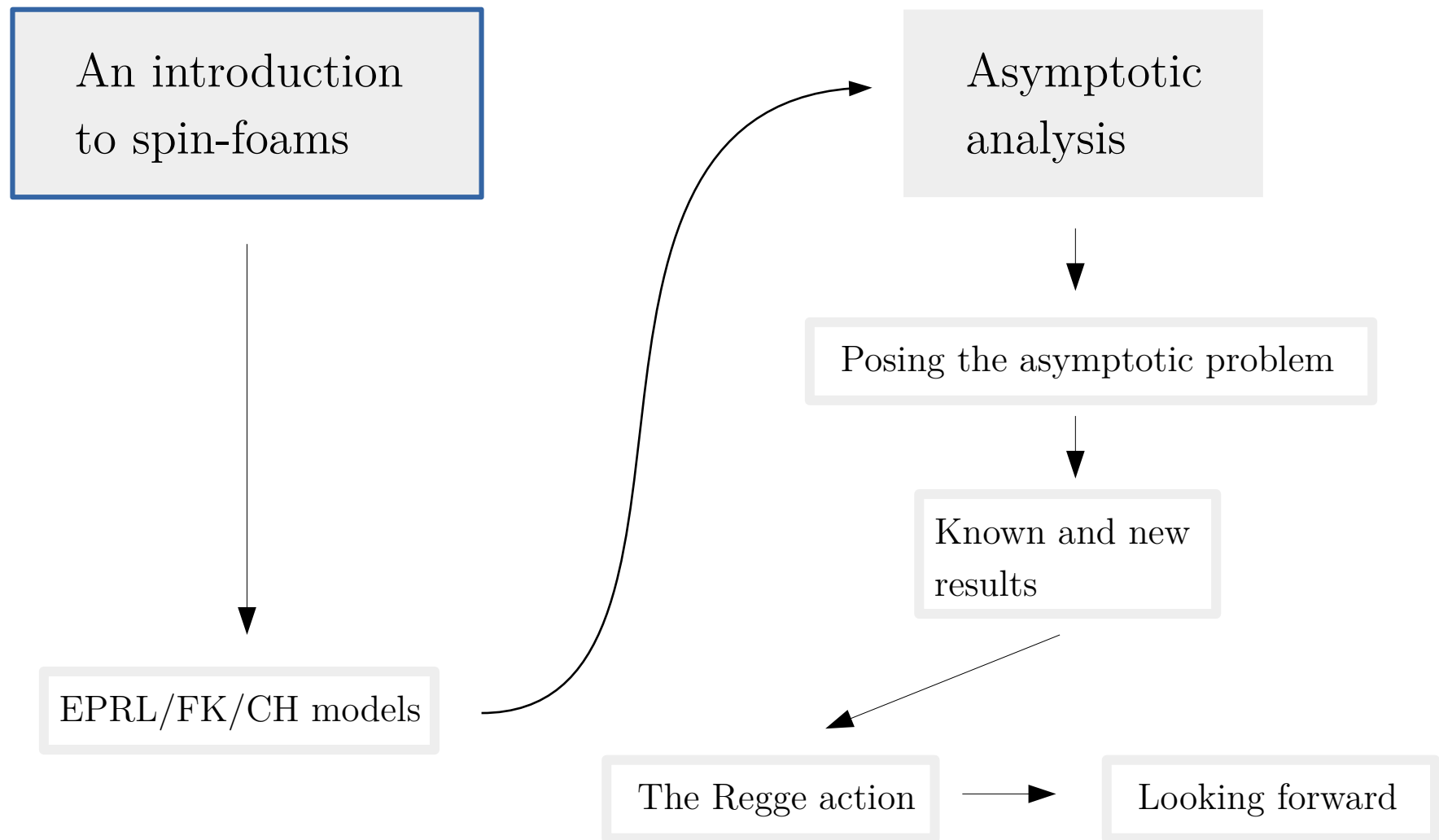
Lorentzian space-times

A possible way forward:
the Spin-foam framework

The road-map



The road-map



- GR = BF + constraints

$$(g = \eta^{IJ} \theta_I \otimes \theta_J)$$

$$S_T = \int_M F^{IJ} \wedge \left(\star + \frac{1}{\gamma} \right) \theta_I \wedge \theta_J$$

$$\int_M \text{Tr } B \wedge F,$$

“Simplicity constraints”
 $B = (\star + \gamma^{-1}) \theta \wedge \theta$

- The general idea: (Engle et al. ‘07; Freidel, Krasnov ‘07; Conrady, Hnibyda ‘10)

Discretize BF



Quantize



Apply constraints
at quantum level

EPRL/FK/CH spin foam model

- GR = BF + constraints

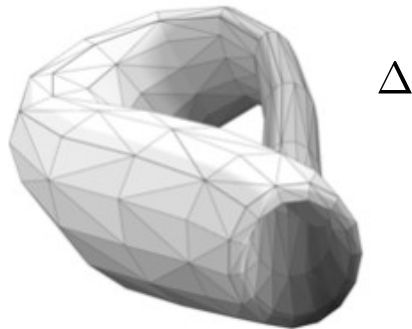
$$S_T = \int_M F^{IJ} \wedge \left(\star + \frac{1}{\gamma} \right) \theta_I \wedge \theta_J \quad \int_M \text{Tr } B \wedge F, \quad \text{“Simplicity constraints”}$$

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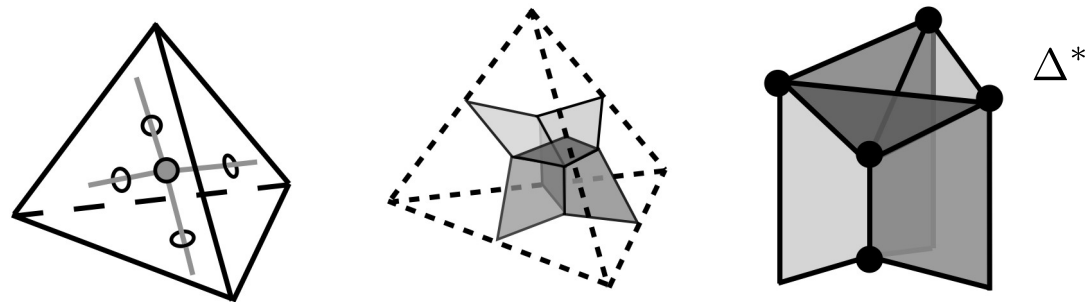
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Cellular decomposition of M



Dual 2d cell complex

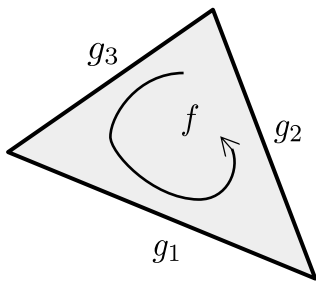
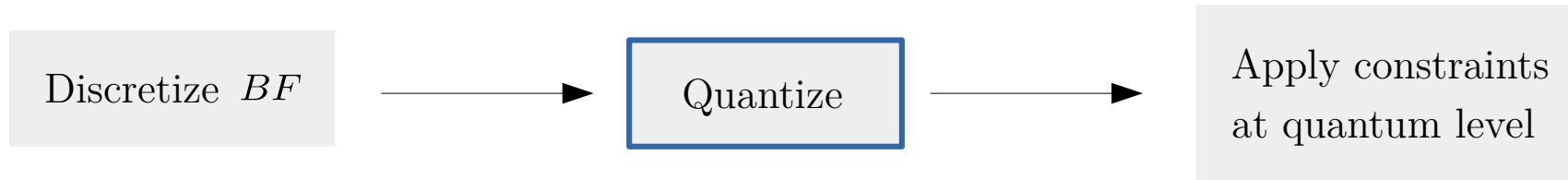


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$$Z(M) \sim \int \mathcal{D}A \mathcal{D}B e^{i \int_M \text{Tr} B \wedge F} \sim \int \mathcal{D}A \delta(F)$$

$$Z(\Delta^*) = \int_{\text{SL}(2, \mathbb{C})} \prod_{e \in \mathcal{E}} dg_e \prod_{f \in \mathcal{F}} \delta \left(\prod_{e \in \partial f} g_e \right)$$

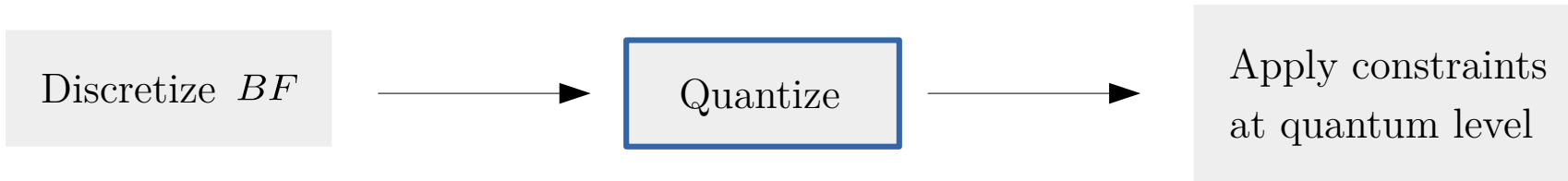
EPRL/FK/CH spin foam model

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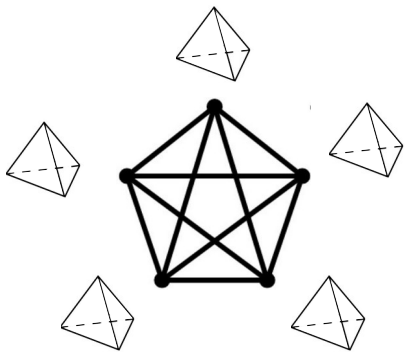
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SL(2,C) unitary irreducible reps.




$$Z(\Delta^*) = \sum_{\{\rho\} \rightarrow \{f\}} \prod_{f \in \mathcal{F}} \dim(\rho) \prod_{v \in \mathcal{V}} \text{[Diagram of a vertex with four faces and associated representations]}$$

$$\begin{aligned} | &= \rho(g) \\ \text{[Diagram of a vertex with four faces]} &= \int dg \rho_1(g) \dots \rho_n(g) \end{aligned}$$

EPRL/FK/CH spin foam model

- The general idea:



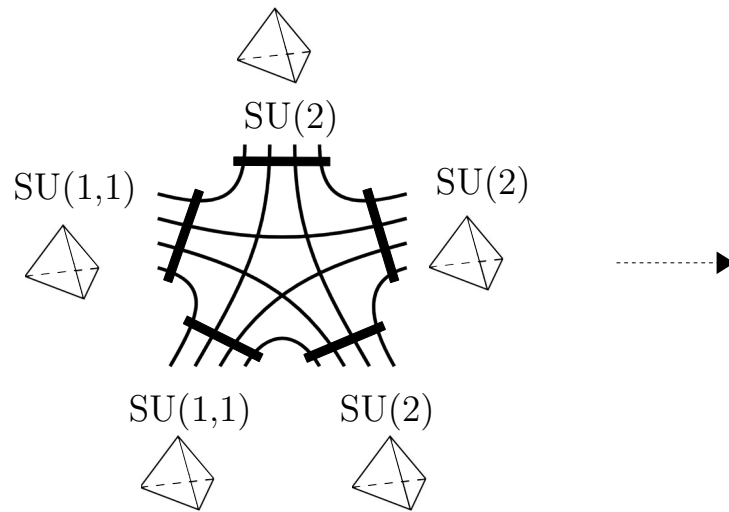
- Discretized constraints require choice of causal character of each ;
- Completeness relations from $SU(2)$ and $SU(1,1)$ unitary irreps.

$$\mathbb{1}_\rho = \sum_{j m} |j m\rangle \langle j m|$$

$SU(2)$, space-like

$$\mathbb{1}_\rho = \sum_{k m} |k m\rangle \langle k m| + \sum_m \int ds |s m\rangle \langle s m|$$

$SU(1,1)$, time-like




- Constraints single out a set of admissible representations $\rho^{(p,n)}$ depending on causal character
- $$n = f(p)$$

EPRL/FK/CH spin foam model

- The general idea:



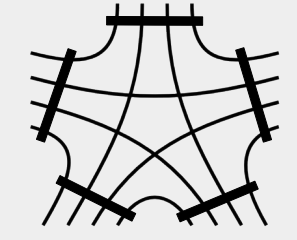
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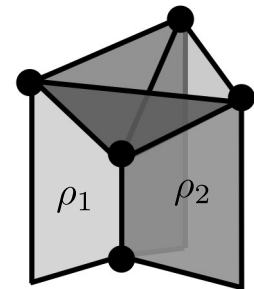
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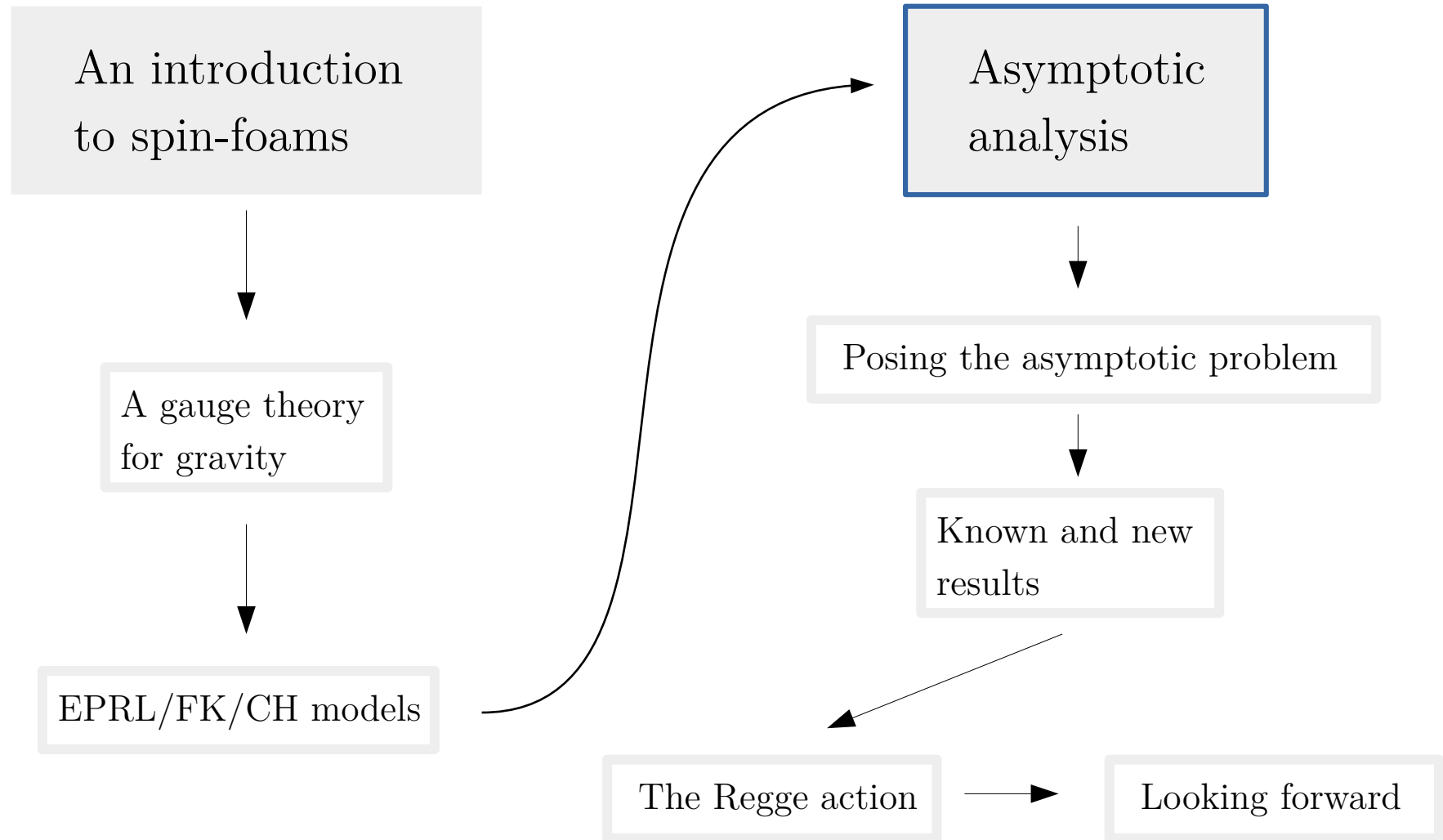
$SU(1,1)$, time-like

$$Z(\Delta^*) = \sum_{\{\rho\} \rightarrow \{f\}} \prod_{f \in \mathcal{F}} \dim(\rho) \prod_{v \in \mathcal{V}}$$


Restrict to admissible representations

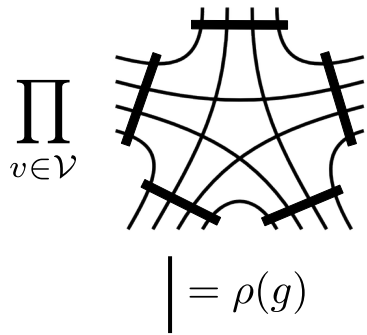


The road-map



Posing the asymptotic problem

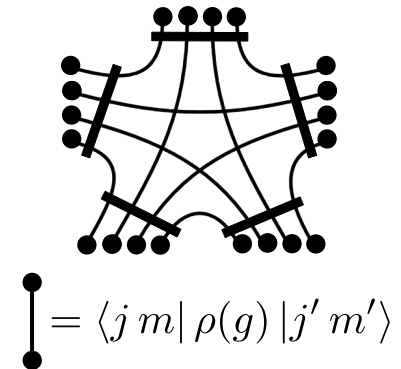
- States of the theory live in spaces of unitary irreps. of $SU(2)$ and $SU(1,1)$, labeled by spins $\{|j m\rangle, |k, m\rangle, |s m\rangle\}$;



$$\mathbb{1}_\rho = \sum_{j m} |j m\rangle \langle j m|$$

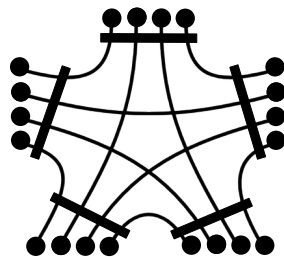
$$\mathbb{1}_\rho = \sum_{k m} |k m\rangle \langle k m| + \sum_m \int ds |s m\rangle \langle s m|$$

Vertex amplitude



- What happens at large j, k and s ?
- Focus on vertex amplitude, take uniform scaling $\{j, k, s\} \rightarrow \{\Lambda j, \Lambda k, \Lambda s\}$, and rewrite:

Fix boundary states



$$= \int_{\text{SL}(2, \mathbb{C})} \left[\prod_i dg_i \right] f(\{g_i\}) e^{i\Lambda S(\{g_i\})}$$

Look at critical points of action

- Asymptotic analysis for entirely space-like foams done in (Barret et al. '09);

$$S|_{\text{critical}} = S_{\text{Regge}}$$

$$A \sim e^{iS_{\text{Regge}}} + e^{-iS_{\text{Regge}}} \quad \text{“cosine problem”}$$

- Inclusion of time-like tetrahedra, but only space-like triangle interfaces in (Kamiński et al. '17);
- Extension to time-like triangles by (Liu, Han '18);

Unanswered questions in time-like analyses

- Minkowski's theorem and rigidity for Lorentzian polyhedra?
- Extension to non-simplicial cells?
- Really need to assume at least one time-like triangle in each time-like tetrahedron?
(Liu, Han '18)
- Explicit asymptotic formula?

Known and new results

- Asymptotic analysis for entirely space-like foams done in (Barret et al. '09);

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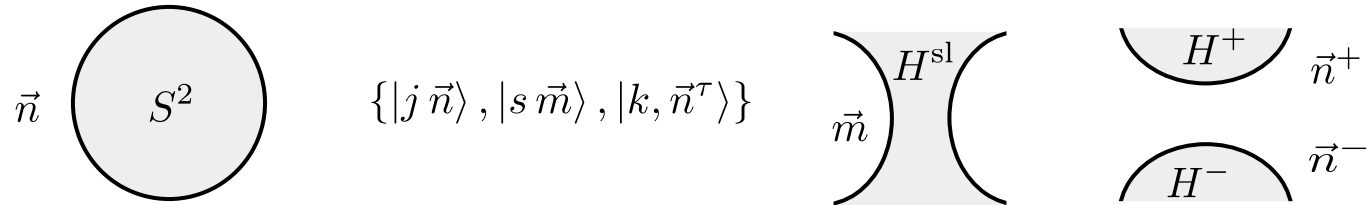
- Minkowski's theorem and rigidity for Lorentzian polyhedra? ✓, **holds**
- Extension to non-simplicial cells? ✓
- Really need to assume at least one time-like triangle in each time-like tetrahedron? ✓, **no!**
- Explicit asymptotic formula? **wip**

+ No “cosine problem” if 4-cell includes both time-like and space-like cells.

(J.D.S., Steinhaus, soon)

Critical points of the action

- Boundary states can be characterized by a spin and by a vector in S^2 , H^\pm , H^{sl} .



- After a lot of massaging,

$$\sum_f j_f \vec{n}_f = 0$$

Space-like 3-cells

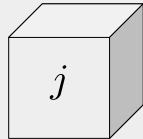
$$\sum_{s.l.f} k_f \vec{n}_f^T + \sum_{t.l.f} s_f \vec{m}_f = 0.$$

Time-like 3-cells

- Minkowski's theorem:** there exists a unique polyhedron satisfying any given closure condition

$$\sum_f A_f \vec{n}_f = 0 \quad (\text{Aleksandrov 1950})$$

with A_f the areas of faces and $\{n_f\}$ a set of l.i. normals to the faces.

Can reconstruct  with spin areas

Critical points of the action

- One more set of critical point equations, prescribing **gluing rules for faces of polyhedra**,



following combinatorics of cellular decomposition;

- One obtains (under well-behaved data) a 4-dimensional polytope with curvature at the boundary;
- The asymptotic action generally becomes the Regge action associated with the polytope,

$$S \sim \sum_f \theta_f A_f \quad (\text{discrete curvature})$$

with θ_f the Euclidian/Lorentzian dihedral angle between normals of polyhedra sharing a face.

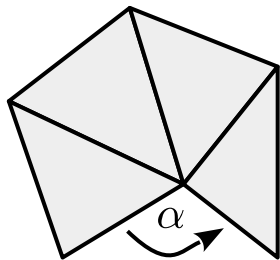
The Regge action

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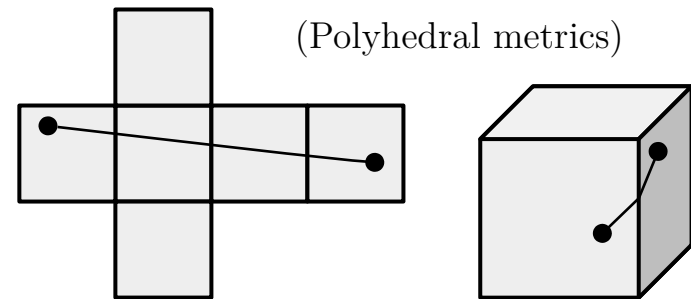
$$ds^2 = dr^2 + \left(1 - \frac{\alpha}{2\pi}\right)d\phi^2 - dt^2 + dz^2$$



(Sorkin 1974)

$$\frac{\sqrt{-g}}{2} R = \frac{\alpha \delta(r)}{2\pi}$$

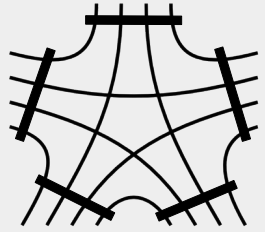
$$\frac{1}{2} \int \sqrt{-g} R \, dr d\phi dz dt = \alpha A$$



(Polyhedral metrics)

(Aleksandrov 1950)

- The spin-foam frameworks provides a **well-defined notion of a sum-over-histories** for gravity;
- The model is structured around **combinatorics** and **algebraic data** – geometry is *diluted*;

$$Z(\Delta^*) = \sum_{\{\rho\} \rightarrow \{f\}} \prod_{f \in \mathcal{F}} \dim(\rho) \prod_{v \in \mathcal{V}}$$


- Triangulated gravity is found in the large-spin regime, where one recovers **geometry**;

Many open questions: inclusion of matter, cellular-decomposition dependence, continuum limit,...

Looking forward

- States associated to time-like 2-cells are generalized eigenstates of a non-compact operator, admitting complex eigenvalues;
 - Possible inclusion of causal ordering between time-like cells?
 - Simpler parametrization of the model?
- Numerical analysis of renormalization flow using hypercuboids;
 - Can one find a phase transition signaling lattice independence?
- Inclusion of matter fields at the foam level;
- Operational meaning of the framework;
- ...?

Thank you for your attention!