Introduction to spin foams and background independent renormalization

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Gravity and quantum matter?

- General relativity \implies dynamical space-time
 - **Curvature** of space-time \iff gravitation
 - Matter curves space-time, curvature determines motion of matter
- Quantum matter (quantum field theory)
 - Visible matter content of universe
 - Space-time acts as a fixed "stage"



• Opposing assumptions, but theories relevant at different scales?

What happens at smallest length scales?

Quantum matter on classical space-time **inconsistent**.

Quantum nature of gravity cannot be ignored!

"Zooming in" to the Planck scale

• What is space-time at the **Planck scale** $l_P = \sqrt{\frac{\hbar G}{c^3}}$?



- Bottom-up approach: ansatz for UV physics
- Use renormalization group as "microscope"
 - "Zoom out" by relating theories at different scales

Replace space-time by a "quantum space-time".

One promising approach: spin foam quantum gravity

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Path integral for gravity



• Formal path integral of geometries

$$Z = \int \mathcal{D}[g] \; e^{\frac{i}{\hbar}S_{\rm EH}[g]}$$

- $S_{\rm EH}$: Einstein-Hilbert action
- $\mathcal{D}[g]$: Measure on space of geometries
- Regulator: Lattice
 - Diffeomorphism symmetry?
 - Continuum limit?
 - Renormalizability?

Spin foam quantum gravity:

Background independent, non-perturbative quantization of gravity

Outline



- 2 Ponzano Regge Model: Spin foams in 3D
- 3 4D spin foam models
- 4 Background independent renormalization
- **5** Summary and Outlook

Riemannian continuum gravity in 3D

- Reviews on spin foams and BF theory [Baez '99, Perez '12]
- \bullet Compact, orientable 3D Riemannian manifold ${\cal M}$
- 3D gravity described by (topological) SU(2) BF theory:

$$S[B,\omega] = \int_{\mathcal{M}} \operatorname{Tr}[B \wedge F(\omega)]$$

- $B \ \mathfrak{su}(2)$ valued 1-form, F curvature 2-form of connection ω
- ω connection on SU(2) principal bundle
- Equations of motion:

$$F(\omega) = 0$$
 and $d_{\omega}B = 0$

• Path integral:

$$Z = \int \mathcal{D}[B]\mathcal{D}[\omega] \ e^{i\int_{\mathcal{M}} \operatorname{Tr}[B \wedge F(\omega)]} \ " = \int \mathcal{D}[\omega] \ \delta(F(\omega))"$$

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Triangulation and its dual complex

• Triangulation Δ (of \mathcal{M}) and its dual Δ^*

- $\bullet \ v \in \Delta^* \ \leftrightarrow \ \tau \in \Delta$
- $\bullet \ e \in \Delta^* \ \leftrightarrow \ T \in \Delta$
- $\bullet \ f \in \Delta^* \ \leftrightarrow \ E \in \Delta$



Discretizing BF theory

• Discretize *B*-field:

•
$$(d-2)$$
-dim face $f \in \Delta^*$

$$B_f := \int_{1-\mathrm{cell}} B$$

 $\bullet~B\mbox{-field}$ smeared over an edge

Edge vector

- **Discretize** connection ω :
 - 1-dim edge $e \in \Delta^*$

$$g_e := \mathcal{P} \exp\left(-\int_e \omega\right)$$

• $g_e \in \mathrm{SU}(2)$: parallel transport along e

Parallel transport to another tetrahedron



Discretized curvature and partition function

- **Discretize** curvature $F(\omega)$:
 - Holonomy around closed loop $f \in \Delta^*$:

$$U_f := \vec{\prod}_{e \subset f} g_e$$

- Matching orienation assumend
- Partition function:

$$Z(\Delta) = \int \prod_{e \in \Delta^*} dg_e \prod_{f \in \Delta^*} dB_f \, e^{iB_f U_f} = \int \prod_{e \in \Delta^*} dg_e \prod_{f \in \Delta^*} \delta\left(\vec{\prod}_{e \subset f} g_e\right)$$

Curvature measured around edges $E \in \Delta$.

Flatness imposed: closed holonomies must be equal to identity.



Toward the spin foam representation

- Peter-Weyl's theorem: $\delta(g) = \sum_j d_j \operatorname{Tr}[D^j(g)]$
 - $j \in \frac{\mathbb{N}}{2}$: labels irrep of SU(2)
 - D_{mn}^j : Wigner-D matrix for irrep j
- **Expand** each δ -function in Z:
 - Irrep j_f per face $f \in \Delta^*$
 - Use $D^{j}(g_1 g_2) = D^{j}(g_1)D^{j}(g_2)$

$$Z(\Delta) = \sum_{j_f} \int \prod_{e \in \Delta^*} dg_e \prod_{f \in \Delta^*} d_{j_f} \operatorname{Tr}[D^{j_f}(g_{e_1}) D^{j_f}(g_{e_2}) \dots D^{j_f}(g_{e_N})]$$
$$= \sum_{j_f} \prod_{f \in \Delta^*} d_{j_f} \operatorname{Tr} \underbrace{\left[\int \prod_{e \in \Delta^*} dg_e \bigotimes_{f \supset e} D^{j_f}(g_e) \right]}_{\mathbf{P}_{\operatorname{Haar}}(\{j_f\})}$$

 \mathbf{P}_{Haar} : Haar-projector onto invariant subspace $\text{Inv}(\bigotimes_f j_f)$.

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Towards spin networks

- 3D triangulation: $e \in \Delta^*$ is 3-valent
 - Triangle is bound by 3 edges
 - $e \in \Delta^*$ part of 3 faces $f \in \Delta^*$

$$\mathbf{P}_{\text{Haar}} = \int dg \, D^{j_{f_1}}(g) \otimes D^{j_{f_2}}(g) \otimes D^{j_{f_3}}(g) = \bigcup_{i=1}^{g} |i\rangle \langle i|$$

• Intertwiner ι : normalized unique vector in $Inv(j_1 \otimes j_2 \otimes j_3)$



Ponzano-Regge model

• Ingredients:

- Irrep j_f : length of edge E (dual to f)
- Intertwiner ι_e : "shape" of triangle (dual to e)
- Final form of **partition function**:

$$Z(\Delta) = \sum_{j_f} \prod_{f \in \Delta^*} d_{j_f} \prod_{e \in \Delta^*} (-1)^{\cdots} \prod_{v \in \Delta^*} \underbrace{j_1 \qquad j_2 \qquad j_3 \qquad j_5 \qquad j_4}^{j_1 \qquad j_6}$$

- Local assignment of amplitudes
 - Vertex amplitude: evaluation of spin network function
 - $\{6j\}$ -symbol: "quantum tetrahedron"

Sum over quantum geometric building blocks

Derived from (constrained) topological quantum field theory

Properties of Ponzano-Regge model



• Triangulation independent

- Invariant under Pachner moves
- Discrete remnant of diffeomorphism symmetry [Freidel, Louapre '02]
- Turaev-Viro model: $\Lambda > 0$ [Turaev, Viro '92]
 - Quantum deformed $SU(2)_k$

•
$$j_{\text{max}} = \frac{\mathbf{k}}{2}$$

• $\Lambda = \frac{1}{G^2 \hbar^2 \mathbf{k}^2}$

• Asymptotic expansion of vertex amplitude

[Ponzano, Regge '68]

$$\underbrace{\frac{j_1}{j_2}}_{j_3} \underbrace{\frac{j_6}{j_5}}_{j_4} \sim \frac{1}{\sqrt{12\pi V}} \cos\left(S_{\rm R} + \frac{\pi}{4}\right)$$

 $S_{\rm R}:$ Regge / discrete gravity action for a single tetrahedron



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Plebanski-Holst action

• 4D gravity as a constrained BF theory:

$$S[B,\omega,\lambda] = \frac{1}{\kappa} \int_{\mathcal{M}} \left[\left(*B^{IJ} + \frac{1}{\gamma} B^{IJ} \right) \wedge F_{IJ}(\omega) + \lambda_{IJKL} B^{IJ} \wedge B^{KL} \right]$$

- $\lambda_{IJKL} = -\lambda_{JIKL} = -\lambda_{IJLK} = \lambda_{KLIJ}$ Lagrange multipliers
- $\rightarrow B = \pm * (e \wedge e)$ or $B = \pm e \wedge e$
- γ : Immirzi parameter
- Compare to **BF Yang-Mills**:

$$S_{\mathrm{BFYM}}[B,A] = \int_{\mathcal{M}} \mathrm{Tr}(B \wedge F(A)) + g^2 \int_{\mathcal{M}} \mathrm{Tr}(B \wedge \star B)$$

BF theory corresponds to Yang-Mills theory in **weak coupling limit** Understand spin foam models as **generalized lattice gauge theories**

4D spin foam strategy

- As before, discretize and quantise $SU(2) \times SU(2)$ BF theory
 - B_f associated to faces $f \in \Delta^*$ (dual to triangles)
 - g_e associated to edges $e \in \Delta^*$ (dual to tetrahedra)



• Interpretation of variables:

- j_f : area of triangle (dual to f)
- ι_e : shape of tetrahedron (dual to e)

Different 4D spin foam models

- Implement simplicity constraints to break topological symmetry
 - Restrictions on irreps (j^+, j^-) and intertwiners ι
- Barrett-Crane model [Barrett, Crane '98]
 - Simple bivector $\rightarrow j^+ = j^-$
- EPRL-FK model [Engle, Pereira, Rovelli, Livine '08, Freidel, Krasnov '08]
 - Weak imposition of constraints (in discrete):

$$j^{\pm} = \frac{1}{2} \left| 1 \pm \gamma \right| j$$

• Contact with LQG Hilbert space (spin networks)

Simplicity constraints: restrictions on representations and intertwiners

Spin foams: interpretation and results



- Amplitude functional $\mathcal{A} : \mathcal{H}_b \to \mathbb{C}$
 - \mathcal{H}_b : boundary Hilbert space on a discretisation b
- Transition amplitude for $\mathcal{H}_b = \mathcal{H}_i \otimes \mathcal{H}_f^*$

•
$$\langle s'|s \rangle_{\mathcal{A}} = \mathcal{A}(s \otimes s')$$

- Evolution of 3D geometries
- Asymptotic expansion of vertex amplitude [Conrady, Freidel '08, Barrett, Dowdall, Fairbarn, Gomes, Hellmann

'09]

$$\mathcal{A}_v \sim \cos(S_{\mathrm{R}}) + \dots$$

• Implementation of Cosmological Constant [Han '11, Fairbarn, Meusburger '11, Haggard, Han, Kaminski, Riello '15]

At discrete level: relation to classical gravity

Can we go **beyond** a **few simplices**?

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Why we need renormalization

- Definition **not unique**
- Lattice is a choice.
 - Should the lattice be "fine" or "coarse"?
 - Can we **relate** the **physics** on different lattices?
 - Diffeomorphism symmetry? [Bahr, S.St. '15]



Physics must be independent of choice of regulator!

- How do we bridge the enormous gap to observable physics?
- Do we recover **general relativity** in a suitable limit?

Relating theories across discretisations

• Wilsonian renormalization in absence of a scale

- No background geometry
- Use the **lattice** itself as a **relative scale**
 - "Fine" vs. "coarse" degrees of freedom
- Coarse grain by summing over "fine" degrees of freedom
 - E.g. decimation procedure
 - Regge calculus for $\Lambda \neq 0$ in 3D [Bahr, Dittrich '09]
- Interactiong theories: non-local interactions arise
 - Example: Ising model in 2D
 - Difficult to iterate

Allow for more general boundary data

Complex building blocks that interact locally

Consistent boundary formulation I $_{[Dittrich, S.St. '14]}$

• Recall: amplitude functionals $\mathcal{A}_b : \mathcal{H}_b \to \mathbb{C}$

- Simplest example: single simplex
- Different boundaries $b, b' \to \text{different } \mathcal{H}_b, \mathcal{H}_{b'}$
- How can we compare \mathcal{A}_b and $\mathcal{A}'_{b'}$?



- Idea: relate states in \mathcal{H}_b to states $\mathcal{H}_{b'}$
 - Embedding maps: $\iota_{b'b} : \mathcal{H}_b \to \mathcal{H}_{b'}$
 - Boundaries partially ordered set $b \prec b'$

Represent the same state on different Hilbert spaces

$$\mathcal{H}_b \ni \psi_b \to \iota_{b'b}(\psi_b) \in \mathcal{H}_{b'}$$

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Consistent boundary formulation II $_{[Dittrich, S.St. '14]}$





Assign different amplitudes to different lattices Renormalization group flow of spin foam amplitudes

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Tensor network renormalization $_{[{\rm Levin}, \ Nave \ '07]}$



Study flow of tensors under coarse graining

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Results from Tensor networks



• Numerical algorithm

- Works for oscillating amplitudes
- Finite dimensional Hilbert spaces
- 2D analogue models for SU(2)_k [Dittrich, Martin-Benito, S.St. '13, Cameron, Dittrich, Schnetter, S.St. '16]
- 3D lattice gauge theories for \mathbb{Z}_2 and S_3 [Dittrich, Mizera, S.St. '14, Delcamp, Dittrich '16]
- Fusion basis for 3D $SU(2)_k$ lattice gauge theory [Dittrich, Geiller '16, Cunningham, Dittrich, S.St. '20]
 - Phase transition: **confining** and **deconfing** phase
 - Critical g_c drops as we increase k.
 - Expectation values of Wilson loops

Hint towards no deconfining phase for continuous group SU(2).

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Results from 4D restricted spin foam models

• Complementary route: Study restricted path integral [Bahr, S.St. '15, Bahr, Klöser, Rabuffo '17]



- Proof of principle [Bahr, S.St. PRL '16]
 - 4D restricted spin foam model
 - Several simplifications / assumptions

Indication for a **UV-attractive fixed point** and restored **diffeomorphism symmetry**

- Diffusion process on space-time
 - Spectral dimension D_S of restricted spin foam [S.St., Thürigen '18]
 - D_S ≤ 4 due to superposition of geometries



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Summary

• Spin foam quantisation strategy

- Discretisation and quantisation of topological BF theory
- Derivation of spin foam representation
- 4D spin foam models: simplicity constraint
- Quantum geometric building blocks
- Background independent renormalization
 - Tackle open questions (diffeomorphism symmetry, discretisation dependence,...)
 - Assign a **family of amplitudes** across lattices
- Example: Tensor network renormalization
 - Fusion basis for 3D lattice gauge theories

Promising results in this line of research

Intriguing open questions remains

Outlook

Renormalization

- $\bullet\,$ Apply tensor networks (or simplified algorithms) to ${\bf 4D}$ spin foams
- Which other numerical methods useful (Monte Carlo)?

• Numerical methods

- Develop algorithms to compute spin foam amplitudes
- Hybrid algorithm, using semi-classical results

• Matter and observables

- Couple matter to spin foams
- Renormalize both gravity and matter
- Geometric observables (spectral dimension, curvature)

Thank you for your attention!